TWO REMARKS ON RECTANGULAR MIXED FINITE ELEMENTS FOR ELASTICITY

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ABSTRACT. The lowest order nonconforming rectangular element in three dimensions involves 54 degrees of freedom for the stress and 12 degrees of freedom for the displacement. With a modest increase in the number of degrees of freedom (24 for the stress), we obtain a conforming rectangular element for linear elasticity in three dimensions. Moreover, unlike the conforming plane rectangular or simplicial elements, this element does not involve any vertex degrees of freedom. Second, we remark that further low order elements can be constructed by approximating the displacement with rigid body motions. This results in a pair of conforming elements with 72 degrees of freedom for the stress and 6 degrees of freedom for the displacement.

1. Introduction

Let $\Omega \subset \mathbb{R}^3$ be a contractible polygonal domain occupied by a linearly elastic body and let $L^2(\Omega, \mathbb{R}^3)$ be the space of square integrable vector fields. We denote by \mathbb{S} the space of 3×3 symmetric matrix fields and by $H(\text{div}, \Omega, \mathbb{S})$ the space of square integrable symmetric matrix fields with divergence, taken row-wise, square integrable. The stress field σ and the displacement u are the unknowns in the elasticity equations

$$A\sigma = \epsilon(u), \quad \text{div } \sigma = f, \quad \text{in } \Omega.$$

Here $\epsilon(u) = (\partial u_i/\partial x_j + \partial u_j/\partial x_i)_{i,j=1,\dots,3}$ is the symmetric part of the gradient and f encodes the body forces. The compliance tensor $A = A(x) : \mathbb{S} \to \mathbb{S}$ is given, bounded and symmetric positive definite uniformly with respect to $x \in \Omega$. The pair (σ, u) is the unique solution of the Hellinger-Reissner variational formulation: Find $(\sigma, u) \in H(\operatorname{div}, \Omega, \mathbb{S}) \times L^2(\Omega, \mathbb{R}^3)$ such that

(1.1)
$$\int_{\Omega} (A\sigma : \tau + \operatorname{div} \tau \cdot u) \, dx = 0, \qquad \tau \in H(\operatorname{div}, \Omega, \mathbb{S}),$$
$$\int_{\Omega} \operatorname{div} \sigma \cdot v \, dx = \int_{\Omega} f \cdot v \, dx, \qquad v \in L^{2}(\Omega, \mathbb{R}^{3}),$$

where $\sigma : \tau = \sum_{i,j=1}^{3} \sigma_{ij} \tau_{ij}$ is the Frobenius product of σ and τ and we assume u = 0 on $\partial \Omega$.

For several decades, the existence of stable, conforming mixed elements for elasticity was an open problem. Starting with the seminal work of Arnold and Winther, [8, 9], for triangular elements, extended to tetrahedral elements in [1, 4, 18], conforming and nonconforming mixed finite elements for elasticity on rectangular meshes have been constructed by several authors [3, 11, 27, 28, 20, 21, 15]. We refer to [8, 11, 7, 12]

and the references therein for alternative approaches to mixed formulations of the elasticity equations. In particular, we note advances on mixed finite elements for elasticity where the symmetry of the stress field is enforced weakly using Lagrange multipliers, [2, 5, 24, 26, 25, 22, 6, 7, 14, 19, 17, 16, 10]. In general, it requires a great deal of ingenuity to construct mixed finite elements for elasticity.

The element described in this paper is the only one known, in three dimension on rectangular meshes, which has strong symmetry for the stress field and is H(div)conforming. The tetrahedral analogues have a high number of degrees of freedom and involve vertex degrees of freedom. In general, on an arbitrary triangulation, vertex degrees of freedom are unavoidable for conforming elements for the stress field with strong symmetry [8, 4]. Although it could be expected that on some triangulation, elements can be constructed with no vertex degrees of freedon [17], it was a surprise to us that this is possible on rectangular meshes in three dimensions which are prefered over tetrahedral elements by some engineers. Progress towards the discovery of the element was made possible by the observation in [15] that for plane elasticity, no vertex degrees of freedom are necessary for the normal normal component of the stress field. Although the resulting element is relatively complicated to use, 72 degrees of freedom for the stress or 12 degrees of freedom on average for each component of the stress field, we believe the strong symmetry property of the element justifies the additional degrees of freedom (the lowest order nonconforming element involves 54 degress of freedom for the stress). An example of a preference in some situations for elements with strong symmetry for two dimensional elasticity is [23] p. 4. We quote "Unfortunately, there exists no simple low order conforming element for ..." $(H(\text{div},\Omega,\mathbb{S}))$. "Possible remedies are to use macro-elements or to weaken the symmetry condition, ..." "However, if we weaken the symmetry, an additional factor due to the nonconformity of the approach occurs". We claim that for the three dimensional problem, the element described in this paper is the only reasonable option available.

