L^p boundedness of maximal averages over hypersurfaces in $\mathbf{R}^3$

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1. Introduction and statement of results

Let Q be a smooth hypersurface in $\mathbf{R}^3$ and let q_0 be a point on Q. Let $d\sigma(q)$ denote the standard surface measure on Q. For a small cutoff function ϕ supported near q_0, we define the maximal operator M, initially on Schwarz functions, by

$$Mf(x) = \sup_{t>0} \left| \int_Q f(x - tq)\phi(q)\, d\sigma(q) \right| \tag{1.1}$$

Our goal is to determine for which p is the maximal operator M bounded on L^p. Note that by subadditivity of maximal operators of the form (1.1), one can prove L^p boundedness of the nonlocalized analogue of (1.1) over a compact surface by doing a partition of unity and reducing to (1.1). If L is an invertible linear transformation and M_L denotes the maximal operator corresponding to the surface $L(Q)$, then one can easily check from the definitions that $M_L f(x) = |det(L)| M(f \circ L)(L^{-1}x)$. Hence when studying L^p boundedness properties one may always replace M by M_L. In particular, we may assume that near q_0, Q is given as the graph of a function $g(x, y)$ such that if (x_0, y_0) denotes the point in the x-y plane which q_0 lies above, then $\nabla g(x_0, y_0) = 0$.

The earliest work in this area was done in the case where Q is an n-dimensional sphere in $\mathbf{R}^{n+1}$, when Stein [St1] showed M is bounded on L^p iff $p > \frac{n+1}{n}$ for $n > 1$. This was later generalized by Greenleaf [Gr] to surfaces of nonvanishing Gaussian curvature, with some further results for when the Hessian has rank between 1 and n. The $n = 1$ case was later proven by Bourgain [B]. In [So] Sogge showed in any dimension that whenever Q has at least one nonvanishing principal curvature, M is bounded on L^p for all $p > 2$. The case of convex surfaces of finite line type has been extensively analyzed; we refer to [IoSe2] and [NaSeWa] for more information on these situations.

Although there are many interesting issues when $p \leq 2$, for the purposes of this paper we always assume $p > 2$. M is trivially bounded on L^∞, and if M is bounded on some L^p, by interpolating with the L^∞ case one has that M is bounded on $L^{p'}$ for $p' > p$. Hence our goal is to determine the optimal $p_0 \geq 2$ for which M is bounded on L^p for $p > p_0$. If Q is tangent to the tangent plane $T_{q_0}(Q)$ to infinite order at q_0, then as long as

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$0 \notin T_{q_0}(Q)$ a relatively straightforward argument shows that M is unbounded on L^p for all finite p. Conversely, if the Gaussian curvature of Q does not vanish to infinite order at q_0, then by [SoSt] M is bounded on L^p for some finite p. (See [CoMa] for another result of this kind.) It is entirely possible that in any dimension, M is bounded on some finite L^p whenever Q is not tangent to $T_{q_0}(Q)$ to infinite order at q_0. Hence in general we expect p_0 to be finite.

Definition 1.1. Let $d(x, y)$ denote the vertical distance between Q and $T_{q_0}(Q)$ above (x, y). The *height* $h(q_0)$ is defined to be the reciprocal of the supremum of all ϵ for which the integral of $|d(x, y)|^{-\epsilon}$ is finite on at least one neighborhood of q_0.

For the case $n = 2$ considered in this paper, a good L^p boundedness theorem for $p > 2$ was proven in [IkKeMu2]. Their main theorem can be stated as follows.

Theorem [IkKeMu2]. Suppose the origin is not contained in $T_{q_0}(Q)$. If $\phi(q)$ is supported on a sufficiently small neighborhood of q_0 then M is bounded on L^p for $p > \max(h(q_0), 2)$. When $h(q_0) \geq 2$ and $\phi(q_0) \neq 0$ this exponent is sharp in that M is unbounded on L^p for $p < h(q_0)$ and if Q is real-analytic, then M is unbounded on $L^{h(q_0)}$ as well.

The purpose of this paper is to provide a relatively short alternative approach to the $p > 2$, $n = 2$ situation by extending the methods of [G1] and using facts about the adapted coordinate systems described below. There will once again be exceptional cases not covered, but due to the differences in the methods the exceptional cases will be quite different from and not mutually exclusive to the exceptional cases of [IkKeMu2], which occur when $0 \in T_{q_0}(Q)$.

Newton Polygons and Adapted Coordinates.

We now give some relevant terminology which will be used throughout this paper. Below, $R(x, y)$ denotes a smooth function defined on a neighborhood of the origin with nonvanishing Taylor expansion at the origin.

Definition 1.2. Let $R(x, y) = \sum_{a,b} R_{ab} x^a y^b$ denote the Taylor expansion of $R(x, y)$ at the origin. For any (a, b) for which $R_{ab} \neq 0$, let Q_{ab} be the quadrant $\{(x, y) \in \mathbf{R}^2 : x \geq a, y \geq b\}$. Then the *Newton polygon* $N(R)$ of $R(x, y)$ is defined to be the convex hull of the union of all Q_{ab}.

In general, a Newton polygon consists of finitely many (possibly zero) bounded edges of negative slope as well as an unbounded vertical ray and an unbounded horizontal ray.

Definition 1.3. The *Newton distance* $d(R)$ of $R(x, y)$ is defined to be $\inf\{t : (t, t) \in N(R)\}$.

One often uses (t_1, t_2) coordinates to write equations of lines relating to Newton polygons, so as to distinguish from the x-y variables of the domain of $R(x,y)$. The line in the t_1-t_2 plane with equation $t_1 = t_2$ comes up so frequently it has its own name:

Definition 1.4. The *bisectrix* is the line in the t_1-t_2 plane with equation $t_1 = t_2$.

A key role in the above theorems as well as our theorems to follow is played by the following polynomials.

Definition 1.5. Suppose e is a compact edge of $N(R)$. Define $R_e(x,y)$ by $R_e(x,y) = \sum_{(a,b)\in e} R_{ab}x^a y^b$. In other words $R_e(x,y)$ is the sum of the terms of the Taylor expansion of f corresponding to $(a,b) \in e$.

Definition 1.6. Suppose $R(x,y)$ has nonvanishing Taylor expansion at the origin such that $R(0,0) = 0$ and $\nabla R(0,0) = 0$. Then $R(x,y)$ is said to be in *nonadapted coordinates* if the bisectrix intersects $N(R)$ in the interior of a compact edge e of $N(R)$ such that $R_e(1,y)$ has a zero of order greater than $d(R)$. If $R(x,y)$ is not in nonadapted coordinates, then $R(x,y)$ is said to be *adapted coordinates.*

The significance of adapted and nonadapted coordinates was first discovered by Varchenko [V] for the real-analytic case and for the general smooth case by Ikromov-Müller [IkMu]. Namely, define ϵ_0 to be the supremum of all ϵ for which $|R|^{-\epsilon}$ is integrable in at least one neighborhood of the origin. Equivalently, ϵ_0 is the supremum of the epsilon such that on some neighborhood of $(0,0)$ one has $|\{(x,y) : |R(x,y)| < t\}| < Ct^\epsilon$ for some C. Then by [V] and [IkMu] one has $d(R) \leq \frac{1}{\epsilon_0}$, with equality holding if and only if $R(x,y)$ is in adapted coordinates. Furthermore, one has the following. Suppose $R(x,y)$ is not in adapted coordinates and let e be the edge of $N(R)$ intersecting the bisectrix in its interior. By [V] and [IkMu], if the slope m_e of e is at least -1, then there is a smooth $\psi(x)$ with $\psi(0) = 0$ such that $R(x, y+\psi(x))$ is in adapted coordinates. By switching the roles of the x and y axes, this means that if $m_e \leq -1$, there is a smooth $\psi(y)$ with $\psi(0) = 0$ such that $R(x+\psi(y), y)$ is in adapted coordinates. Thus [V] and [IkMu] show that there necessarily is a "nice" coordinate change after which the growth rate of R at the origin is given in the above way by $\frac{1}{d(R)}$ and where R is in adapted coordinates; these facts in turn are used in [IkKeMu2] in their proof of the L^p boundedness properties of M in their theorem above. Another useful aspect of adapted coordinates proven in [V] and [IkMu] is that if one is in adapted coordinates, then the order of any zero of any $R_e(1,y)$ (cf. Definition 1.5) is at most $d(R)$.

Let a denote the order of the zero of $R(x,y)$ at $(0,0)$. Then for a generic linear transformation T, $\partial_y^a(R\circ T)(0,0)$ and $\partial_x^a(R\circ T)(0,0)$ are both nonzero. Thus the Newton polygon of $R\circ T$ is entirely on or above the line $t_1 + t_2 = a$. But no point of $N(R\circ T)$ can be below this line; otherwise $r \circ T$ would have a zero of order less than a at the origin. We conclude that $N(R\circ T)$ is exactly $\{(t_1,t_2) \in \mathbf{R}^2 : t_1 + t_2 \geq a\}$. Since the compact edge of $N(R\circ T)$ has slope -1, by the above discussion there is a $\psi(x)$ such that $R\circ T(x, y+\psi(x))$

is in adapted coordinates. Note that the Newton polygon of $R \circ T(x, y+\psi(x))$ still has its upper vertex at $(0, a)$ and that the slope of each edge of this Newton polygon is at least -1. This motivates the following definition.

Definition 1.7. Suppose $R(0,0) = 0$ and $\nabla R(0,0) = 0$. Then $R(x, y)$ is said to be in *generic adapted coordinates* if $R(x, y)$ is in adapted coordinates, each edge of $N(R)$ has slope at least -1, and $N(R)$ intersects the y-axis at some point $(0, a)$.

Note that the above definition implies that $R(x, y)$ has a zero of order a at the origin. We now are in a position to state the main theorem of this paper. Recall we are working in the situation where Q is a surface with a distinguished point $q_0 = (x_0, y_0, z_0)$, such that near q_0 the surface Q is the graph of some $g(x, y)$ with $g(x_0, y_0) = z_0$ and $\nabla g(x_0, y_0) = (0, 0)$. Let $f(x, y) = g(x + x_0, y + y_0) - z_0$, and define $D(x, y)$ to be the Hessian determinant of f at (x, y). Then our main theorem is as follows

Theorem 1.1. Suppose M is as defined in (1.1). If $\phi(q)$ is supported on a sufficiently small neighborhood of q_0, then M is bounded on L^p for $p > \max(h(q_0), 2)$ as long as neither of the following two exceptional situations occurs.

a) $D(x, y)$ has a zero of infinite order at $(0, 0)$.

b) Whenever T is an invertible linear transformation and $\psi(x)$ a smooth function with a zero of order $b > 0$ at the origin such that $F(x, y) = (f \circ T)(x, y + \psi(x))$ is in generic adapted coordinates, then the bisectrix intersects $N(F)$ in the interior of a compact edge e with slope m_e with $|m_e| < \frac{1}{b}$ such that $p(y) = \partial_y(F_e(1, y))$ or $\partial_y(F_e(-1, y))$ has a zero of order greater than $\max(1, d(F) - 1)$ at some y_0 for which $p(y_0) \neq 0$.

Here $F_e(x, y)$ is as in Definition 1.5. If $\psi(x)$ has a zero of infinite order at $x = 0$, then we take $\frac{1}{b} = 0$ in Theorem 1.1b); in other words, the exceptional condition of b) will not be satisfied. As is explained in [IkKeMu2], using a theorem in [IoSa1] one can show the exponent $h(q_0)$ is best possible when $q_0 \notin T_{q_0}(Q)$, assuming $\phi(q_0) \neq 0$. If $q_0 \in T_{q_0}(Q)$, then sometimes one can do better as the maximal operator starts to resemble a traditional Hardy-Littlewood maximal operator in two dimensions.

The exceptional situation a) of this paper is necessitated by our use of damping functions in conjunction with the theorem of Sogge and Stein [SoSt] that we will describe below. It is unlikely that it can be avoided without using substantial additional ideas. A canonical example of when the first kind of exception situation occurs is when $F(x, y) = p(ax + by)$, in other words, when $F(x, y)$ is effectively a function of one variable.

Exceptional situation b) may be viewed as rare in the sense that it requires a certain polynomial to have a zero of high order among other things; however, a simple example of where it happens is the function $(y + x^a)^b + x^c$ for $c \geq b > 2$, $a > 1$, and $ab < c$. Although we will not prove it here, further manipulations involving Newton polygons can be used to show that if the exceptional condition holds for one of $\partial_y(F_e(1,y))$ and $\partial_y(F_e(-1,y))$, it holds for the other. Hence in the statement of Theorem 1.1 we could have just used one of the two functions.

The exceptional situation b) arises for the following reason. The main theorem of [G1] gives Theorem 1.1 if $f(x,y)$ is already in adapted coordinates. Much of the analysis of [G1] carries over even in nonadapted coordinates; the y-variable shift by $\psi(x)$ does not interfere with most of the argument. The exception occurs when one cannot avoid using integrations by parts in the x-variable in arguments resembling the proof of the Van der Corput lemma. This happens when the bisectrix intersects $N(F)$ in the interior of a bounded edge satisfying the above conditions on $F_e(1,y)$ and its y derivatives. If the order b of the zero of $\psi(x)$ at the origin is at least $\frac{1}{|m_e|}$, then $\psi(x)$ is too small to cause any serious problems in such integrations by parts. If b is less than $\frac{1}{|m_e|}$ then it seems difficult to adapt the arguments of [G1] to these situations. So as long as we can find some linear T such that $f \circ T$ avoids such situations, then the methods of [G1] can be adapted to the current situation and Theorem 1.1 can be proved.

