

## Stationary Phase

Given

$$F(k) = \int_a^b g(x)e^{ik\phi(x)} dx \quad (1)$$

A question one might ask is how does this integral behave for large  $k$ . The simplest way to attack this is to just integrate by parts. If  $\phi'(x) \neq 0$  rewrite (1)

$$F(k) = \int_a^b g(x) \frac{ik\phi'(x)}{ik\phi'(x)} e^{ik\phi(x)} dx$$

Recall integration by parts

$$\int u dv = vu - \int v du$$

Then clearly  $u = g(x)/ik\phi'(x)$ , and  $dv = ik\phi'(x)e^{ik\phi(x)}$  will allow us to write

$$F(k) = \frac{g(x)}{ik\phi'(x)} e^{ik\phi(x)} \Big|_a^b - \int_a^b \frac{d}{dx} \left( \frac{g(x)}{ik\phi'(x)} \right) e^{ik\phi(x)} dx$$

Repeating this procedure we will get a power series for  $F(k)$ , so long as  $\phi^n(x) \neq 0$

$$F(k) = a/k + b/k^2 + \dots$$

This is a very simple procedure, but it depends highly on  $\phi^n(x) \neq 0$ . The procedure for the finding the leading order behavior when  $\phi'(x) = 0$  is called stationary phase. We can simplify our problem here by presuming  $\phi'(a) = 0$  and  $\phi'(x) \neq 0$  elsewhere in  $(a, b)$  (if not we will sum over intervals of this form.) Next split the domain into an interval about  $a$  and the rest

$$F(k) = \int_a^{a+\epsilon} g(x)e^{ik\phi(x)} dx + \int_{a+\epsilon}^b g(x)e^{ik\phi(x)} dx$$

By the previous statement the second integral is  $O(1/k)$ . In the first integral we will replace  $g(x)$  with  $g(a)$ , and  $\phi(x)$  with  $\phi(a) + \phi^n(a)(x-a)^n/n!$ ,  $\phi^n(a)$  is the first nonzero derivative of  $\phi$  at  $a$  giving

$$F(k) = \int_a^{a+\epsilon} g(a)e^{ik(\phi(a) + \phi^n(a)(x-a)^n/n!)} dx + O(1/k)$$

Next replace  $a + \epsilon$  with  $\infty$  and introduce only another term proportionate to  $1/k$ . So now we can try to get the leading order behavior from

$$F(k) \sim g(a)e^{ik\phi(a)} \int_0^\infty e^{ik/n!\phi^n(a)s^n} ds$$

Now use the substitution  $u = -ik\frac{\phi^n(a)}{n!}s^n$ . This changes the integral to

$$F(k) \sim g(a)e^{ik\phi(a) \pm i\pi/2} \left( \frac{n!}{k|\phi^n(a)|} \right)^{1/n} \int_0^\infty e^{-u} u^{1/n-1} du$$

All that is left is to recognize that  $\int_0^\infty e^{-s} s^{z-1} = \Gamma(z)$  and we see that in fact we now have

$$F(k) \sim C/k^{1/n}$$

And we can even explicitly write the constant.

# 1 Example

Given a fourier integral representation of a solution below, with  $W(k)$  monotone in  $k$ , find the behavior at long times at a fixed speed, ie where  $x = ct$ .

$$\phi(x, t) = \int_R F(k) e^{ikx - iW(k)t} dk$$

Replace  $x$  with  $ct$  in the integral

$$\phi(x, t) = \int_R F(k) e^{i\chi(k)t} dk$$

Now we know that when  $\chi'(k) \neq 0$  this integral will decay like  $O(1/t)$  (b/c we could integrate by parts). Thus we look at the places where  $\chi'(k) = 0$ . In this problem this is where  $c = W'(k)$  or where  $c$  is the phase speed. Now we make the approximations,  $F = F(K)$ ,  $\chi = \chi(K) + 1/2(k - K)^2 \chi''(K)$ .

$$\phi(x, t) \sim F(K) e^{i\chi(K)t} \int_R e^{i/2(k-K)^2 \chi''(K)t} dk$$

Now we can evaluate this integral either by a change of variables to a gaussian, or to the gamma function. Either way we get

$$\phi \sim F(K) \sqrt{\frac{2\pi}{t|\chi''(K)|}} e^{i\chi(K)t + i\pi/4 \text{sign}(\chi''(K))}$$

Ref: See Bender and Orszag, Whitham  
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