The paper is organized as follows: In the next section we describe the finite element spaces on each rectangle after outlining a general framework for the construction of mixed finite elements for elasticity. We show how one may obtain shape functions for three dimensional rectangular elements for elasticity from the corresponding plane elements using tensor products. In Section 3 we prove stability and convergence of the mixed finite elements using arguments now standard. We conclude with some remarks on low order elements using the space of rigid body motions to approximate the displacement and the extension to higher order elements.

2. Description of the finite elements

We denote by $H^k(K,X)$ the space of functions with domain $K \subset \mathbb{R}^3$, taking values in the finite dimensional space X, and with all derivatives of order at most k square integrable. For our purposes, X will be either \mathbb{S} , \mathbb{R}^3 , \mathbb{R}^2 or \mathbb{R} , and in the latter case, we simply write $H^k(X)$. The norms in $H^k(K,X)$ and $H^k(K)$ are denoted respectively by $||\cdot||_{H^k}$ and $||\cdot||_k$. Let \mathcal{T}_h denote a conforming partition of Ω into rectangles Kof size bounded by h, which is quasi-uniform in the sense that the aspect ratio of the rectangles is bounded by a fixed constant. Whether K is a plane rectangle or a three dimensional brick will be clear from the context.

We first outline a set of general criteria to satisfy in order to obtain stable mixed finite elements for the linear elasticity problem. A conforming mixed finite element approximation of (1.1) consists in choosing spaces $\Sigma_h \subset H(\text{div}, \Omega, \mathbb{S})$ and $V_h \subset L^2(\Omega, \mathbb{R}^3)$ and determining $(\sigma_h, u_h) \in \Sigma_h \times V_h$ such that:

(2.1)
$$\int_{\Omega} (A\sigma_h : \tau_h + \operatorname{div} \tau_h \cdot u_h) \, dx = 0, \qquad \tau_h \in \Sigma_h,$$
$$\int_{\Omega} \operatorname{div} \sigma_h \cdot v_h \, dx = \int_{\Omega} f \cdot v_h \, dx, \qquad v_h \in V_h.$$

We recall that a matrix field $\tau \in \Sigma_h$ if and only if τn is continuous across the internal faces of \mathcal{T}_h where n denotes one of the unit vectors normal to the internal face. It is now known that for stability, [8, 9, 11, 4], the following two conditions are sufficient:

- $\operatorname{div} \Sigma_h \subset V_h$
- There exists a linear operator $\Pi_h: H^1(\Omega, \mathbb{S}) \to \Sigma_h$, such that there is a constant c independent of h with $||\Pi_h \tau||_0 \le c||\tau||_1$ for all $\tau \in H^1(\Omega, \mathbb{S})$, and such that $\operatorname{div} \Pi_h \tau = P_h \operatorname{div} \tau$ for all $\tau \in H^1(\Omega, \mathbb{S})$ where $P_h: L^2(\Omega, \mathbb{R}^2) \to V_h$ denotes the L^2 projection.

If we assume that an interpolation operator $\Pi_h: H^1(\Omega, \mathbb{S}) \to \Sigma_h$ can be defined which reproduces the degrees of freedom of $\tau \in \Sigma_h$, the second condition is satisfied if the degrees of freedom for the stress space are chosen such that

$$(2.2) \int_K (\operatorname{div} \Pi_h \tau - \operatorname{div} \tau) \cdot v \, dx = -\int_K (\Pi_h \tau - \tau) : \epsilon(v) \, dx + \int_{\partial K} (\Pi_h \tau - \tau) n \cdot v \, ds,$$

vanishes. Therefore in addition to require normal continuity of the degrees of freedom, it is sufficient that the degrees of freedom include $\int_f \tau n \cdot v \, ds$ and $\int_K \tau : \epsilon(v)$, for each face f of an element K and each $v \in V_K$. If in addition, the interpolation operator reproduces piecewise constant matrix fields, by a standard scaling argument, it is bounded and the error analysis also follows from now classical arguments.

Before describing the new three dimensional conforming elements, we first note that the shape functions for the low order three dimensional nonconforming element in [21], can be obtained from their two dimensional analogue by a tensor product construction.