The stipulation that $b < \frac{1}{|m_e|}$ is more than a technical improvement in the statement of Theorem 1.1. For example, if $f(x,y)$ is a mixed homogeneous function, then it is not hard to show that in converting to adapted coordinates $\psi(x)$ can always be taken to be of the form cx^b, and if the bisectrix of $F(x,y) = f(x, y + cx^b)$ intersects $N(F)$ in the interior of a compact edge of slope m_e then $b = \frac{1}{|m_e|}$. Hence Theorem 1.1 covers the mixed homogeneous case (as long as $D(x,y)$ does not have a zero of infinite order at the origin) because we include the condition. For $h(q_0) \geq 2$ this result was first proved in [IkKeMu1].

Strategy of the Proof.

Ever since [SoSt] one successful method of proving L^p boundedness theorems for maximal operators such as (1.1) has involved embedding M in an analytic family M_z. The idea is that one replaces the standard surface measure $d\sigma(q)$ with the damped surface measure $e^{z^2}|h|^z d\sigma(q)$, and then defines the maximal operator M_z to be the analogue of (1.1) with $d\sigma(q)$ replaced by $e^{z^2}|h|^z d\sigma(q)$. One shows that for some $s_0 < 0$, M_z is bounded on L^∞ whenever $Re(z) > s_0$, uniformly in $|Im(z)|$ for fixed $Re(z)$, and that for some $s_1 > 0$, M_z is bounded on L^2 whenever $Re(z) > s_1$, again uniformly in $|Im(z)|$ for fixed $Re(z)$. The e^{z^2} factor is present to ensure uniform L^2 bounds on each vertical line. Then by a well known interpolation technique for maximal operators (see Ch 11 of [St2] for details), one obtains a $p_0 > 2$ such that $M = M_0$ is bounded on L^p for $p > p_0$. The hope is that the damping function $h(z)$ can be chosen so that p_0 is optimal.

In the above interpolation, the L^∞ bounds are typically obtained using the observation that $||M_z||_{L^\infty \to L^\infty}$ is bounded by the L^1 norm of $C|h|^{Re(z)}$. Thus as long as $|h|^{s_0} d\sigma(q)$ is a finite measure, one obtains the desired uniform L^∞ bounds on any vertical line $Re(z) = s$ for $s > s_0$. For the L^2 boundedness, we will use the following consequence of a theorem of Sogge and Stein:

Theorem 1.2. [SoSt]. Suppose the surface measure $\phi(q)d\sigma(q)$ is as in (1.1) and there are $C, \epsilon > 0$ such that for all multiindices α with $|\alpha| = 0, 1$ the Fourier transform of the measure $d\sigma_z(q) = e^{z^2}|h(x)|^z\phi(q)d\sigma(q)$ satisfies

$$|\partial^\alpha \hat{\sigma_z}(\lambda)| < C(1 + |\lambda|)^{-\frac{1}{2}-\epsilon} \tag{1.2}$$

Then there is a constant C' depending on C and ϵ such that $||M_z f||_2 \leq C'||f||_2$ for all $f \in L^2$.

In practice, if one has (1.2) for $\alpha = 0$, it will generally automatically hold for all the first derivatives since the effect of taking such a derivative is to replace the cutoff function by another one. In the proof of Theorem 1.1 in this paper, we will work in this framework. We will first write the surface Q near q_0 as the union of finitely many "slivers" containing q_0 on their boundaries. These slivers will be defined using the Newton polygon of the function $f(x, y)$ defined above Theorem 1.1 when put in generic adapted coordinates; we will effectively be doing a coarse resolution of singularities using Lemmas 2.2 and 2.3 of the next section. We will then define the damping function separately on each sliver. The damping functions will be analogues of those of [G1]. There will be four types of slivers and showing (1.2) holds will be done by separately estimating the contribution to $\hat{\sigma_z}(\lambda)$ coming from each type of sliver, again using methods analogous to those of [G1]. It should be pointed out that the idea of dividing a neighborhood of a point into slivers with respect to a Newton polygon for the purpose of proving oscillatory integral estimates such as (1.2) is quite old; it appears in [PSt] and its predecessors for example and was also used in analyzing such maximal operators in [IkKeMu1] and [IkKeMu2].

2. Lemmas about Newton polygons; subdivisions into slivers

In the analysis of this paper, for the appropriate $F(x, y)$ one treats on similar footing all compact edges of $N(F)$ intersecting the set $\{(t_1, t_2) : t_2 > t_1\}$. To avoid exceptional situations such as those of part b) in the statement of Theorem 1.1 for any such edge not intersecting the bisectrix in its interior, we have the following lemma.

Lemma 2.1. Suppose $R(x, y)$ is a smooth function on a neighborhood of the origin that is in generic adapted coordinates. Let d^* denote $\max(2, d(R))$. Suppose e is any compact edge of $N(R)$ lying entirely on or above the bisectrix. Then if there is $y_0 \neq 0$ such that if $R_e(1, y_0) \neq 0$ and $\partial_y R_e(1, y)$ has a zero at y_0 of order greater than $d^* - 1$, then e has slope -1 and upper vertex lying on the y axis.

Proof. Since $R(x, y)$ is in generic adapted coordinates, the uppermost vertex of $N(R)$ is $(0, a)$ for some $a \geq 2$. The point $(d(R), d(R))$ is on $N(R)$, and the line of slope -1 containing this point intersects $N(R)$ at $(0, 2d(R))$. Since all edges of $N(R)$ have slope at least -1 and $(d(R), d(R)) \in N(R)$, we conclude that $a \leq 2d(R)$ with equality holding iff there is an edge of $N(R)$ connecting $(d(R), d(R))$ to $(0, 2d(R))$. Hence if e is a compact edge of $N(R)$ lying entirely on or above the bisectrix, either e is the segment $(d(R), d(R))$ to $(0, 2d(R))$ or $a < 2d(R)$. Thus if e satisfies the assumptions of this lemma and we can show that $R_e(1, y)$ has degree at least $2d(R)$, then in particular e has slope -1 and upper vertex lying on the y axis as needed.

So assume e is a compact edge lying entirely on or above the bisectrix such that there is $y_0 \neq 0$ with $R_e(1, y_0) \neq 0$ and such that $\partial_y R_e(1, y)$ has a zero at y_0 of order greater than d^*. First consider the case where e's lower vertex is $(d(R), d(R))$. In particular, $d(R)$ is an integer. We will show that $R_e(1, y)$ has degree at least $2d(R)$. Note that $R_e(1, y)$ has a zero at $y = 0$ of order $d(R)$ and therefore $\partial_y R_e(1, y)$ has a zero at $y = 0$ of order $d(R) - 1$. Since we are assuming it also has a zero at y_0 of order greater than $d(R) - 1$, $\partial_y R_e(1, y)$ must have degree least $2d(R) - 1$. Hence $R_e(1, y)$ has degree at least $2d(R)$. This means that the upper vertex of e is at least $2d(R)$ and by the previous paragraph we are done.

Next, we consider the case where e does not contain $(d(R), d(R))$. Hence e lies entirely above the bisectrix, and $R_e(1, y)$ can be written as $y^{d'} p(y)$ where $d' > d(R)$. This means that $\partial_y R_e(1, y)$ is of the form $y^{d'-1} q(y)$. Since $\partial_y R_e(1, y)$ is assumed to have a zero $y_0 \neq 0$ of order greater than $d(R) - 1$, we can write

$$\partial_y R_e(1, y) = y^{d'-1} (y - y_0)^{d''-1} r(y) \tag{2.1}$$

Here $d'' > d(R)$. Note that (2.1) implies that the degree of $\partial_y R_e(1, y)$ is at least $d' + d'' - 2 > 2d(R) - 2$, so the degree of $R_e(1, y)$ is greater than $2d(R) - 1$. Hence the upper vertex of e must be the upper vertex of $N(R)$; otherwise $N(R)$ would have a vertex at height greater than $2d(R)$ and since $a \leq 2d(R)$ this can't happen. We conclude the upper vertex of e is given by $(0, a)$ for some $2d(R) - 1 < a \leq 2d(R)$. As a result, the degree of $\partial_y R_e(1, y)$ is greater than $2d(R) - 2$ and at most $2d(R) - 1$.

If r had positive degree, then the degree of $\partial_y R_e(1, y)$ would be at least $d' + d'' - 1 > 2d(R) - 1$, contradicting the above. So $r(y)$ is constant and we may write

$$\partial_y R_e(1, y) = c y^{d'-1} (y - y_0)^{d''-1} \tag{2.2}$$

Similarly, if d' or d'' were equal to $d(R) + 1$ or greater then the degree of $\partial_y R_e(1, y)$ would be greater than $2d(R) - 1$, again giving a contradiction. So we have $d(R) < d', d'' < d(R) + 1$. Since d' and d'' are both integers, this means $d' = d''$. "Homogenizing" (2.2), for some k we get that

$$\partial_y R_e(x, y) = c y^{d'-1} (y - y_0 x^k)^{d'-1} \tag{2.3}$$

Looking at the term of (2.3) whose degree in y is second-highest we see that k is an integer. Hence e is an edge of slope $-\frac{1}{k}$ whose upper vertex is $(0, 2d' - 1)$ and our objective is to show that $k = 1$. Assume $k \geq 2$; we will arrive at a contradiction. Since we are dealing with the case that e lies entirely above the bisectrix, the point $(d(R), d(R))$ lies above the line containing e. Since this line intersects the bisectrix at $(\frac{k}{k+1}(2d'-1), \frac{k}{k+1}(2d'-1))$, we conclude that

$$\frac{k}{k+1}(2d'-1) < d(R) \tag{2.4}$$

Since $d' > d(R)$ and $\frac{k}{k+1} \geq \frac{2}{3}$, this in turn implies that

$$\frac{2}{3}(2d(R)-1) < d(R) \tag{2.5}$$

Equivalently, $d(R) < 2$. Since $d' < d(R)+1$ and $d'-1$ is an integer at least one, $d'-1 = 1$ and (2.3) just becomes

$$\partial_y R_e(x,y) = cy(y - y_0 x^k) \tag{2.6}$$

Now note that $\partial_y R_e(1,y)$ no longer has a zero of order greater than one, so it no longer falls under the assumptions of this lemma. So this situation cannot happen; we have arrived at a contradiction and we are done.

Lemma 2.2. Suppose $R(x,y)$ is a smooth function on a neighborhood of the origin such that $R(0,0) = 0$. Suppose v is a vertex of $N(R)$ that is the intersection of compact edges e_1 and e_2 with slopes $0 > m_1 > m_2$. Let $M_1 = -\frac{1}{m_1}$ and $M_2 = -\frac{1}{m_2}$. Let $R_{cd}x^c y^d$ denote the term of the Taylor expansion of $R(x,y)$ at the origin corresponding to v. Then on a sufficiently small neighborhood of the origin, there is an $N > 0$ such that if $N|x|^{M_1} < |y| < \frac{1}{N}|x|^{M_2}$, then we have

$$\frac{1}{2}|R_{cd}x^c y^d| < |R(x,y)| < 2|R_{cd}x^c y^d| \tag{2.7}$$

If v lies on the y axis and is the upper vertex of a compact edge e_1 with slope $-\frac{1}{M_1}$, then there is an $N > 0$ such that if $N|x|^{M_1} < |y|$ then once again (2.7) holds.

Proof. Without loss of generality, we may restrict our attention to (x,y) in the upper right quadrant. Write the Taylor expansion of $R(x,y)$ at the origin as $\sum_{a,b} R_{ab}x^a y^b$. We first prove (2.7) in the case where v is the intersection of two compact edges, whose equations we denote by $t_1 + M_1 t_2 = \alpha_1$ and $t_1 + M_2 t_2 = \alpha_2$. For a large K we can write

$$R(x,y) - R_{cd}x^c y^d = \sum_{(a,b): c \leq a < M,\ d \leq b < M,\ (a,b) \neq (c,d)} R_{ab}x^a y^b$$

$$+ \sum_{(a,b): a < c,\ d < b < M,\ a + M_2 b \geq \alpha_2} R_{ab}x^a y^b + \sum_{(a,b): c < a < M,\ b < d,\ a + M_1 b \geq \alpha_1} R_{ab}x^a y^b + E_K(x,y) \tag{2.8}$$

Here $E_K(x,y)$ satisfies

$$|E_K(x,y)| < C(|x|^K + |y|^K) \tag{2.9}$$

We start by noting that the first sum in (2.8) is less than $\frac{1}{8}|r_{cd}|x^c y^d$ in absolute value if (x,y) is in a sufficiently small neighborhood of the origin, which we may assume. As for the second sum, if one changes coordinates from (x,y) to (x,y'), where $y' = x^{M_2}y$, then $(x,y') \in [0,1] \times [0,\frac{1}{N}]$ whenever $y < \frac{1}{N}x^{M_2}$. Observe that under this coordinate change, a given term $R_{ab}x^a y^b$ of the second sum becomes $R_{ab}x^{a+M_2 b}(y')^b$. Since $a + M_2 b \geq \alpha_2$ and $b > d$ in each term in the second sum, the entire sum can be written as $x^{\alpha_2}(y')^d(y'f(x,y'))$ for some $f(x,y')$ which is a polynomial in y' and a fractional power of x. Thus if N is sufficiently large, whenever $y' < \frac{1}{N}$ the sum is of absolute value less than $\frac{1}{8}|R_{cd}|x^\alpha(y')^d = \frac{1}{8}|R_{cd}|x^c y^d$. Since $y' < \frac{1}{N}$ is equivalent to $y < \frac{1}{N}x^{M_2}$, these are the bounds we need.