2.1. **Tensor product construction.** We use the usual notation of $\mathcal{P}_k(K, X)$ for the space of polynomials on K with values in X of total degree less than k and $\mathcal{P}_{k_1,k_2}(K,X)$ for the space of polynomials of degree at most k_1 in x_1 and of degree at most k_2 in x_2 . Similarly, $\mathcal{P}_{k_1,k_2,k_3}(K,X)$ denotes the space of polynomials of degree at most k_1 in x_1 , of degree at most k_2 in x_2 and of degree at most k_3 in x_3 . We write \mathcal{P}_k , \mathcal{P}_{k_1,k_2} and $\mathcal{P}_{k_1,k_2,k_3}$ respectively when $X = \mathbb{R}$. When the domain K in \mathbb{R}^3 is obvious, $\mathcal{P}_k(x_i, x_j)$ and $\mathcal{P}_{k_1,k_2}(x_i, x_j)$ denote respectively the space of polynomials of total degree k in the variables x_i and x_j and of degree at most k_1 in x_i and at most k_2 in x_j for i not necessarily smaller than j. We also use a similar notation for

spaces of polynomials on a face. We will denote by $\begin{pmatrix} \mathcal{P}_{k_1,k_2} & \mathcal{P}_{k_3,k_4} \\ \mathcal{P}_{k_3,k_4} & \mathcal{P}_{k_5,k_6} \end{pmatrix}_{\mathbb{S}}$ the space of symmetric matrix fields $\sigma = \begin{pmatrix} \sigma_{11} & \sigma_{12} \\ \sigma_{21} & \sigma_{22} \end{pmatrix}$ such that $\sigma_{11} \in \mathcal{P}_{k_1,k_2}$, $\sigma_{12} = \sigma_{2,1} \in \mathcal{P}_{k_3,k_4}$ and $\sigma_{22} \in \mathcal{P}_{k_5,k_6}$. Similarly the notation $\begin{pmatrix} \mathcal{P}_{k_1,k_2} \\ \mathcal{P}_{k_3,k_4} \end{pmatrix}$ for $\mathcal{P}_{k_1,k_2} \times \mathcal{P}_{k_3,k_4}$ will also be used. By $\begin{pmatrix} \mathcal{P}_{k_1,k_2,k_3} & \mathcal{P}_{k_4,k_5,k_6} & \mathcal{P}_{k_7,k_8,k_9} \\ \mathcal{P}_{k_4,k_5,k_6} & \mathcal{P}_{k_{10},k_{11},k_{12}} & \mathcal{P}_{k_{13},k_{14},k_{15}} \\ \mathcal{P}_{k_7,k_8,k_9} & \mathcal{P}_{k_{13},k_{14},k_{15}} & \mathcal{P}_{k_{16},k_7,k_{18}} \end{pmatrix}_{\mathbb{S}}$ we mean the space of symmetric matrix fields $\sigma = \begin{pmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{pmatrix}$ such that $\sigma_{11} \in \mathcal{P}_{k_1,k_2,k_3}$, $\sigma_{22} \in \mathcal{P}_{k_7,k_8,k_9}$, $\sigma_{33} \in \mathcal{P}_{k_{13},k_{14},k_{15}}$, $\sigma_{12} = \sigma_{21} \in \mathcal{P}_{k_4,k_5,k_6}$, $\sigma_{13} = \sigma_{31} \in \mathcal{P}_{k_7,k_8,k_9}$ and $\sigma_{23} = \sigma_{32} \in \mathcal{P}_{k_{13},k_{14},k_{15}}$. We will also write $\begin{pmatrix} \mathcal{P}_{k_1,k_2,k_3} \\ \mathcal{P}_{k_4,k_5,k_6} \\ \mathcal{P}_{k_7,k_8,k_9} \end{pmatrix}$ in place of $\mathcal{P}_{k_1,k_2,k_3} \times \mathcal{P}_{k_4,k_5,k_6} \times \mathcal{P}_{k_7,k_8,k_9}$. Often at the place of the spaces $\mathcal{P}_{k_1,m}(K,X)$, we will write explicit vector spaces as span (K, X), we will write explicit vector spaces as span.

For two finite dimensional spaces E and F, we denote by $E \otimes F$ the vector space equal to span $\{ef\}, e \in E, f \in F$. A low order plane nonconforming element described in [20] uses as shape functions on an element K for the stress

$$\sigma \in \Sigma_{\mathrm{ncf}}^{\mathrm{2D}}(K) = \begin{pmatrix} P_{1,1} & \mathrm{span}\{1, x_1, x_2, x_1^2, x_2^2\} \\ \mathrm{span}\{1, x_1, x_2, x_1^2, x_2^2\} & P_{1,1} \end{pmatrix}_{\mathbb{S}},$$

and for the displacement $v \in V^{\text{2D}}(K) = \begin{pmatrix} \operatorname{span}\{1, x_2\} \\ \operatorname{span}\{1, x_1\} \end{pmatrix}$. We postulate that for the three dimensional generalization, the choice of finite dimensional spaces for each component of the stress field $\begin{pmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{12} & \sigma_{22} & \sigma_{23} \\ \sigma_{13} & \sigma_{23} & \sigma_{33} \end{pmatrix}$ and each component of the displacement $\begin{pmatrix} v_1 \\ v_2 \\ v_2 \end{pmatrix}$ should conform to the following requirements:

- $\begin{pmatrix} \sigma_{11} & \sigma_{12} \\ \sigma_{12} & \sigma_{22} \end{pmatrix}$ and $\begin{pmatrix} v_1 \\ v_2 \end{pmatrix}$ belong to the plane elements in the variables x_1, x_2 when restricted to the planes $x_3 = 0, 1$,
- $\begin{pmatrix} \sigma_{11} & \sigma_{13} \\ \sigma_{13} & \sigma_{33} \end{pmatrix}$ and $\begin{pmatrix} v_1 \\ v_3 \end{pmatrix}$ belong to the plane elements in the variables x_1, x_3 when restricted to the planes $x_2 = 0, 1$,
- $\begin{pmatrix} \sigma_{22} & \sigma_{23} \\ \sigma_{23} & \sigma_{33} \end{pmatrix}$ and $\begin{pmatrix} v_2 \\ v_3 \end{pmatrix}$ belong to the plane elements in the variables x_2, x_3 when restricted to the planes $x_1 = 0, 1$.

We note that the above statements can be summarized as follows: for $i < j, i, j \in \{1, 2, 3\}, k \notin \{i, j\}, \begin{pmatrix} \sigma_{ii} & \sigma_{ij} \\ \sigma_{ij} & \sigma_{jj} \end{pmatrix}$ and $\begin{pmatrix} v_i \\ v_j \end{pmatrix}$ belong to the plane elements in the variables x_i, x_j when restricted to the planes $x_k = 0, 1$, that is $\sigma_{ii} \in \mathcal{P}_{1,1}(x_i, x_j), \sigma_{i,j} \in \text{span}\{1, x_i, x_j, x_i^2, x_j^2\}$ and $v_i \in \text{span}\{1, x_j\}, v_j \in \text{span}\{1, x_i\}$ on the planes $x_k = 0, 1$.

A set of shape functions for the three dimensional elasticity problem is obtained by taking as finite element spaces for the components of the stress and the displacement, tensor products of the plane finite element spaces in the variables x_i, x_j with span $\{1, x_k\}$. For this note that for $i < j, i, j \in \{1, 2, 3\}, k \notin \{i, j\}$,

$$\mathcal{P}_{1,1}(x_i, x_j) \otimes \text{span}\{1, x_k\} = \mathcal{P}_{1,1,1},$$

$$\text{span}\{1, x_i, x_j, x_i^2, x_j^2\} \otimes \text{span}\{1, x_k\} = \text{span}\{1, x_i, x_j, x_i^2, x_j^2, x_k,$$

$$x_i x_k, x_j x_k, x_i^2 x_k, x_j^2 x_k)\}$$

$$\text{span}\{1, x_i\} \otimes \text{span}\{1, x_k\}, = \text{span}\{1, x_i, x_k, x_i x_k\} = \mathcal{P}_{1,1}(x_i, x_k).$$

We thus obtain the shape functions for the low order three dimensional element described in [21]. That is if we denote by $\Sigma_{\text{ncf}}^{\text{3D}}(K)$ the finite element for the stress and $V^{\text{3D}}(K)$ the finite element space for the displacement, $\sigma \in \Sigma_{\text{ncf}}^{\text{3D}}(K)$ if and only if $\sigma_{ii} \in \mathcal{P}_{1,1,1}, i = 1, 2, 3, \ \sigma_{i,j} \in \text{span}\{1, x_i, x_j, x_i^2, x_j^2, x_k, \ x_i x_k, x_j x_k, x_i^2 x_k, x_j^2 x_k)\}$ and $v \in V^{\text{3D}}(K)$ if and only if $v_i \in \mathcal{P}_{1,1}(x_j, x_k)$ for $i < j, i, j \in \{1, 2, 3\}, k \notin \{i, j\}$. Alternatively, if we define

$$\tilde{\mathcal{P}}_2(x_i, x_j) = \text{span}\{1, x_i, x_j, x_i^2, x_i^2, x_k, x_i x_k, x_j x_k, x_i^2 x_k, x_i^2 x_k)\},\,$$

then

$$\Sigma_{\text{ncf}}^{\text{3D}}(K) = \begin{pmatrix} \mathcal{P}_{111} & \tilde{\mathcal{P}}_{2}(x_{1}, x_{2}) & \tilde{\mathcal{P}}_{2}(x_{1}, x_{3}) \\ \tilde{\mathcal{P}}_{2}(x_{1}, x_{2}) & \mathcal{P}_{111} & \tilde{\mathcal{P}}_{2}(x_{2}, x_{3}) \\ \tilde{\mathcal{P}}_{2}(x_{1}, x_{3}) & \tilde{\mathcal{P}}_{2}(x_{2}, x_{3}) & \mathcal{P}_{111} \end{pmatrix}_{\mathbb{S}}.$$

2.2. **Description of the elements.** We first describe the shape functions of a conforming rectangular element recently introduced in [15]. Then, following the strategy described in Section 2.1, we give a set of shape functions for the finite element spaces in three dimensions. We give a more concise description of the stress finite element space before giving a unisolvent set of degrees of freedom. Let us denote by $\Sigma_{\rm cf}^{\rm 2D}(K)$ the space of shape functions for the stress for the plane conforming element and by $\Sigma_{\rm cf}^{\rm 3D}(K)$ the corresponding three dimensional element. Then [15],

$$\begin{split} \Sigma_{\mathrm{cf}}^{\mathrm{2D}}(K) &= \Sigma_{\mathrm{ncf}}^{\mathrm{2D}}(K) + \mathrm{span} \bigg\{ \begin{pmatrix} -\frac{1}{2}x_1^2 & x_1x_2 \\ x_1x_2 & -\frac{1}{2}x_2^2 \end{pmatrix}, \begin{pmatrix} -\frac{1}{3}x_1^3 & x_1^2x_2 \\ x_1^2x_2 & -x_1x_2^2 \end{pmatrix}, \begin{pmatrix} -x_1^2x_2 & x_1x_2^2 \\ x_1x_2^2 & -\frac{1}{3}x_2^3 \end{pmatrix}, \\ \begin{pmatrix} -\frac{2}{3}x_1^3x_2 & x_1^2x_2^2 \\ x_1^2x_2^2 & -\frac{2}{3}x_1x_2^3 \end{pmatrix} \bigg\}. \end{split}$$

We choose as finite element space for the displacement the same space $V^{\rm 3D}(K)$ used in the nonconforming case. Note that it can equivalently be written as $\mathcal{P}_{011} \times \mathcal{P}_{101} \times \mathcal{P}_{110}$ and will be denoted by V(K) in the remaining part of this paper. We define $\Sigma_{\rm cf}^{\rm 3D}(K)$ as the direct sum of $\Sigma_{\rm ncf}^{\rm 3D}(K)$ and the span of 24 additional shape functions $\sigma^{ij,(k)}$, i < j, i, j = 1, 2, 3 and $k = 1, \ldots, 8$ obtained from the additional shape functions in the definition of $\Sigma_{\rm cf}^{\rm 2D}(K)$ in the variables x_i and x_j and the tensor product construction

outlined in Section 2.1. For a given matrix, we will refer to the entry on the *i*th row and *j*th column as the (i, j) component.

$$\sigma^{12,(1)} = \begin{pmatrix} -\frac{1}{2}x_1^2 & x_1x_2 & 0 \\ x_1x_2 & -\frac{1}{2}x_2^2 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \ \sigma^{12,(2)} = \begin{pmatrix} -\frac{1}{2}x_1^2x_3 & x_1x_2x_3 & 0 \\ x_1x_2x_3 & -\frac{1}{2}x_2^2x_3 & 0 \\ 0 & 0 & 0 \end{pmatrix},$$

$$\sigma^{12,(3)} = \begin{pmatrix} -\frac{1}{3}x_1^3 & x_1^2x_2 & 0 \\ x_1^2x_2 & -x_1x_2^2 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \ \sigma^{12,(4)} = \begin{pmatrix} -\frac{1}{3}x_1^3x_3 & x_1^2x_2x_3 & 0 \\ x_1^2x_2x_3 & -x_1x_2^2x_3 & 0 \\ 0 & 0 & 0 \end{pmatrix},$$

$$\sigma^{12,(5)} = \begin{pmatrix} -x_1^2 x_2 & x_1 x_2^2 & 0 \\ x_1 x_2^2 & -\frac{1}{3} x_2^3 x_3 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \ \sigma^{12,(6)} = \begin{pmatrix} -x_1^2 x_2 x_3 & x_1 x_2^2 x_3 & 0 \\ x_1 x_2^2 x_3 & -\frac{1}{3} x_2^3 x_3 & 0 \\ 0 & 0 & 0 \end{pmatrix},$$

$$\sigma^{12,(7)} = \begin{pmatrix} -\frac{2}{3} x_1^3 x_2 & x_1^2 x_2^2 & 0 \\ x_1^2 x_2^2 & -\frac{2}{3} x_1 x_2^3 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \ \sigma^{12,(8)} = \begin{pmatrix} -\frac{2}{3} x_1^3 x_2 x_3 & x_1^2 x_2^2 x_3 & 0 \\ x_1^2 x_2^2 x_3 & -\frac{2}{3} x_1 x_2^3 x_3 & 0 \\ 0 & 0 & 0 \end{pmatrix}.$$