The third sum is dealt with in exactly the same way, reversing the roles of the x and y axes and the edges e_1 and e_2. Lastly, since $\frac{1}{N}x^{M_2} > y > Nx^{M_1}$ the error term $E_K(x,y)$ is less than $\frac{1}{8}|r_{cd}|x^c y^d$ in absolute value for small $|x|, |y|$ if K is chosen sufficiently large. Putting these all together, we get that $|R(x,y) - R_{cd}x^c y^d| < \frac{1}{2}|R_{cd}|x^c y^d$ as needed. This completes the proof of Lemma 2.2 for the case where (c,d) is the intersection of two compact edges of $N(R)$.

We now move to the case where (c,d) is on the y-axis and is the upper vertex of a compact edge e_1 of $N(R)$. We again examine the sum (2.8). In the case at hand, since (c,d) is on the y-axis, the second sum in (2.8) is empty. The second sum is where the condition $y < \frac{1}{N}x^{M_2}$ was used above, and the third sum is where the condition $y > Nx^{M_1}$ was used. Since we have no second sum, the lack of a condition $y < \frac{1}{N}x^{M_2}$ holding does not cause any problem in repeating the above argument. For the third sum we use the condition $y > Nx^{M_1}$ exactly as before, and for the first and fourth sum the previous argument works unmodified. Hence (2.7) holds again and we are done.

Lemma 2.3. Suppose $R(x,y)$ is a smooth function of y and a fractional power of x on a neighborhood of the origin such that $R(0,0) = 0$. Write the Taylor expansion of $R(x,y)$ at the origin as $\sum_{a,b} R_{ab}x^a y^b$. For a given $M > 0$ let $R_M(x,y)$ denote the sum of the nonzero terms of this Taylor expansion for which $a + Mb$ is minimized; in particular $R_M(x,y)$ is either of the form $R_e(x,y)$ for a compact edge e of $N(R)$ or is equal to $R_{cd}x^c y^d$ for a vertex (c,d) of $N(R)$. Denote this minimal value of $a + Mb$ by α. Then for any $r \in \mathbf{R}$ and any $\epsilon > 0$, there is a $\delta > 0$ such that on the set $\{(x,y) \in \mathbf{R}^2 : 0 < x < \delta, (r-\delta)x^M < y < (r+\delta)x^M\}$ we have

$$|R(x,y) - R_M(x,y)| < \epsilon x^\alpha \tag{2.10}$$

Proof. On the region $\{(x,y) \in \mathbf{R}^2 : 0 < x < \delta, (r-\delta)x^M < y < (r+\delta)x^M\}$, we do the coordinate change $(x,y) = (x, x^M y')$, converting the region into the box $(0,\delta)\times(r-\delta, r+\delta)$. In the new coordinates, the finite Taylor expansion $R(x,y) = \sum_{a,b<K} R_{ab}x^a y^b + O(|x|^K + |y|^K)$ becomes of the form

$$R(x, x^M y') = x^\alpha R_M(1,y') + x^{\alpha+\zeta}s(x,y') + O(|x|^K + |x|^{KM}|y'|^K) \tag{2.11}$$

Here $\zeta > 0$ and $f(x,y')$ is a polynomial in y' and a fractional power of x. For any $\epsilon' > 0$, if δ is sufficiently small we have $|R_M(1,y') - R_M(1,r)| < \epsilon'$ for all $|y' - r| < \delta$. Equivalently, $|R_M(1,y') - R_M(1,r)|x^\alpha < \epsilon' x^\alpha$. Furthermore, if δ is sufficiently small $x^{\alpha+\zeta}|s(x,y')|$ and the $O(|x|^K + |x|^{KM}|y'|^K)$ term are less than $\epsilon' x^\alpha$ whenever x and $|y' - r|$ are sufficiently smalll. Combining, if δ is sufficiently small then on our domain we have

$$|R(x, x^M y') - x^\alpha R_M(1, r)| < 3\epsilon' x^\alpha R_M(1, r) \tag{2.12}$$

Translating this back into the original coordinates, we have

$$|R(x, y) - R_M(x, y)| < 3\epsilon' x^\alpha R_M(1, r) \tag{2.13}$$

Taking $\epsilon = 3\epsilon' R_M(1, r)$ gives us the lemma and we are done.

We now are in a position to set up the proof of the main theorem, Theorem 1.1. Recall we have a surface Q with a distiguished point (x_0, y_0, z_0) that is the graph of some smooth function $g(x, y)$ defined near (x_0, y_0) such that $\nabla g(0, 0) = (0, 0)$. Suppose the assumptions of Theorem 1.1 hold. Then, after a linear coordinate change if necessary, we may assume $f(x, y) = g(x_0 + x, y_0 + y) - z_0$ has a generic adapted coordinate system on a neighborhood of the origin for which the exceptional situation b) of Theorem 1.1 does not occur. Therefore there is a smooth $\psi(x)$ with $\psi(0) = 0$ such that $F(x, y) = f(x, y + \psi(x))$ is in generic adapted coordinates, and if $N(F)$ intersects the bisectrix in the interior of a compact edge e then e does not satisfy the exceptional situation b) of the statement of Theorem 1.1.

Definition of slivers for $F(x, y)$

We now use Lemmas 2.2 and 2.3 on $F(x, y)$ and its various y derivatives to subdivide a small neighborhood B of $(0, 0)$ into "slivers" containing the origin. The case where $N(F)$ has exactly one vertex (which is therefore on the y-axis) is easier and will be treated separately, so in the following we always assume $N(F)$ contains multiple vertices. Denote the vertices of $N(F)$ above the bisectrix by $v_1, ..., v_k$ where if $i < j$ then v_i is below v_j. Let e_i denote the edge of $N(F)$ whose upper vertex is v_i; if v_1 is the lowest vertex of $N(F)$ then we just do not define e_1. Let $m_i = -\frac{1}{M_i}$ denote the slope of e_i. Write $v_i = (a_i, b_i)$; observe $b_i \geq 2$ for all i since (a_i, b_i) lies above the bisectrix and we are assuming $F(0, 0) = 0$ and $\nabla F(0, 0) = (0, 0)$.

For v_i that is the intersection of two compact edges of $N(F)$, define D_i to be the set $\{(x, y) \in B : N_0|x|^{M_i} < |y| < \frac{1}{N_0}|x|^{M_{i+1}}\}$. Here N_0 is large enough so that we may invoke Lemma 2.2 and say there are $c_0, c_1 > 0$ such that (assuming B is small enough) for $m = 0, 1, 2$ on D_i we have

$$c_1|x|^{a_i}|y|^{b_i - m} > |\frac{\partial^m F}{\partial y^m}(x, y)| > c_0|x|^{a_i}|y|^{b_i - m} \tag{2.14}$$

It should be pointed out that if $b_i = m = 2$, then (2.14) holds by applying Lemma 2.3, reversing the roles of the x and y axes and setting $r = 0$. If v_i is the upper vertex of $N(F)$, we define D_i to be the points where $N_0|x|_i^M < |y|$, in which case (2.14) still holds on D_i by Lemma 2.2.

We next subdivide the set $B-\cup_i D_i$ into some slivers touching the origin amenable to the analysis of this paper. We only describe the slivers for $x>0$; the ones where $x<0$ are defined analogously. Note that the points of $B-\cup_i D_i$ where $x>0$ can be written as $\cup_i C_i$, where

$$C_1 = \{(x,y) \in B : x > 0,\ |y| < N_0 x^{M_1}\} \tag{2.15a}$$

$$C_i = \{(x,y) \in B : x > 0,\ \frac{1}{N_0}x^{M_i} < |y| < N_0 x^{M_i}\} \qquad (i > 1) \tag{2.15b}$$

(In the event that v_1 is the lowest vertex of $N(F)$, one takes M_2 in (2.15a) and then (2.15b) is valid for $i>2$). Suppose r is such that $F_{e_i}(1,r) \neq 0$, but $\partial_y F_{e_i}(1,r)$ has a zero of order greater than $d^* - 1 = \max(2, d(F)) - 1$ at r. Then if e_i intersects the bisectrix in its interior, by assumption the exceptional case of Theorem 1.1 part b does not occur. If e_i does not intersect the bisectrix in its interior, then by Lemma 2.1 e_i has slope -1 and intersects the y-axis. In either event, if $i=1$ and $|r| < N_0$ or $i>1$ and $\frac{1}{N_0} < |r| < N_0$, we define E_{ir} to be the sliver

$$E_{ir} = \{(x,y) \in B : x > 0,\ (r-\delta_r)x^{M_i} < y < (r+\delta_r)x^{M_i}\} \tag{2.16}$$

Here δ_r is a small constant to be determined by our future arguments. We will refer to the (finitely many) E_{ir} occurring as E_{ij} in the rest of this paper.

For any r other than these, we may let $1 \leq k \leq d^*$ and $\delta_r, C_r > 0$ be such that on $[r-\delta_r, r+\delta_r]$ we have $|\partial_y^k F_{e_i}(1,y)| > C_r$. If $F_{e_i}(1,r) \neq 0$ this follows from the above definition of the D_i and if $F_{e_i}(1,r) = 0$ it follows from the fact that any zero of any $F_{e_i}(1,y)$ has order at most $d(F)$ in adapted coordinates. As a result, on the set $B_r = \{(x,y) \in B : x > 0,\ (r-\delta_r)x^{M_i} \leq y \leq (r+\delta_r)x^{M_i}\}$, given that e_i contains (a_i, b_i) we have

$$|\partial_y^k F_{e_i}(x,y)| > C_r x^{a_i + M_i(b_i - k)} \tag{2.17}$$

By applying Lemma 2.3 to $\partial_y^k F(x,y)$ we can assume B is small enough that we also have

$$|\partial_y^k F(x,y) - \partial_y^k F_{e_i}(x,y)| < \frac{C_r}{2} x^{a_i + M_i(b_i - k)} \tag{2.18}$$

Putting (2.17) and (2.18) together, on B_r we have

$$|\partial_y^k F(x,y)| > \frac{C_r}{2} x^{a_i + M_i(b_i - k)} \tag{2.19}$$

By compactness, we can write $B - \cup_i D_i - \cup_{ij} E_{ij}$ as the union of finitely many slivers on which (2.19) is satisfied. For a given edge e_i, we write the slivers for which $k=1$ as F_{ij} and the slivers for which $k>1$ by G_{ij}. We denote the value of k corresponding to a given G_{ij} by k_{ij}. Note that each $k_{ij} \leq d^*$.

The above decompositions were for the case where $N(F)$ had more than one vertex. When $N(F)$ has exactly one vertex, since F it is in generic adapted coordinates it is of the form $(0,k)$. In this case, we simply designate a neighborhood of the origin as a single G_{ij}, with $k_{ij}=k$. In general, the arguments for this G_{ij} will be simplified versions of the G_{ij} arguments for the multivertex case.

Let $v(F)$ denote the set of vertices of $N(F)$, and define $F^*(x,y)$ by

$$F^*(x,y) = \Big(\sum_{(v_1,v_2)\in v(F)} (x^{v_1}y^{v_2})^2\Big)^{\frac{1}{2}}$$

The function $F^*(x,y)$ will be used in defining the damping function. To this end, first note that $N((F^*)^2)$ is the double of $N(F)$ and therefore $d((F^*)^2) = 2d(F)$. As a result, by [V], $F^*(x,y)^{-t}$ is integrable on a neighborhood of the origin iff $t < \frac{1}{2}d((F^*)^2) = d(F)$. We apply Lemmas 2.2 and 2.3 to $F^*(x,y)^2$ in place of $F(x,y)$ and obtain that if the N_0 in the definition of D_i were chosen sufficiently large, then there is a constant $C_0 > 0$ such that on each D_i we have

$$\frac{1}{C_0}|x^{a_i}y^{b_i}| < F^*(x,y) < C_0|x^{a_i}y^{b_i}| \tag{2.20}$$

Similarly, by Lemma 2.3 on each E_{ij}, F_{ij} and G_{ij}, the constant C_0 can be taken so that we have

$$\frac{1}{C_0}|x|^{a_i+M_ib_i} < F^*(x,y) < C_0|x|^{a_i+M_ib_i} \tag{2.21}$$

We now subdivide the surface Q near our distinguished point q_0 in accordance with the above subdivisions, applied to the function $f(x,y)$ in generic adapted coordinates. In other words, we let $\psi(x)$ be such that $F(x,y) = f(x, y+\psi(x))$ is in generic adapted coordinates such that the exceptional cases of Theorem 1.1 do not hold, and define D_i, E_{ij}, F_{ij}, and G_{ij} to be the above slivers as defined for $F(x,y)$. We next transfer these slivers into the original coordinates of the surface Q; let D_i' be the portion of Q above the set $\{(x,y) : (x-x_0, y-y_0-\psi(x-x_0)) \in D_i\}$, with the analogous definitions for E_{ij}', F_{ij}', and G_{ij}'. We also will have use for D_i in the original nonadapted coordinates, centered at (x_0,y_0). To that end we let $D_i'' = \{(x,y) : (x, y-\psi(x)) \in D_i\}$, making the analogous definitions for E_{ij}'', F_{ij}'', and G_{ij}''.

The next lemma will be useful in bounding the contribution of our integrals over E_{ij}'' in the L^2 estimates of section 4.

Lemma 2.4. There is a constant C such that on E_{ij}'' we have $|\partial_{xx}f(x,y)| \geq C|x|^{a_i+b_iM_i-2}$.