Note that the third column of all the above matrices is the zero vector. We conclude that for k = 1, ..., 8, $\sigma^{12,(k)}$ has its (1,3) and (2,3) component 0, but the (1,2) component non-zero.

$$\sigma^{13,(1)} = \begin{pmatrix} -\frac{1}{2}x_1^2 & 0 & x_1x_3 \\ 0 & 0 & 0 \\ x_1x_3 & 0 & -\frac{1}{2}x_3^2 \end{pmatrix}, \sigma^{13,(2)} = \begin{pmatrix} -\frac{1}{2}x_1^2x_2 & 0 & x_1x_2x_3 \\ 0 & 0 & 0 \\ x_1x_2x_3 & 0 & -\frac{1}{2}x_2x_3^2 \end{pmatrix},$$

$$\sigma^{13,(3)} = \begin{pmatrix} -\frac{1}{3}x_1^3 & 0 & x_1^2x_3 \\ 0 & 0 & 0 \\ x_1^2x_3 & 0 & -x_1x_3^2 \end{pmatrix}, \sigma^{13,(4)} = \begin{pmatrix} -\frac{1}{3}x_1^3x_2 & 0 & x_1^2x_2x_3 \\ 0 & 0 & 0 \\ x_1^2x_2x_3 & 0 & x_1x_2x_3^2 \end{pmatrix},$$

$$\sigma^{13,(5)} = \begin{pmatrix} -x_1^2x_3 & 0 & x_1x_3^2 \\ 0 & 0 & 0 \\ x_1x_3^2 & 0 & -\frac{1}{3}x_3^3 \end{pmatrix}, \sigma^{13,(6)} = \begin{pmatrix} -x_1^2x_2x_3 & 0 & x_1x_2x_3^2 \\ 0 & 0 & 0 \\ x_1x_2x_3^2 & 0 & -\frac{1}{3}x_2x_3^3 \end{pmatrix},$$

$$\sigma^{13,(7)} = \begin{pmatrix} -\frac{2}{3}x_1^3x_3 & 0 & x_1^2x_2^2 \\ 0 & 0 & 0 \\ x_1^2x_2^2 & 0 & -\frac{2}{3}x_1x_3^2 \end{pmatrix}, \sigma^{13,(8)} = \begin{pmatrix} -\frac{2}{3}x_1^3x_2x_3 & 0 & x_1^2x_2x_3^2 \\ 0 & 0 & 0 \\ x_1^2x_2x_3^2 & 0 & -\frac{2}{3}x_1x_2x_3^3 \end{pmatrix}.$$

Note that the second row or second column of all the matrices $\sigma^{13,(k)}$, k = 1, ..., 8 is the zero vector. We conclude that for k = 1, ..., 8, $\sigma^{13,(k)}$ has its (1,2) and (2,3) component 0, but the (1,3) component non-zero.

$$\sigma^{23,(1)} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & -\frac{1}{2}x_2^2 & x_2x_3 \\ 0 & x_2x_3 & -\frac{1}{2}x_3^2 \end{pmatrix}, \ \sigma^{23,(1)} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & -\frac{1}{2}x_1x_2^2 & x_1x_2x_3 \\ 0 & x_1x_2x_3 & -\frac{1}{2}x_1x_3^2 \end{pmatrix},$$

$$\sigma^{23,(3)} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & -\frac{1}{3}x_2^3 & x_2^2x_3 \\ 0 & x_2^2x_3 & -x_2x_3^2 \end{pmatrix}, \ \sigma^{23,(4)} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & -\frac{1}{3}x_1x_2^3 & x_1x_2^2x_3 \\ 0 & x_1x_2^2x_3 & -x_1x_2x_3^2 \end{pmatrix},$$

$$\sigma^{23,(5)} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & -x_2^2 x_3 & x_2 x_3^2 \\ 0 & x_2 x_3^2 & -\frac{1}{3} x_3^3 \end{pmatrix}, \, \sigma^{23,(6)} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & -x_1 x_2^2 x_3 & x_1 x_2 x_3^2 \\ 0 & x_1 x_2 x_3^2 & -\frac{1}{3} x_1 x_3^3 \end{pmatrix},$$

$$\sigma^{23,(7)} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & -\frac{2}{3} x_2^3 x_3 & x_2^2 x_3^2 \\ 0 & x_2^2 x_3^2 & -\frac{2}{3} x_2 x_3^3 \end{pmatrix}, \, \sigma^{23,(8)} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & -\frac{2}{3} x_1 x_2^3 x_3 & x_1 x_2^2 x_3^2 \\ 0 & x_1 x_2^2 x_3^2 & -\frac{2}{3} x_1 x_2 x_3^3 \end{pmatrix}.$$

Note that the first row or first column of all the matrices $\sigma^{23,(k)}$, $k=1,\ldots,8$ is the zero vector. We conclude that for $k = 1, \ldots, 8, \sigma^{23,(k)}$ has its (1,2) and (1,3) component 0, but the (1,3) component non-zero.

We can summarize the above observations by stating that

(2.3)
$$\sigma_{pq}^{ij,(k)}, k = 1, \dots, 8 = 0 \text{ for } i < j \text{ unless } (p,q) = (i,j).$$

Note that (2.3) can also be seen directly from the tensor product construction. We also note that like their two dimensional analogue, the above additional shape functions are divergence free and that the dimension of $\Sigma_{\rm cf}^{\rm 3D}(K)$ is 54+24=78. This is implicit in the tensor product construction and also follows from Lemma 2.1 below. We can describe $\Sigma_{\mathrm{cf}}^{\mathrm{3D}}(K)$ in a more concise form as follows: $\tau \in \Sigma_{\mathrm{cf}}^{\mathrm{3D}}(K)$ if and only if

(2.4)
$$\tau \in \Sigma(K) := \left\{ \begin{pmatrix} \mathcal{P}_{311} & \mathcal{P}_{221} & \mathcal{P}_{212} \\ \mathcal{P}_{221} & \mathcal{P}_{131} & \mathcal{P}_{122} \\ \mathcal{P}_{212} & \mathcal{P}_{122} & \mathcal{P}_{113} \end{pmatrix}_{\mathbb{S}}, \operatorname{div} \tau \in \begin{pmatrix} \mathcal{P}_{011} \\ \mathcal{P}_{101} \\ \mathcal{P}_{110} \end{pmatrix} \right\}.$$

Clearly $\Sigma_{\rm cf}^{\rm 3D}(K)$ is contained in the latter space. Notice that the dimension of $S_2 :=$

Clearly
$$\mathcal{L}_{cf}^{-}(K)$$
 is contained in the latter space. Notice that the dimension of $S_2 := \begin{pmatrix} \mathcal{P}_{311} & \mathcal{P}_{221} & \mathcal{P}_{212} \\ \mathcal{P}_{221} & \mathcal{P}_{131} & \mathcal{P}_{122} \\ \mathcal{P}_{212} & \mathcal{P}_{122} & \mathcal{P}_{113} \end{pmatrix}$ is 102. And for $\tau \in S_2$, div $\tau \in \begin{pmatrix} \mathcal{P}_{211} \\ \mathcal{P}_{121} \\ \mathcal{P}_{112} \end{pmatrix}$ which has dimension 36.

Hence the constraint div $\tau \in V(K)$ imposes 36-12=24 constraints. It follows that the dimension of $\Sigma(K)$ is at least 78. That its dimension is 78 follows from the unisolvency of the degrees of freedom we give in Lemma 2.1 below. Hence, $\Sigma(K) = \Sigma_{\rm cf}^{\rm 3D}(K)$.

Before we give the degrees of freedom of $\Sigma_{\rm cf}^{\rm 3D}(K)$, we explain why unlike the conforming plane rectangular or simplicial elements, this element does not involve any vertex degrees of freedom. In general, on an arbitrary simplicial triangulation, vertex degrees of freedom are necessary [8, 4]. It will be enough to consider the unit cube $[0,1]^3$. We first write explicitly the condition of normal continuity of

$$\sigma = \begin{pmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{12} & \sigma_{22} & \sigma_{23} \\ \sigma_{13} & \sigma_{23} & \sigma_{33} \end{pmatrix}.$$

Clearly on horizontal faces $x_3 = 0, 1, \sigma n = \pm(\sigma_{13}, \sigma_{23}, \sigma_{33})$. On vertical faces $x_1 = 0, 1, \sigma_{13} = 0$ $\sigma n = \pm(\sigma_{11}, \sigma_{12}, \sigma_{13})$ and on lateral faces $x_2 = 0, 1$ $\sigma n = \pm(\sigma_{12}, \sigma_{22}, \sigma_{23})$. We therefore need σ_{33} continuous on faces $x_3 = 0, 1, \sigma_{22}$ continuous on faces $x_2 = 0, 1$ and σ_{11} continuous on faces $x_1 = 0, 1$. Since for i = 1, 2, 3 the faces $x_i = 0$ and $x_i = 1$ do not meet, no edge or vertex degrees of freedom are necessary for σ_{ii} . Next, we need σ_{23} continuous on faces $x_3 = 0, 1, x_2 = 0, 1$, so we need σ_{23} continuous on the edges $x_2 = 0, x_3 = 0; x_2 = 1, x_3 = 0; x_2 = 0, x_3 = 1; x_2 = 1, x_3 = 1.$ And similar conditions for σ_{13} and σ_{12} , that is for $i < j, i, j \in \{1, 2, 3\}, k \notin \{i, j\}$ we need σ_{ij} continuous on the faces $x_i = 0, 1$ and $x_j = 0, 1$. This requires the continuity of σ_{ij} on the four edges $x_i, x_i \in \{0, 1\}$. No vertex degrees of freedom are therefore necessary here.