Proof. We consider slivers for which $x>0$ as the $x<0$ slivers are entirely analogous. Recall each E_{ij} is a region of the form $\{(x,y) : 0 < x < \mu, (r-\nu)x^{M_i} < y < (r+\nu)x^{M_i}\}$, where $F_{e_i}(1,r) \neq 0$, but where $\partial_y^l F_{e_i}(1,r) = 0$ for at least $l = 1,2$. Under the map $(x,y) \rightarrow (x, y - rx^{M_i})$, the set E_{ij} becomes the region $E_{ij}''' = \{(x,y) : 0 < x < \mu, |y| < \nu x^{M_i}\}$, and if $G(x,y)$ denotes $F(x, y+rx^{M_i})$ then $G_{e_i}(1,0) \neq 0$ but $\partial_y^l G_{e_i}(1,0) = 0$ for $l = 1,2$.

In terms of Newton polygons, the above can be translated as follows. Since e_i is an edge of $N(F)$ with equation $x + M_i y = a_i + M_i b_i$, $N(G)$ has an edge with the same equation which goes all the way to the x-axis since $G_{e_i}(1,0) \neq 0$. As a result, for $l = 1,2$ $N(\partial_x^l G)$ has an edge with equation $x + M_i y = a_i + M_i b_i - l$ which extends to the x axis. On the other hand, since $\partial_y G_{e_i}(1,y)$ has a zero of order at least two at 0, for $l = 1,2$ $N(\partial_y^l G)$ intersects the line with equation $x + M_i y = a_i + (M_i - l)b_i$ but does not contain $(a_i + (M_i - l)b_i, 0)$. Using this fact for $l = 1$ and taking an x derivative shows that $N(\partial_{xy} G)$ intersects the line with equation $x + M_i y = a_i - 1 + (M_i - 1)b_i$ but does not contain $(a_i - 1 + (M_i - 1)b_i, 0)$. Using (2.11) in conjunction with these latter observations concerning the Newton polygons, we obtain that for any $\eta > 0$, if ν were chosen sufficiently small, then on E_{ij}''' for $l = 1,2$ we have

$$|\partial_y^l G(x,y)| < \eta |x|^{a_i + (M_i - l)b_i} \tag{2.22a}$$

$$|\partial_{xy} G(x,y)| < \eta |x|^{a_i - 1 + (M_i - 1)b_i} \tag{2.22b}$$

On the other hand Lemma 2.3 in conjuction with the above conditions on the Newton polygons of the x-derivatives ensures that for some $c_1 > 0$, on E_{ij}''' for $l = 1,2$ we have

$$|\partial_x^l G(x,y)| \geq c_1 |x|^{a_i + M_i b_i - l} \tag{2.22c}$$

Next, as mentioned after (2.15), E_{ij} is only defined in two situations. The first is when e_i intersects the bisectrix in its interior and the function $\psi(x)$ such that $F(x,y) = f(x, y + \psi(x))$ has a zero of order at least M_i at $x = 0$. In the second situation, e_i does not intersect the bisectrix in its interior, but by Lemma 2.1 $M_i = 1$ and thus $\psi(x)$ still has a zero of order at least M_i at $x = 0$. In either case, since $f(x,y) = F(x, y - \psi(x)) = G(x, y - \psi(x) + rx^{M_i})$, we can write $f(x,y) = G(x, y + \xi(x))$ where $\xi(x)$ has a zero of order at least M_i at $x = 0$. Applying the chain rule, we get that on E_{ij}'' we have

$$\partial_{xx} f(x,y) = \partial_{xx} G(x, y + \xi(x)) + 2\xi'(x)\partial_{xy} G(x, y + \xi(x))$$

$$+\xi''(x)\partial_y G(x, y + \xi(x)) + (\xi'(x))^2 \partial_{yy} G(x, y + \xi(x))$$

Equation (2.22c) ensures that $|\partial_{xx} G(x, y+\xi(x))| \geq c_1 |x|^{a_i + M_i b_i - 2}$, and equations (2.22$a$)–(2.22$b$) coupled with the fact that $\xi(x)$ has a zero of order at least M_i at $x = 0$ ensure that the remaining terms can be made less than any $\eta' |x|^{a_i + M_i b_i - 2}$. We conclude that $|\partial_{xx} f(x,y)|$ is at least $\frac{c_1}{2} |x|^{a_i + M_i b_i - 2}$ on E_{ij}'' if ν were chosen appropriately small. This completes the proof of the lemma.

Definition of the damping factor.

We now define the damping factor on the surface Q in a neighborhood of $q_0 = (x_0, y_0, g(x_0, y_0))$. Above a point (x,y) it will be of the form $e^{z^2} |h(x,y)|^z |D(x,y)|^{\delta z}$, where $D(x,y)$ is the Hessian determiant of g at (x,y) and δ is a small positive number to be determined by our arguments. The function $h(x,y)$ will be defined in the form $H(x - x_0, y - y_0 - \psi(x - x_0))$, where $H(x,y)$ is expressed in terms of $F(x,y) = f(x, y + \psi(x))$.

To this end, once again let $d^* = \max(2, d(F))$. Since $F(x,y)$ is in adapted coordinates, $d(F) = h(q_0)$ and thus equivalently we have $d^* = \max(2, h(q_0))$.

On each D_i, E_{ij}, F_{ij}, as well as on each G_{ij} with $k_{ij} = 2$, we define $H(x,y)$ to be $F^*(x,y)^{\frac{1}{2}-\frac{1}{d^*}}$. On the remaining G_{ij}, if $N(F)$ has just one vertex we let $H(x,y) = |\frac{\partial^2 F}{\partial y^2}(x,y)|^{\frac{1}{2}}$, while if $N(F)$ has multiple vertices let $H(x,y) = |x|^{M_i - \frac{a_i + M_i b_i}{d^*}} |\frac{\partial^2 F}{\partial y^2}(x,y)|^{\frac{1}{2}}$. One thing worth mentioning concerning these latter $H(x,y)$ is the following. Since $|y| < C|x|^{M_i}$ on G_{ij}, by (2.11) with $\alpha = a_i + M_i b_i$, on G_{ij} one has $|\frac{\partial^2 F}{\partial y^2}(x,y)| < C|x|^{a_i + M_i(b_i - 2)}$. As a result, $|H(x,y)| \leq C|x|^{(a_i + M_i b_i)(\frac{1}{2} - \frac{1}{d^*})}$ and in view of (2.21) this gives

$$|H(x,y)| \leq CF^*(x,y)^{\frac{1}{2}-\frac{1}{d^*}} \tag{2.23}$$

Note the right-hand side of (2.23) is exactly $CH(x,y)$ for the other regions.

3. L^∞ estimates

Define the operator M_z to be the maximal operator (1.1) with respect to the measure $e^{z^2}|h(x,y)|^z|D(x,y)|^{\delta z}\phi(q)d\sigma(q)$ in place of $\phi(q)d\sigma(q)$. Note that for $z = 0$, M_z is exactly M. These will be the analytic family of maximal operators used in proving Theorem 1.1 as described at the end of section 1. We first prove the L^∞ to L^∞ boundedness properties of the M_z we need.

Theorem 3.1. Write $d = d(F)$. If $d > 2$, then for any $s > -\frac{2}{d-2}$, if δ is sufficiently small (depending on s) then there exists a constant C such that $||M_z||_{L^\infty \to L^\infty} < C$ for all z with $Re(z) = s$. If $d \leq 2$, the same holds for any $s \in \mathbf{R}$.

Proof. We will use the fact that $||M_z||_{L^\infty \to L^\infty}$ is bounded by the L^1 norm of the damping function of M_z. We consider the case $d = d(F) \leq 2$ first. Note that for each D_i, E_{ij} or F_{ij} the function $h(x,y)$ is given by $F^*(x - x_0, y - y_0 - \psi(x - x_0))^{\frac{1}{2}-\frac{1}{d^*}} = 1$. Since in adapted coordinates each k_{ij} is at most $d^* \leq 2$ here, for each G_{ij} that may appear $h(x,y)$ is also just 1. We conclude that the damping factor always equal to $e^{z^2}|D(x,y)|^{-\delta z}$. On a given vertical line $Re(z) = s$, this has magnitude bounded by $C_s|D(x,y)|^{-\delta s}$. Since $D(x,y)$ is assumed to be of finite type at the origin, there is some $\epsilon > 0$ for which $|D(x,y)|^{-\epsilon}$ is integrable on a neighborhood of the origin. Hence as long as $\delta < \frac{\epsilon}{s}$, the damping factor is integrable, with integral uniformly bounded on $Re(z) = s$. This is exactly what we needed to prove.

Now suppose $d > 2$. On any D_i', E_{ij}', F_{ij}', or G_{ij}' with $k_{ij} = 2$, $h(x,y) = F^*(x - x_0, y - y_0 - \psi(x - x_0))$ and thus the damping factor has magnitude $F^*(x - x_0, y - y_0 - \psi(x - x_0))^{(\frac{1}{2}-\frac{1}{d})Re(z)}|D(x,y)|^{\delta Re(z)}$. As mentioned at the end of section 2, a result of Varchenko says that $F^*(x,y)^{-t}$ is integrable on a neighborhood of the origin iff $t < \frac{1}{d}$. Thus the same is true for $F^*(x - x_0, y - y_0 - \psi(x - x_0))^{-t}$. As a result, $F^*(x - x_0, y - y_0 - \psi(x - x_0))^{(\frac{1}{2}-\frac{1}{d})Re(z)}$ is integrable on a neighborhood of the origin iff $Re(z) > -\frac{2}{d-2}$. Consequently, by Holder's inequality, for fixed $s > -\frac{2}{d-2}$, by choosing δ sufficiently small we have that when $Re(z) = s$ the damping factor is integrable over any D_i', E_{ij}', F_{ij}', or G_{ij}' with $k_{ij} = 2$, with integral uniformly bounded in $Im(z)$. This is what we need here.

We now move on to the G_{ij}' with $k_{ij} > 2$. If $N(F)$ has one vertex, we use the fact that $\frac{\partial^2 F}{\partial y^2}(x,y)$ has nonvanishing $(k_{ij}-2)$th derivative in the y direction, where $k_{ij} \leq d$. Thus if $Re(z) = s > -\frac{2}{d-2}$, then $|H(x,y)|^z = |\frac{\partial^2 F}{\partial y^2}(x,y)|^{\frac{z}{2}}$ is integrable in y with integral uniformly bounded in $Im(z)$. Making δ sufficiently small and using Holder's inequality again gives the desired result. Suppose now $N(F)$ has multiple vertices. We will consider those G_{ij} for which $x > 0$ as the G_{ij} for which $x < 0$ are done in the same way. In the G_{ij} coordinates we can write the damping function as

$$|\tilde{D}(x,y)|^{\delta z} x^{M_i - \frac{a_i + M_i b_i}{d}} |\frac{\partial^2 F}{\partial y^2}(x,y)|^{\frac{1}{2}}$$

Recall that G_{ij} is of the form $\{(x,y) : 0 < x < \eta, (r-\delta)x^{M_i} < y < (r+\delta)x^{M_i}\}$. Analogous to (2.11), on the box $(0,\eta) \times (r-\delta, r+\delta)$ for any K one can write

$$F(x, x^{M_i}Y) = x^{\alpha_i} F_{M_i}(1,Y) + x^{\alpha_i+\zeta} s(x,Y) + O(|x|^K + |x|^{KM_i}|Y|^K) \tag{3.1a}$$

Here $s(x,Y)$ is a polynomial in Y and a fractional power of x. Analogous expressions hold for various y derivatives of F; for example, we can write

$$\frac{\partial^2 F}{\partial y^2}(x, x^{M_i}Y) = x^{\alpha_i - 2M_i} \partial_{yy} F_{M_i}(1,Y) + x^{\alpha_i - 2M_i + \zeta} \tilde{s}(x,Y) + O(|x|^K + |x|^{KM_i}|Y|^K) \tag{3.1b}$$

The constant δ was chosen small enough that $|\partial_Y^{k_{ij}} F_{M_i}(1,Y)| > C$ on $(r-\delta, r+\delta)$ for some positive C, where $k_{ij} \leq d$. As a result, shrinking η if necessary we can assume that on $(0,\eta) \times (r-\delta, r+\delta)$ we have $\frac{\partial^2 F}{\partial y^2}(x, x^{M_i}Y) = x^{\alpha_i - 2M_i} S(x,Y)$ where

$$|\partial_Y^{k_{ij}-2} S(x,Y)| > C \tag{3.2}$$

We now let $x = X^{\frac{1}{M_i+1}}$, so that $\frac{\partial^2 F}{\partial y^2}(X^{\frac{1}{M_i+1}}, X^{\frac{M_i}{M_i+1}}Y)$ is of the form $X^{\frac{\alpha_i - 2M_i}{M_i+1}} T(X,Y)$ where

$$|\partial_Y^{k_{ij}-2} T(X,Y)| > C \tag{3.2$'$}$$

We make this coordinate change so that the change from (x,y) coordinates to (X,Y) has constant Jacobian determinant. This ensures that a power of the damping function is integrable in the (x,y) coordinates iff it is integrable in the (X,Y) coordinates. In the old coordinates, the damping function is $|\tilde{D}(x,y)|^{\delta z}$ times $x^{(\frac{dM_i - a_i - M_i b_i}{d})z} |\frac{\partial^2 F}{\partial y^2}(x,y)|^{\frac{z}{2}}$ so in the new coordinates it is of the form

$$|\bar{D}(X,Y)|^{\delta z} \times X^{\frac{dM_i - a_i - M_i b_i}{d(1+M_i)} z} |X^{\frac{\alpha_i - 2M_i}{M_i+1}} \frac{\partial^2 T}{\partial Y^2}(X,Y)|^{\frac{z}{2}}$$

$$= |\bar{D}(X,Y)|^{\delta z} \times X^{\frac{d\frac{\alpha_i}{2} - a_i - M_i b_i}{d(M_i+1)} z} |\frac{\partial^2 T}{\partial Y^2}(X,Y)|^{\frac{z}{2}} \tag{3.3}$$

Note that by (3.2′) the function $\frac{\partial^2 T}{\partial Y^2}(X,Y)$ has $(k_{ij}-2)$th derivative uniformly bounded below in the y variable, for fixed x, where $k_{ij} \leq d$. As a result, if $Re(z) = s > -\frac{2}{d-2}$, then $|\frac{\partial^2 T}{\partial Y^2}(X,Y)|^{\frac{z}{2}}$ is integrable in y with integral uniformly bounded in $Im(z)$. Thus for such z we have

$$\int_{r-\delta}^{r+\delta} \int_0^{\eta} |x^{\left(\frac{d\frac{\alpha_i}{2} - a_i - M_i b_i}{d(M_i+1)}\right)}|^{\frac{z}{2}} |\frac{\partial^2 T}{\partial Y^2}(X,Y)|^{\frac{z}{2}}\, dx\, dy < C \int_0^{\eta} x^{s\frac{d\frac{\alpha_i}{2} - a_i - M_i b_i}{d(M_i+1)}}\, dx \tag{3.4}$$

This will be finite if the exponent of x is greater than -1. Substituting $\alpha_i = a_i + M_i b_i$, the exponent in (3.4) is $s\frac{d-2}{2d}\frac{a_i+M_i b_i}{1+M_i}$. Note that $\frac{a_i+M_i b_i}{1+M_i}$ is the x-coordinate of the intersection of the bisectrix with the edge e_i, which is at most d. Hence $\frac{d-2}{2d}\frac{a_i+M_i b_i}{1+M_i} < \frac{d-2}{2}$. Therefore if $s > \frac{2}{d-2}$, the exponent in (3.4) is greater than -1 and thus the right-hand factor of (3.3) is integrable for $Re(z) = s$, uniformly in $Im(z)$. As in the previous argument, by making δ in the $|\bar{D}(X,Y)|^{\delta z}$ factor sufficiently small by Holder's inequality the same will be true for the entire damping factor (3.3). This what we needed to prove and we are done.

4. L^2 estimates.

We now move to proving the L^2 bounds needed in the proof of Theorem 1.1. As indicated in section 1, we will be utilizing Theorem 1.2 that follows from [SoSt]. Letting $\phi(q)d\sigma(q)$ be the surface measure of (1.1), we define the measure σ_z by $d\sigma_z(q) = e^{z^2}|h(x,y)|^z|D(x,y)|^{\delta z}\phi(q)d\sigma(q)$. Since we will using Theorem 1.2, we examine its Fourier transform $\hat{\sigma_z}(\lambda)$, given by

$$\hat{\sigma_z}(\lambda) = e^{z^2}\int e^{-i\lambda_1 g(x,y) - i\lambda_2 x - i\lambda_3 y}|h(x,y)|^z|D(x,y)|^{\delta z}\phi(x,y)\, dx\, dy \tag{4.1}$$

Here $\phi(x,y)$ denotes some cutoff function on a neighborhood of the origin. We always shift by (x_0, y_0), so that our integrals are over a small neighborhood of the origin. Thus up to an ignorable factor of magnitude 1, (4.1) is given by

$$\hat{\sigma_z}(\lambda) = e^{z^2}\int e^{-i\lambda_1 f(x,y) - i\lambda_2 x - i\lambda_3 y}|H(x, y-\psi(x))|^z|D(x,y)|^{\delta z}\phi^*(x,y)\, dx\, dy \tag{4.2}$$

In (4.2), $\psi(x)$ is the function taking $f(x,y)$ into its adapted coordinates and $D(x,y)$ now denotes the Hessian determinant of f at (x,y). For the analysis of the part of (4.2) coming from the D_i'', F_{ij}'', and G_{ij}'', we will transfer into the adapted coordinates of $f(x,y)$ and use Van der Corput-type arguments in the y-variable. For the E_{ij}'' we will remain in the original coordinates, and use a Van der Corput-type argument in the x-variable in conjunction with Lemma 2.4.

We will prove that the conditions of Theorem 1.2 hold by virtue of the following theorem, whose proof will comprise most of the rest of this paper.

Theorem 4.1. Suppose $s > 1$. Then if the constant δ used in the exponent of $|D(x,y)|$ is sufficiently small, then there are constants C, ϵ independent of $Im(z)$ such that if $Re(z) = s$ then for any multiindex α with $|\alpha| = 0, 1$ we have

$$|\partial^\alpha \hat{\sigma_z}(\lambda)| < C(1 + |\lambda|)^{-\frac{1}{2}-\epsilon} \tag{4.3}$$

Proof. We will only prove (4.3) for $|\alpha| = 0$ as the $|\alpha| = 1$ cases are identical other than having a different cutoff function $\phi(x,y)$. Recall that we are assuming that $f(0,0) = 0$ and $\nabla f(0,0) = 0$. So if $|\lambda_2|$ or $|\lambda_3|$ is the maximal $|\lambda_i|$ one may integrate by parts in x or y respectively and get that $|\hat{\sigma}(\lambda)| < C|\lambda|^{-1}$, which is better than the estimate that we need. Hence for the remainder of this paper we will always assume that $|\lambda_1|$ is at least as large as $|\lambda_2|$ and $|\lambda_3|$.

Let $\alpha(x)$ be an even function on $\mathbf{R}$ that is equal to 1 for $|x| \leq \frac{1}{2}$, zero for $|x| > 1$, and is monotone decreasing on $\mathbf{R}^+$. Let $\beta(x) = 1 - \alpha(x)$. For constants δ_1 and N_1 to be determined by our arguments, we express (4.2) as $I_1 + I_2 + I_3$, where

$$I_1(\lambda) = e^{z^2} \int e^{-i\lambda_1 f(x,y) - i\lambda_2 x - i\lambda_3 y} |H(x, y - \psi(x))|^z |D(x,y)|^{\delta z}$$

$$\times \alpha(|\lambda|^{N_1} D(x,y)) \phi^*(x,y)\, dx\, dy \tag{4.4a}$$

$$I_2(\lambda) = e^{z^2} \int e^{-i\lambda_1 f(x,y) - i\lambda_2 x - i\lambda_3 y} |H(x, y - \psi(x))|^z |D(x,y)|^{\delta z}$$

$$\times \big(\alpha(|\lambda|^{\delta_1} D(x,y) - \alpha(|\lambda|^{N_1} D(x,y))\big) \phi^*(x,y)\, dx\, dy \tag{4.4b}$$

$$I_3(\lambda) = e^{z^2} \int e^{-i\lambda_1 f(x,y) - i\lambda_2 x - i\lambda_3 y} |H(x, y - \psi(x))|^z |D(x,y)|^{\delta z}$$

$$\times \beta(|\lambda|^{\delta_1} D(x,y)) \phi^*(x,y)\, dx\, dy \tag{4.4c}$$

The analysis of $I_2(\lambda)$ will be the crux of the argument. The contribution to (4.3) due to $I_1(\lambda)$ is easily shown to decrease rapidly in $|\lambda|$. Specifically, since $D(x,y)$ is being assumed to be of finite-type in a neighborhood of the origin, if N_1 is large enough the measure of the points where $|D(x,y)| < |\lambda|^{-N_1}$ will be less than $\frac{1}{|\lambda|}$. As a result, the integrand of $I_1(\lambda)$ is nonzero on a set of measure at most $\frac{1}{|\lambda|}$. Since all the factors in (4.4a) are uniformly bounded on a line $Re(z) = s$ with $s > 1$, this gives that $|I_1(\lambda)| < C|\lambda|^{-1}$, better than what is needed.

Bounding $|I_3(\lambda)|$.

Note that on the support of the integrand of $I_3(\lambda)$, the Hessian determinant $D(x,y)$ is at least $\frac{1}{2}|\lambda|^{-\delta_1}$. The idea is that if δ_1 were actually zero, then on this support the Hessian would be bounded below and we would get an estimate $|I_3(\lambda)| < C|\lambda|^{-1}$.

Although δ_1 is not zero, if it is sufficiently small we still get an estimate $|I_3(\lambda)| < C|\lambda|^{-1-t}$ for any given small t, an estimate better than what is needed.

We proceed as follows. For a sufficiently small $c > 0$ to be determined by our arguments, we divide the support of the integrand of $I_3(\lambda)$ into squares of diameter $c|\lambda|^{-\delta_1}$. We will show that the contribution to $I_3(\lambda)$ from each such square is at most $c|\lambda|^{-\frac{3}{5}}$ if δ_1 is sufficiently small. Adding this over all these squares, this gives an estimate better than needed.

Let S be any such square. Since $D(x,y)$ is of finite type, we may let u and v be nonparallel directions such that for some k, $\partial_u^k D(x,y)$, $\partial_v^k D(x,y)$, and $\partial_u \partial_v^{k-1} D(x,y)$ are nonvanishing on the support of the integrand of $I_3(\lambda)$. We can similarly assume that there are k', k'' such that $\partial_u^{k'}\big(F^*(x, y-\psi(x))^2\big)$, $\partial_v^{k'}\big(F^*(x, y-\psi(x))^2\big)$, $\partial_u^{k''} f(x,y)$, and $\partial_u^{k''} f(x,y)$ are nonvanishing on any S. Let $a_1 = -\partial_u(\frac{\lambda_2}{\lambda_1}x + \frac{\lambda_3}{\lambda_1}y)$ and $a_2 = -\partial_v(\frac{\lambda_2}{\lambda_1}x + \frac{\lambda_3}{\lambda_1}y)$. Note a_1 and a_2 are constants. Define the sets S_1, S_2, and S_3 by

$$S_1 = \{(x,y) \in S : |\partial_u f(x,y) - a_1| > |\lambda|^{-\frac{1}{3}}\} \tag{4.5a}$$

$$S_2 = \{(x,y) \in S : |\partial_u f(x,y) - a_1| \leq |\lambda|^{-\frac{1}{3}}, |\partial_v f(x,y) - a_2| > |\lambda|^{-\frac{1}{3}}\} \tag{4.5b}$$

$$S_3 = \{(x,y) \in S : |\partial_u f(x,y) - a_1| \leq |\lambda|^{-\frac{1}{3}}, |\partial_v f(x,y) - a_2| \leq |\lambda|^{-\frac{1}{3}}\} \tag{4.5c}$$

Correspondingly, write the contributions to $I_3(\lambda)$ from S_1, S_2, and S_3 as $J_1(\lambda)$, $J_2(\lambda)$, and $J_3(\lambda)$ respectively. To analyze $J_1(\lambda)$, integrate the integrand of (4.4c) by parts in the u direction, integrating $\lambda_1(\partial_u f(x,y) - a_1)e^{-i\lambda_1 f(x,y) - i\lambda_2 x - i\lambda_3 y}$ in and differentiating $\frac{1}{\lambda_1(\partial_u f(x,y) - a_1)}$ times the remainder of the integrand. We get several terms depending on where the derivative lands. If it lands on the $\phi(x,y)$ factor, then each factor in the term is bounded above by a constant, with the exception of the $\frac{1}{\lambda_1(\partial_u f(x,y) - a_1)}$ factor, which is bounded in absolute value by $|\lambda_1|^{-\frac{2}{3}}$ on S_1. Hence $|J_1(\lambda)|$ is at most $C|\lambda_1|^{-\frac{2}{3}+2\delta_1}$, which is bounded by $C|\lambda|^{-\frac{3}{5}}$, the desired estimate.

Next, we consider the term where the derivative lands on the $\frac{1}{\lambda_1(\partial_u f(x,y) - a_1)}$ factor. We take absolute values of the entire integrand and bound it above by $\frac{C}{|\lambda|}\frac{\partial_u^2 f(x,y)}{(\partial_u f(x,y) - a_1)^2}$. We integrate this in the u direction as in the proof of the Van der Corput lemma; the assumed condition that $|\partial_u^{k''} f(x,y)|$ is bounded below ensures that we integrate over boundedly many intervals on which $\partial_u f(x,y) - a_1$ is monotone and thus $\frac{\partial_u^2 f(x,y)}{(\partial_u f(x,y) - a_1)^2}$ integrates back to $\frac{1}{(\partial_u f(x,y) - a_1)}$. We end out with a bound of $C|\lambda_1|^{-\frac{2}{3}+\delta_1} < C|\lambda_1|^{-\frac{3}{5}}$.

If the derivative lands on the $|D(x,y)|^{\delta z}$ factor, we argue similarly. We take absolute values and integrate in the u direction, this time using that $|\partial_u^k D(x,y)|$ is bounded below on the integrand to ensure that there are boundedly many intervals on which $D(x,y)$ is monotone and thus on which we can integrate back its u-derivative. In the (extremely rare) case that only $k = 0$ can be used, $|D(x,y)|^{\delta z}$ is a smooth function and the term behaves as in the case where the derivative lands on $\phi(x,y)$. If the derivative lands on $\beta(|\lambda|^{\delta_1} D(x,y))$ the argument we just used for the $|D(x,y)|^{\delta z}$ case works. One thing worth pointing out is that in these cases the presence of the z in the exponent leads to an additional factor of $C|Im(z)|$ upon differentiation; however, the presence of the e^{z^2} in the damping factor is more than enough to compensate.

Lastly, we consider the case where the derivative lands on the factor $|H(x, y - \psi(x))|^{z(\frac{1}{2} - \frac{1}{d^*})}$. Since this factor was defined differently on the different D_i, E_{ij}, etc, we split the square S into its intersections with the D_i'', E_{ij}'', F_{ij}'', and G_{ij}''. For anything other than a G_{ij}'' with $k_{ij} > 2$, the damping factor is a power of $F^*(x, y-\psi(x))$. The directions u and v were defined so that $F^*(x, y - \psi(x))^2$ has some nonvanishing higher order derivative in the u and v directions, so one can argue as above, breaking up the one-dimensional integration in the u or v variables into boundedly many intervals on which $F^*(x, y - \psi(x))$ is monotone.

On a G_{ij}'' with $k_{ij} > 2$, the damping factor was defined as $x^{M_i - \frac{a_i + M_i b_i}{d^*}} \partial_{yy} f(x,y)$. We can actually assume that u and v are such that the $(k_{ij} - 2)$th u and v derivatives of $x^{M_i - \frac{a_i + M_i b_i}{d^*}} \partial_{yy} f(x,y)$ are nonvanishing. To see why, first note that (2.19) gives that the $(k_{ij} - 2)$th y-derivative of $\partial_{yy} f(x,y)$ is bounded below by $Cx^{a_i + M_i(b_i - k_{ij})}$. On the other hand, by Lemma 2.3, (remembering that M_i is always at least 1 in generic adapted coordinates) on G_{ij}'' we have $|\partial^\alpha f(x,y)| \leq C_\alpha^{a_i + M_i b_i - M_i|\alpha|}$. Using these facts with the product rule , if u and v are close enough to the y direction, the $(k_{ij} - 2)$th derivative in the u or v direction of $x^{M_i - \frac{a_i + M_i b_i}{d^*}} \partial_{yy} f(x,y)$ will also be nonvanishing. Hence one can argue as in the previous paragraph and get the same upper bounds as before. We have now considered all possible places the derivative lands, concluding the proof of the desired upper bounds for $|J_1(\lambda)|$.

The bounds for $|J_2(\lambda)|$ are proven exactly as they were for $|J_1(\lambda)|$, replacing the roles of the u and v variables. The presence of the added condition $|\partial_u f(x,y) - a_1| \leq |\lambda|^{-\frac{1}{3}}$ in the domain, which does not have an analogue above, does not interfere with any of the above estimates; the condition that $\partial_u \partial_v^{k-1} D(x,y)$ is nonvanishing ensures that in any of the situations where one takes absolute values and does a Van der Corput type argument in the v direction, one still has boundedly many intervals.

We now move on to $J_3(\lambda)$. Consider the level sets of $\partial_u f(x,y)$ and $\partial_v f(x,y)$. The gradients of both functions are bounded below in absolute value by $C|D(x,y)|$, which is at least $\frac{1}{3}|\lambda|^{-\delta_1}$ on the square S if we chose the constant c in the diameter $c|\lambda|^{-\delta_1}$ of the squares sufficiently small. As a result, if c is small enough the level sets of both $\partial_u f(x,y)$ and $\partial_v f(x,y)$ do not self-intersect on S. Hence we may use $\partial_u f(x,y)$ and $\partial_v f(x,y)$ as coordinates on S. In particular, we may evaluate the measure of the set S_3 of $(4.5c)$ by changing into these coordinates in the integral of its characteristic function. The result is

$$|S_3| < C \min_S |D(x,y)|^{-1} |\lambda|^{-\frac{2}{3}} \tag{4.6}$$

So we conclude that $|S_3| < C'|\lambda|^{-\frac{2}{3}+\delta_1}$. Since the integrand of $J_3(\lambda)$ is uniformly bounded on $Re(z) = s$ for any $s > 1$, we conclude that

$$|J_3(\lambda)| \leq C''|\lambda|^{-\frac{2}{3}+\delta_1} < C''|\lambda|^{-\frac{3}{5}} \tag{4.7}$$

This gives the needed estimate. Adding the contributions from $|J_1(\lambda)|$, $|J_2(\lambda)|$, and $|J_3(\lambda)|$, we conclude that the contribution to $|I_3(\lambda)|$ from the square S is at most $C'''|\lambda|^{-\frac{3}{5}}$, and since $-\frac{3}{5} < -\frac{1}{2}$ we conclude that $|I_3(\lambda)|$ satisfies the bounds we need so long as δ_1 was chosen sufficiently small.

Estimating $|I_2(\lambda)|$.

We focus our attention on the main term $I_2(\lambda)$, given by $(4.4b)$. We divide the domain of $(4.4b)$ into squares of diameter $c|\lambda|^{-\delta_2}$, where c and δ_2 are small constants. For a given such square S, denote the corresponding term of $I_2(\lambda)$ by $I_2^S(\lambda)$. We will show that if c and δ_2 are sufficiently small, then for any such S we have $|I_2^S(\lambda)| < C|\lambda|^{-\frac{1}{2}-\epsilon}$, where ϵ is independent of c and δ_2, and C is independent of $Im(z)$ for $Re(z) = s > 1$. Since there are at most $c'|\lambda|^{2\delta_2}$ squares, as long as we make sure $\delta_2 < \frac{\epsilon}{2}$, this is enough to show that $I_2(\lambda)$ itself satisfies the bounds needed for Theorem 4.1. This subdivision into squares is useful because it allows us to replace $D(x,y)$ by a polynomial approximation of bounded degree which is therefore piecewise monotone in a direction in which we are integrating by parts, enabling us to use Van der Corput type arguments in such a direction.

We now perform this polynomial replacement. For a given S and positive integer N, let $D_S^N(x,y)$ be the polynomial in x and y consisting of the sum of the terms of degree at most N of $D(x,y)$'s Taylor expansion centered about the center of S. Thus on S we have

$$|D_S^N(x,y) - D(x,y)| < C|\lambda|^{-\delta_2 N} \tag{4.8}$$

As a result, on S we have

$$|\big(\alpha(|\lambda|^{\delta_1}D(x,y)) - \alpha(|\lambda|^{N_1}D(x,y)\big) - \big(\alpha(|\lambda|^{\delta_1}D_S^N(x,y)) - \alpha(|\lambda|^{N_1}D_S^N(x,y))\big)| < C|\lambda|^{N_1-\delta_2 N}$$

In particular, if N is chosen large enough we can make the exponent $N_1 - \delta_2 N$ appearing in (4.9) less than -1. Consequently, for the purposes of the analysis of $I_2^S(\lambda)$ we may replace $D(x, y - \psi(x))$ by $D_S^N(x,y)$ in the $\big(\alpha(|\lambda|^{\delta_1}D(x,y) - \alpha(|\lambda|^{N_1}D(x,y))\big)$ factor; the difference will contribute no more than $C|\lambda|^{-1}$ to $I_2^S(\lambda)$, and adding over all squares gives a result smaller than the bounds needed in Theorem 4.1.

We can do something similar for the $|D(x,y)|^{\delta z}$ factor. Namely, suppose N is taken large enough that in (4.8) we have

$$|D_S^N(x,y) - D(x,y)| < C|\lambda|^{-2N_1}$$

Then since $|D(x,y)| \geq \frac{1}{2}|\lambda|^{-N_1}$ when the integrand of $(4.4b)$ is nonzero, if $|\lambda|$ is large enough we may use the Taylor expansion of $|x|^{\delta z}$ about $x = D(x,y)$ to obtain

$$\big||D(x,y)|^{\delta z} - |D_S^N(x,y)|^{\delta z}\big| < C|Im(z)||\lambda|^{-N_1(\delta Re(z)-1)-2N_1} \tag{4.9}$$

As a result, since $Re(z) > 1$, as long as $N_1 > 1$, we have an estimate

$$\big||D(x,y)|^{\delta z} - |D_S^N(x,y)|^{\delta z}\big| < C|Im(z)||\lambda|^{-1} \tag{4.10}$$

The e^{z^2} is more than enough to take care of the $|Im(z)|$ factor in (4.10), and the exponent -1 is less than $-\frac{1}{2}$. Consequently, we may replace $|D(x,y)|^{\delta z}$ by $|D_S^N(x,y)|^{\delta z}$ in the analysis of $I_2^S(\lambda)$; the difference added over all squares S contributes less than the bounds needed for Theorem 4.1.

We have now shown that for the purposes of our future arguments, we may adjust our notation and assume $I_2^S(\lambda)$ is given by

$$I_2^S(\lambda) = e^{z^2}\int e^{-i\lambda_1 f(x,y) - i\lambda_2 x - i\lambda_3 y}|H(x, y-\psi(x))|^z|D_S^N(x,y)|^{\delta z}$$

$$\times\big(\alpha(|\lambda|^{\delta_1}D_S^N(x,y)) - \alpha(|\lambda|^{N_1}D_S^N(x,y))\big)\phi^*(x,y)\,dx\,dy \tag{4.11}$$

We divide the domain of integration of (4.11) into the intersections of S with the D_i'', E_{ij}'', F_{ij}'', and G_{ij}'' and denote the corresponding term of $I_2(\lambda)$ by $I_2^{D_i}(\lambda)$, $I_2^{E_{ij}}(\lambda)$, $I_2^{F_{ij}}(\lambda)$, and $I_2^{G_{ij}}(\lambda)$. (Recall $D_i'' = \{(x,y) : (x, y-\psi(x)) \in D_i\}$ with analogous definitions for the other regions). We suppress the S since the bounds we will prove, given in the statement of Theorem 4.1, are independent of S. We will only consider those regions for which $x > 0$ as the $x < 0$ ones are entirely analogous. We now focus our attention on the analysis of the $I_2^{D_i}(\lambda)$.

Bounds for $|I_2^{D_i}(\lambda)|$.

Recalling that $|H(x,y)| = F^*(x,y)^{\frac{1}{2}-\frac{1}{d^*}}$ on a D_i, if we change coordinates from (x,y) to $(x, y+\psi(x))$ in (4.11) we obtain

$$I_2^{D_i}(\lambda) = e^{z^2}\int_{S\cap D_i} e^{-i\lambda_1 F(x,y) - i\lambda_2 x - i\lambda_3(y+\psi(x))}F^*(x,y)^{(\frac{1}{2}-\frac{1}{d^*})z}|D_S^N(x, y+\psi(x))|^{\delta z}$$

$$\big(\alpha(|\lambda|^{\delta_1}D_S^N(x, y+\psi(x))) - \alpha(|\lambda|^{N_1}D_S^N(x, y+\psi(x)))\big)\phi^{**}(x,y)\,dx\,dy \tag{4.12}$$

Here $\phi^{**}(x,y)$ denotes a new cutoff function on a neighborhood of the origin, and $F(x,y) = f(x, y+\psi(x))$ is in generic adapted coordinates not satisfying the exceptional situations of Theorem 1.1. We slightly abuse notation in (4.12) in that S now denotes the square in the new coordinates. We now decompose the domain of (4.12) into dyadic rectangles. We only consider those rectangles in the upper right quadrant as the other quadrants are done the same way. For a given dyadic rectangle $J_{kl} = [2^{-k-1}, 2^{-k}] \times [2^{-l-1}, 2^{-l}]$, we use the shorthand by I_{kl} to denote the corresponding term of $I_2^{D_i}(\lambda)$, given by

$$I_{kl} = e^{z^2}\int_{S\cap D_i\cap J_{kl}} e^{-i\lambda_1 F(x,y) - i\lambda_2 x - i\lambda_3(y+\psi(x))}F^*(x,y)^{(\frac{1}{2}-\frac{1}{d^*})z}|D_S^N(x, y+\psi(x))|^{\delta z}$$

$$\big(\alpha(|\lambda|^{\delta_1} D_S^N(x, y+\psi(x))) - \alpha(|\lambda|^{N_1} D_S^N(x, y+\psi(x)))\big)\phi^{**}(x,y)\,dx\,dy \tag{4.13}$$

We will analyze (4.13) by imitating the proof of Van der Corput's lemma in the y direction. Our objective is to show that (4.13) is bounded by $C(1+|\lambda|)^{-\frac{1}{2}-\epsilon}$ as in the statement of Theorem 4.1. The second y derivative of the phase function in (4.13) is given by $\lambda_1 \partial_{yy} F(x,y)$, and by (2.14), if the vertex of $N(F)$ corresponding to D_i is written as (a_i, b_i), then on D_i we have $|\partial_{yy} F(x,y)| > c|x|^{a_i}|y|^{b_i - 2}$. Since $x \sim 2^{-k}$ and $y \sim 2^{-l}$ on J_{kl} we can write this as

$$|\partial_{yy} F(x,y)| > c' \frac{1}{(2^{-l})^2}(2^{-k})^{a_i}(2^{-l})^{b_i} \tag{4.14}$$

As in the proof of the Van der Corput theorem for functions with nonvanishing second derivative, we will split the integral (4.13) into two parts. The first is the part where $|\lambda_1 \partial_y F(x,y) + \lambda_3| < |\lambda|^{\frac{1}{2}} 2^{\frac{-a_i k - (b_i - 2)l}{2}}$, and the second is the part where $|\lambda_1 \partial_y F(x,y) + \lambda_3| \geq |\lambda|^{\frac{1}{2}} 2^{\frac{-a_i k - (b_i - 2)l}{2}}$. Call the resulting integrals K_1 and K_2, so that $K_1 + K_2 = I_{kl}$. We will bound K_1 by taking absolute values and integrating, and K_2 by performing an integration by parts.

We start with K_1. The integrand of (4.13) is bounded in absolute value by a constant times $F^*(x,y)^{Re(z)(\frac{1}{2} - \frac{1}{d^*})}|D_S^N(x, y+\psi(x))|^{\delta Re(z)}$. By (2.20), $F^*(x,y) < C|x^{a_i} y^{b_i}| \leq C2^{-ka_i - lb_i}$, and on the domain of (4.13) we have $|D_S^N(x, y+\psi(x))|^{\delta z} < C'|\lambda|^{-\delta\delta_1 Re(z)}$ Hence if s denotes $Re(z)$, the integrand of (4.13) is at most

$$C''|\lambda|^{-\delta\delta_1 s} 2^{(-ka_i - lb_i)s(\frac{1}{2} - \frac{1}{d^*})} \tag{4.15}$$

Since by (4.14) the absolute value of the y-derivative of $\lambda_1 \partial_y F(x,y) + \lambda_3$ is at least $c'|\lambda| 2^{-ka_i - l(b_i - 2)}$ we have

$$|\{y : |\lambda_1 \partial_y F(x,y) + \lambda_3| < |\lambda|^{\frac{1}{2}} 2^{\frac{-a_i k - (b_i-2)l}{2}}\}| < |\lambda|^{-\frac{1}{2}} 2^{\frac{ka_i + l(b_i - 2)}{2}} \tag{4.16}$$

Thus bounding the y integral of K_1 by (4.15) times the measure (4.16) and then integrating the result in x, we obtain

$$|K_1| < C'''|\lambda|^{-\frac{1}{2} - \delta\delta_1 s} 2^{(-ka_i - lb_i)(\frac{s-1}{2} - \frac{s}{d^*})} 2^{-k-l} \tag{4.17}$$

We now turn to K_2 and show that K_2 also satisfies the upper bounds of (4.17). We integrate the integrand in (4.13) by parts in y, integrating the factor $(\lambda_1 \partial_y F(x,y) + \lambda_3)e^{-i\lambda_1 F(x,y) - i\lambda_2 x - i\lambda_3(y - \psi(x))}$ and differentiating $\frac{1}{\lambda_1 \partial_y F(x,y) + \lambda_3}$ times the rest of the integrand. We get several terms depending on where the derivative lands. If the derivative lands on $\phi^*(x,y)$, the absolute value of the integrand in the resulting term is bounded by

$$CF^*(x,y)^{s(\frac{1}{2} - \frac{1}{d^*})}|D_S^N(x, y+\psi(x))|^{\delta s}|\lambda|^{-\frac{1}{2}} 2^{\frac{a_i k + (b_i - 2)l}{2}} \tag{4.18}$$

Bounding $F^*(x,y) < C2^{-ka_i - lb_i}$ and $|D_S^N(x, y+\psi(x))| < C|\lambda|^{-\delta_1}$ as in the analysis of K_1, we get that (4.18) is bounded by

$$C' 2^{(-ka_i - lb_i)(\frac{s-1}{2} - \frac{s}{d^*})} 2^{-l}|\lambda|^{-\frac{1}{2} - \delta\delta_1 s} \tag{4.19}$$

Integrating (4.19) over $S \subset I_{kl}$ multiplies this by at most $C2^{-k-l}$, so the resulting term is at most

$$C' 2^{(-ka_i - lb_i)(\frac{s-1}{2} - \frac{s}{d^*})} 2^{-k-2l} |\lambda|^{-\frac{1}{2} - \delta\delta_1 s} \tag{4.20}$$

Note this is better than the estimate (4.17). We next consider the case where the y-derivative lands on the $\frac{1}{\lambda_1 \partial_y F(x,y) + \lambda_3}$ factor, turning it into $-\frac{\lambda_1 \partial_{yy} F(x,y)}{(\lambda_1 \partial_y F(x,y) + \lambda_3)^2}$. We take absolute values and integrate in the y variable as in the proof of the Van der Corput lemma, bounding the other factors as was done for (4.19). Since by (4.14) the function $\partial_{yy} F(x,y)$ is never zero on the domain of integration, we have at most finitely many intervals of integration on each of which $|\frac{\lambda_1 \partial_{yy} F(x,y)}{(\lambda_1 \partial_y F(x,y) + \lambda_3)^2}|$ integrates back into $\pm \frac{1}{\lambda_1 \partial_y F(x,y) + \lambda_3}$. Hence the resulting term, as well as the endpoint terms, will be bounded by (4.20), except divided by the y-width 2^{-l}. We conclude that this term is bounded by (4.17), namely

$$C'' 2^{(-ka_i - lb_i)(\frac{s-1}{2} - \frac{s}{d^*})} 2^{-k-l} |\lambda|^{-\frac{1}{2} - \delta\delta_1 s} \tag{4.21}$$

If the y-derivative lands on either the $|D_S^N(x, y + \psi(x))|^{\delta z}$ or $\big(\alpha(|\lambda|^{\delta_1} D_S^N(x, y + \psi(x))) - \alpha(|\lambda|^{N_1} D_S^N(x, y + \psi(x)))\big)$ factors one estimates the resulting term in very much the same way; the fact that $D_S^N(x,y)$ is a polynomial and α is monotone ensures that the Van der Corput lemma proof still applies and we will have boundedly many intervals of integration on which the appropriate derivative is nonvanishing. Similarly, since $F^*(x,y)^2$ is a polynomial, one can deal with the term where the derivative lands on the damping factor $F^*(x,y)^{z(\frac{1}{2} - \frac{1}{d^*})}$ in a similar fashion. It should be pointed out that in taking these derivatives we do incur a factor of $C|Im(z)|$, but this is more than compensated for by the e^{z^2} factor. Hence we once again get the upper bound (4.21). Adding all terms together, we see that $|K_2|$ and therefore $|I_{kl}|$ is bounded by (4.17), the estimate we need.

We rewrite (4.17) in an especially useful form. Recall that by (2.20), on S we have $C2^{(-ka_i - lb_i)} < F^*(x,y) < C'2^{(-ka_i - lb_i)}$. So we have just shown that

$$|I_{kl}| < C \int_{J_{kl}} |\lambda|^{-\frac{1}{2} - \frac{\delta\delta_1 s}{2}} F^*(x,y)^{\frac{s-1}{2} - \frac{s}{d^*}} \, dx \, dy \tag{4.22}$$

We now break into cases $d^* \leq 2$, and $d^* > 2$, starting with the latter. Adding (4.22) over all rectangles, we obtain that $|I_2^{D_i}(\lambda)|$ is at most

$$C|\lambda|^{-\frac{1}{2} - \frac{\delta\delta_1 s}{2}} \int_{[0,1]\times[0,1]} F^*(x,y)^{\frac{s-1}{2} - \frac{s}{d}} \, dx \, dy \tag{4.23}$$

Note that if $s > 1$, then $\frac{s-1}{2} - \frac{s}{d} > -\frac{1}{d}$, and thus since $F^*(x,y)^t$ is integrable over $[0,1] \times [0,1]$ for all $t > -\frac{1}{d}$, the integral in (4.23) is finite and we obtain that $|I_2^{D_i}(\lambda)|$ is bounded by $C|\lambda|^{-\frac{1}{2} - \frac{\delta\delta_1 s}{2}}$. Since the exponent here is less than $-\frac{1}{2}$, this gives what is needed for Theorem 4.1.

Moving on to the $d^* = 2$ case, (4.22) becomes

$$|I_{kl}| < C \int_{J_{kl}} |\lambda|^{-\frac{1}{2} - \frac{\delta\delta_1 s}{2}} F^*(x,y)^{-\frac{1}{2}} \, dx \, dy \tag{4.24}$$

Since the damping factor is just $|D_S^N(x, y+\psi(x))|^{\delta z}$ when $d^* = 2$, from (4.13) we get

$$I_{kl} = e^{z^2} \int_{S\cap D_i \cap J_{kl}} e^{-i\lambda_1 F(x,y) - i\lambda_2 x - i\lambda_3 (y+\psi(x))} |D_S^N(x, y+\psi(x))|^{\delta z}$$

$$\times \big(\alpha(|\lambda|^{\delta_1} D_S^N(x, y+\psi(x))) - \alpha(|\lambda|^{N_1} D_S^N(x, y+\psi(x)))\big)\phi^{**}(x,y)\, dx\, dy \tag{4.25}$$

Note that due to the cutoff and the presence of the $|D_S^N(x, y+\psi(x))|^{\delta z}$ in the integrand of (4.25), this integrand is at most $|\lambda|^{-\delta\delta_1 s}$. So just by taking absolute values and integrating we get

$$|I_{kl}| < C|\lambda|^{-\delta\delta_1 s} 2^{-k-l} \tag{4.26a}$$

$$< C|\lambda|^{\frac{-\delta\delta_1 s}{2}} 2^{-k-l} \tag{4.26b}$$

Combining this with (4.24), we get

$$|I_{kl}| < C|\lambda|^{\frac{-\delta\delta_1 s}{2}} \int_{[2^{-k-1}, 2^{-k}]\times[2^{-l-1}, 2^{-l}]} \min(1, |\lambda F^*(x,y)|^{-\frac{1}{2}})\, dx\, dy \tag{4.27}$$

Adding this up over all j and k, we obtain that $|I_2^{D_i}(\lambda)|$ is at most

$$C|\lambda|^{\frac{-\delta\delta_1 s}{2}} \int_{[0,1]\times[0,1]} \min(1, |\lambda F^*(x,y)|^{-\frac{1}{2}})\, dx\, dy \tag{4.28}$$

Since $(d,d) \in N(F)$, (d,d) is a convex combination of vertices of $N(F)$. So since $F^*(x,y)$ is comparable to the sum of $|x^a y^b|$ over vertices (a,b) of $N(F)$, we have $F^*(x,y) > C|x^d y^d|$. Since we are assuming $d \leq 2$ here, we conclude that $F^*(x,y) > Cx^2y^2$ and as a result (4.28) is bounded by

$$C|\lambda|^{\frac{-\delta\delta_1 s}{2}} \int_{[0,1]\times[0,1]} \min\left(1, \frac{C'}{|\lambda|^{\frac{1}{2}} xy}\right) dx\, dy \tag{4.29}$$

A direct calculation reveals that the right hand side is bounded above by $C''|\lambda|^{-\frac{1}{2}}(\ln|\lambda|)^2$ (The integral over $[|\lambda|^{-\frac{1}{2}}, 1] \times [|\lambda|^{-\frac{1}{2}}, 1]$ is bounded by a constant times the integral of $\frac{1}{|\lambda|^{\frac{1}{2}} xy}$ over this region, while the integral over the remaining region is bounded by its area). As a result, (4.29) is bounded by

$$C''|\lambda|^{-\frac{1}{2} - \frac{\delta\delta_1 s}{2}} (\ln|\lambda|)^2 \tag{4.30}$$

Since the exponent here is less than $-\frac{1}{2}$ we have proved the desired bounds for the $|I_2^{D_i}(\lambda)|$.

Bounds for $|I_2^{E_{ij}}(\lambda)|$.

Note that $I_2^{E_{ij}}(\lambda)$ is given by

$$I_2^{E_{ij}}(\lambda) = e^{z^2} \int_{S\cap E''_{ij}} e^{-i\lambda_1 f(x,y) - i\lambda_2 x - i\lambda_3 y} |F^*(x, y-\psi(x))|^{z(\frac{1}{2} - \frac{1}{d^*})} |D_S^N(x,y)|^{\delta z}$$

$$\big(\alpha(|\lambda|^{\delta_1} D_S^N(x,y)) - \alpha(|\lambda|^{N_1} D_S^N(x,y))\big)\phi^*(x,y)\,dx\,dy \tag{4.31}$$

As we did with $I_2^{D_i}(\lambda)$, we break the domain of integral (4.31) into rectangles $J_{kl} = [2^{-k-1}, 2^{-k}] \times [2^{-l-1}, 2^{-l}]$. Denote the corresponding term of (4.31) by I_{kl}, so that $\sum_{kl} I_{kl} = I_2^{E_{ij}}(\lambda)$.

Note that by (2.21) we have

$$F^*(x,y) > Cx^{a_i+M_ib_i} \tag{4.32}$$

As before (a_i, b_i) denotes the upper vertex of e_i. Recall that E_{ij} lies between $y = (r-\eta)x^{M_i}$ and $y = (r+\eta)x^{M_i}$ for some r and η such that $F_{e_i}(1,r) \neq 0$, and that by definition of E_{ij}, $\psi(x)$ has a zero of order at least M_i at $x = 0$. Consequently, $|y - \psi(x)| < Cx^{M_i}$ on E_{ij}. Thus by (2.11), $F^*(x, y - \psi(x)) < Cx^{a_i+b_iM_i}$ on E_{ij}. Combining with (4.32) we get

$$F^*(x, y-\psi(x)) < CF^*(x,y) \tag{4.33a}$$

By Lemma 2.4, on the domain of (4.31) we have

$$|\partial_{xx} f(x,y)| > C\frac{1}{x^2}x^{a_i}(x^{M_i})^{b_i} \tag{4.33b}$$

Equation (4.33a) shows that the damping function $F^*(x, y-\psi(x))$ satisfies the same upper bounds that the damping function $F^*(x,y)$ did in the D_i case. Equation (4.33b) shows the same thing for the phase (cf (4.14)), reversing the roles of the x and y derivatives. Furthermore, the functions that need to be piecewise monotone in x with boundedly many pieces in order to perform the Van der Corput argument do satisfy this; $D_S^N(x,y)$ is a polynomial and the second x derivative of $-i\lambda_1 f(x,y) - i\lambda_2 x - i\lambda_3 y$ is nonvanishing by (4.33b). Hence by repeating the D_i argument, reversing the roles of the x and y variables, we get that I_{kl} is bounded by (4.17). Adding this up like before gives that as in (4.30), $|I_2^{E_{ij}}(\lambda)|$ is bounded by $C''|\lambda|^{-\frac{1}{2}-\frac{\delta\delta_1 s}{2}}(\ln|\lambda|)^2$, the estimate we need.

Bounds for $|I_2^{F_{ij}}(\lambda)|$.

Recall the set F_{ij} is of the form $\{(x,y) : 0 < x < \eta, |y - rx^{M_i}| < \nu|x|^{M_i}\}$, where $F_{e_i}(1,y)$ has a zero of order 1 at $y = r$. Define $G(x,y) = F(x, y + rx^{M_i})$. Thus $G(x,y)$ is a function on the set $H_{ij} = \{(x,y) : 0 < x < \eta, |y| < \nu|x|^{M_i}\}$ such that $G_{e_i}(1,y)$ has a zero of order 1 at $y = 0$. Thus $N(\partial_y G)$ has an edge with equation $x + M_iy = a_i + M_ib_i - M_i$ that intersects the x axis. Conseqently, $N(\frac{\partial^2 G}{\partial x \partial y})$ has an edge with equation $x + M_iy = a_i + M_ib_i - M_i - 1$ intersecting the x axis. Hence assuming η was chosen sufficiently small, by Lemma 2.3 we may conclude that on H_{ij} we have

$$\left|\frac{\partial^2 G}{\partial x \partial y}(x,y)\right| > Cx^{a_i+M_ib_i-M_i-1}$$

We rewrite this as

$$\left|\frac{\partial^2 G}{\partial x \partial y}(x,y)\right| > C\frac{1}{x(x^{M_i})}x^{a_i}(x^{M_i})^{b_i} \tag{4.34}$$

Letting $\tilde{\psi}(x) = \psi(x) + rx^{M_i}$, we do a change of variables from y to $y + \tilde{\psi}(x)$ and write $I_2^{F_{ij}}(\lambda)$ as

$$I_2^{F_{ij}}(\lambda) = e^{z^2} \int_{S \cap H_{ij}} e^{-i\lambda_1 G(x,y) - i\lambda_2 x - i\lambda_3 (y + \tilde{\psi}(x))} |F^*(x, y + rx^{M_i})|^{z(\frac{1}{2} - \frac{1}{d^*})}$$

$$\times |D_S^N(x, y + \tilde{\psi}(x))|^{\delta z} \big(\alpha(|\lambda|^{\delta_1} D_S^N(x, y + \tilde{\psi}(x))) - \alpha(|\lambda|^{N_1} D_S^N(x, y + \tilde{\psi}(x)))\big) \phi^{***}(x, y)\, dx\, dy \tag{4.35}$$

As with the D_i, the S under the integral symbol now denotes the square in the new coordinates. By (2.21), on H_{ij} we have

$$F^*(x, y + rx^{M_i}) < Cx^{a_i}(x^{M_i})^{b_i} \tag{4.36}$$

We now break the domain of integration of (4.35) up into rectangles J_k of the form $[2^{-k-1}, 2^{-k}] \times [-\nu 2^{-kM_i}, \nu 2^{-kM_i}]$, and let I_k' the the portion of (4.35) coming from J_k. Equation (4.36) shows that the damping function $F^*(x, y + rx^{M_i})$ in (4.35) satisfies the same upper bounds the damping function did on the I_{kM_k} rectangle for the the D_i. (The $x \sim 2^{-k}$ rectangle of the "lower edge" of D_i). As for the phase, instead of having a lower bound on a second y derivative as in (4.14), we have the substitute (4.34). We still may argue as for the I_{kM_k} rectangle in the D_i case, but with one difference. In the analysis of the term called K_1 below (4.14), instead of bounding the measure of a sublevel set of $|\lambda_1 \partial_y G(x, y) + \lambda_3|$ in the y-variable and integrating with respect to x, one bounds the measure of the same sublevel set in the x variable using (4.34) and then integrates the result with respect to y.

Furthermore, all relevant factors are piecewise monotone with boundedly many pieces. The function $D_S^N(x, y+\tilde{\psi}(x))$ is a polynomial in y of bounded degree, as is $F^*(x, y+rx^{M_i})^2$, while since $G(x, y)$ is just a y shift of $F(x, y)$ by $\psi(x)$, if $(0, a)$ denotes the upper vertex of $N(F)$ then $\partial_y G(x, y)$ has nonvanishing $(a-1)$th y derivative.

Hence after making the above adjustment to the I_{kM_ik} argument of the D_i case, for a given k we get the bounds (4.17) for I_k'. (The arguments there did not require l to be an integer). Adding over all k, as for the $|I_2^{D_i}(\lambda)|$ we get that $|I_2^{F_{ij}}(\lambda)|$ is bounded by $C''|\lambda|^{-\frac{1}{2} - \frac{\delta\delta_1 s}{2}}(\ln|\lambda|)^2$, the needed estimate.

Bounds for $|I_2^{G_{ij}}(\lambda)|$.

For the $I_2^{G_{ij}}(\lambda)$, we separate the $k_{ij} = 2$ and $k_{ij} > 2$ cases as the damping factors are different in these two situations. First, we suppose $k_{ij} = 2$. Then $I_2^{G_{ij}}(\lambda)$ is given by

$$I_2^{G_{ij}}(\lambda) = e^{z^2} \int_{S \cap G_{ij}} e^{-i\lambda_1 F(x,y) - i\lambda_2 x - i\lambda_3 (y + \psi(x))} F^*(x, y)^{z(\frac{1}{2} - \frac{1}{d^*})} |D_S^N(x, y + \psi(x))|^{\delta z}$$

$$\big(\alpha(|\lambda|^{\delta_1} D_S^N(x, y + \psi(x))) - \alpha(|\lambda|^{N_1} D_S^N(x, y + \psi(x)))\big) \phi^{**}(x, y)\, dx\, dy \tag{4.37}$$

Observing that $|y| < Cx^{M_i}$ on G_{ij}, we divide the domain of (4.37) into rectangles J_k of the form $[2^{-k-1}, 2^{-k}] \times [C_0 2^{-kM_i}, C_1 2^{-kM_i}]$, and let I_k' be the corresponding piece of $I_2^{G_{ij}}(\lambda)$, so that $\sum_k I_k' = I_2^{G_{ij}}(\lambda)$.

Note that the integrand in (4.37) is the same as that of (4.13) for the D_i case. In particular, the damping function is the same as in the D_i case. Also, by (2.19) on G_{ij} we have the following analogue of (4.14):

$$|\partial_{yy} F(x,y)| > C\frac{1}{(x_i^M)^2} x^{a_i}(x^{M_i})^{b_i} \tag{4.38}$$

As a result, all estimates used in the D_i case for the I_{kl} rectangle, setting $l = kM_i$ (the lower edge of D_i) hold for the term I_k'. Thus $|I_k'|$ is bounded by $C|I_{kM_ik}|$, and adding over all k we recover $C|\lambda|^{-\frac{1}{2}-\frac{\delta\delta_1 s}{2}}(\ln|\lambda|)^2$ as an upper bound for $|I_2^{G_{ij}}(\lambda)|$. This completes the proof for the $k_{ij} = 2$ case.

We may now assume $k_{ij} > 2$, focusing our attention for now on the case when $N(F)$ has multiple vertices. Here, $I_2^{G_{ij}}(\lambda)$ is given by

$$I_2^{G_{ij}}(\lambda) = e^{z^2}\int_{S\cap G_{ij}} e^{-i\lambda_1 F(x,y) - i\lambda_2 x - i\lambda_3(y+\psi(x))}\big[x^{M_i - \frac{a_i+M_ib_i}{d}}|\frac{\partial^2 F}{\partial y^2}(x,y)|^{\frac{1}{2}}\big]^z$$

$$|D_S^N(x, y+\psi(x))|^{\delta z}\big(\alpha(|\lambda|^{\delta_1} D_S^N(x,y+\psi(x))) - \alpha(|\lambda|^{N_1} D_S^N(x,y+\psi(x)))\big)\phi^{**}(x,y)\,dx\,dy \tag{4.39}$$

We divide the domain of (4.39) into rectangles J_k as in the above $k_{ij} = 2$ case. and again let I_k' be the corresponding piece of $I_2^{G_{ij}}(\lambda)$. Observe that by (2.23), there is some C_0 such that the magnitude of the bracketed expression in (4.39) (which is the same as the $H(x,y)$ in (2.23)) is bounded by $C_0(x^{a_i+M_ib_i})^{\frac{1}{2}-\frac{1}{d}}$. Thus we may write $I_k' = \sum_{l=0}^\infty P_{kl}$, where P_{kl} is the portion of the integral over J_k where $|H(x,y)|$ is between $2^{-l+1}C_0(x^{a_i+M_ib_i})^{\frac{1}{2}-\frac{1}{d}}$ and $2^{-l}C_0(x^{a_i+M_ib_i})^{\frac{1}{2}-\frac{1}{d}}$. We will now bound each P_{kl}. To this end, note that on the domain of P_{kl}, by the definition of $H(x,y)$ and the P_{kl} we have

$$C_0^2 2^{-2l-2}(x^{a_i+M_ib_i})^{1-\frac{2}{d}} < x^{2M_i - \frac{2a_i+2M_ib_i}{d}}|\frac{\partial^2 F}{\partial y^2}(x,y)| < C_0^2 2^{-2l}(x^{a_i+M_ib_i})^{1-\frac{2}{d}} \tag{4.40}$$

Solving for $|\frac{\partial^2 F}{\partial y^2}(x,y)|$, we get

$$C_1 2^{-2l}\frac{1}{(x^{M_i})^2}x^{a_i+M_ib_i} < |\frac{\partial^2 F}{\partial y^2}(x,y)| < C_1' 2^{-2l}\frac{1}{(x^{M_i})^2}x^{a_i+M_ib_i} \tag{4.41}$$

One now bounds P_{kl} by integrating by parts in y in the portion of (4.39) corresponding to P_{kl}. One proceeds exactly as for the I_{kM_ik} term of the D_i (the $x \sim 2^{-k}$ rectangle of the "lower edge" of D_i), except instead of using $|\frac{\partial^2 F}{\partial y^2}(x,y)| > C\frac{1}{(x^{M_i})^2}x^{a_i+M_ib_i}$ from (4.14) one uses (4.41). This gives us an additional factor of $C2^l$ in the resulting bounds for the integral. This however is compensated by the damping factor, which by the definition of P_{kl} is bounded by $C2^{-lRe(z)}$ times the damping factor used for the I_{kM_ik} term in the D_i case. Thus the overall integral is bounded by $C2^{l(1-Re(z))}$ times what is obtained for the I_{kM_ik} term in the D_i case. We do not have to worry about whether each factor in (4.37) is boundedly piecewise monotone in y in our integrations by parts; the only new element in this regard is $\frac{\partial^2 F}{\partial y^2}(x,y)$, whose $(k_{ij}-2)$th y derivative is nonvanishing.

Since $Re(z) > 1$, we conclude $\sum_l P_{kl}$ is bounded by a constant times the estimate obtained for the I_{kM_ik} term in the D_i situation, and adding this over all k gives

$$|I_2^{G_{ij}}(\lambda)| \leq \sum_{kl} |P_{kl}| < C''|\lambda|^{-\frac{1}{2}-\frac{\delta\delta_1 s}{2}}(\ln|\lambda|)^2 \tag{4.42}$$

This is the estimate we seek. The above argument was for when $N(F)$ has multiple vertices, but when $N(F)$ just has one vertex the following simplified version of this argument works. In the one vertex situation, $|H(x,y)| = |\frac{\partial^2 F}{\partial y^2}(x,y)|$. This time we let P_l be the portion of the integral defining $I_2^{G_{ij}}(\lambda)$ over the set where $|H(x,y)|$ is between $C_0 2^{-l-1}$ and $C_0 2^{-l}$, where C_0 denotes the maximum value of $|H(x,y)|$. Like above, for P_l the decreased second y derivative of the phase gives an additional factor of $C2^l$ which is more than compensated by the additional $C2^{-lRe(z)}$ factor coming from the damping function. Adding over all l, we recover (4.42). This completes the proof of the bounds for the $|I_2^{G_{ij}}(\lambda)|$, which in turn completes the proof of Theorem 4.1.

The proof of Theorem 1.1.

We may now finish the proof of Theorem 1.1 in short order. First suppose $d(F) > 2$. For any $\eta > 0$, Theorem 3.1 says that on the line $Re(z) = -\frac{2}{d(F)-2} + \eta$, M_z is bounded on L^∞ with uniform constant, while Theorem 4.1 in conjunction with Theorem 1.2 says that on $Re(z) = 1 + \eta$, M_z is bounded on L^2 with uniform constant. Using interpolation for maximal operators (see Ch. 11 of [St2]), we have that M_0 is bounded on $L^{d(F)+\eta'}$ where $\eta' \rightarrow 0$ as $\eta \rightarrow 0$. Thus we conclude M_0 is bounded on L^p for all $p > d(F)$. Since $d(F) = h(q_0) = \max(2, h(q_0))$, this gives Theorem 1.1 for $d(F) > 2$.

On the other hand, if $d(F) = 2$, Theorem 3.1 says that on any vertical line $Re(z) = s$, M_z is bounded on L^∞ with uniform constant, and Theorem 4.1 still applies on a line $Re(z) = 1 + \eta$. Thus interpolation now gives the result obtained by letting $d(F)$ approach 2 in the previous paragraph, namely that M is bounded on L^p for $p > 2 = \max(h(q_0), 2)$. This completes the proof of Theorem 1.1.

5. References.

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