Let $\hat{K} = [0,1]^3$ be the reference cube and let $F: \hat{K} \to K$ be an affine mapping onto $K, F(\hat{x}) = B\hat{x} + b, \text{ with } b = (b_1, b_2, b_3)^T \text{ and }$

$$B = \begin{pmatrix} h_1 & 0 & 0 \\ 0 & h_2 & 0 \\ 0 & 0 & h_3 \end{pmatrix}.$$

To maintain the analogy with the simplicial case, although B is symmetric we will write B^T . Given a matrix field $\hat{\tau}: \hat{K} \to \mathbb{S}$, define $\tau: K \to \mathbb{S}$ by the matrix Piola transform

$$\tau(x) = B\hat{\tau}(\hat{x})B^T.$$

It is not difficult to verify that $\tau \in \Sigma(K)$ if and only if $\hat{\tau} \in \Sigma(\hat{K})$. We recall that an edge of K is given by $x_i, x_j \in \{0, 1\}$ and we denote by $\int_{\{x_i=0,1, x_j=0,1\}}$ an integral over such an edge. We now give the degrees of freedom on $\Sigma(\hat{K})$.

Lemma 2.1. An element $\hat{\sigma} \in \Sigma(\hat{K})$ is uniquely determined by the following degrees of freedom

- (1) $\int_{\{x_{i}=0,1,\ x_{j}=0,1\}} \hat{\sigma}_{ij} \hat{v} \ dx_{f}, \ \hat{v} \in span\{1,x_{k},\} \ for \ k \notin \{i,j\}, \ i < j \ i,j = 1,2,3,$ (24 degrees of freedom),
- (2) $\int_{x_{i}=0,1}^{\cdot} \hat{\sigma}_{ii}\hat{v} dx, \hat{v} \in \mathcal{P}_{1,1}(x_{l}, x_{k}), i \notin \{l, k\}, i = 1, 2, 3, (24 \text{ degrees of freedom}),$ (3) $\int_{x_{l}=0,1}^{\cdot} \hat{\sigma}_{ij}\hat{v} dx_{f}, \hat{v} \in span\{1, x_{k}\}, k \notin \{i, j\}, i < j, l = i, j, (24 \text{ degrees of freedom})$
- (4) $\int_{K} \hat{\sigma}_{ij} \hat{v} \ dx$, $\hat{v} \in span\{1, x_k\}$, $k \notin \{i, j\}$, i < j, i, j = 1, 2, 3, (6 degrees of freedom).

Proof. We assume that all degrees of freedom vanish. Note that on the faces $x_i = 0, 1$, $\hat{\sigma}_{ij} \in \mathcal{P}_{2,1}(x_j, x_k), \ k \notin \{i, j\}, \ i < j, i, j = 1, 2, 3, \ (\text{e.g. on the faces} \ x_2 = 0, 1, \hat{\sigma}_{23} \in \mathcal{P}_{2,1}(x_j, x_k), \ k \notin \{i, j\}, \ i < j, i, j = 1, 2, 3, \ (i, j) \in \mathcal{P}_{2,1}(x_j, x_k), \ k \notin \{i, j\}, \ i < j, i, j = 1, 2, 3, \ (i, j) \in \mathcal{P}_{2,1}(x_j, x_k), \ k \notin \{i, j\}, \ i < j, i, j = 1, 2, 3, \ (i, j) \in \mathcal{P}_{2,1}(x_j, x_k), \ k \notin \{i, j\}, \ k \in \{i, j\}, \ k$ $\mathcal{P}_{2,1}(x_3, x_1)$) and on an edge $x_i, x_j \in \{0, 1\}$, $\hat{\sigma}_{ij} \in \mathcal{P}_1(x_k)$. By the first set of degrees of freedom, we have $\hat{\sigma}_{ij} = 0$ on such an edge. Thus on a face $x_i = 0, 1, \hat{\sigma}_{ij} = x_j(1-x_j)\tilde{\sigma}_{ij}$, $\tilde{\sigma}_{ij} \in \mathcal{P}_{0,1}(x_j, x_k)$. Since $\int_{x_i=0,1} \hat{\sigma}_{ij} \hat{v} \, dx_f = 0$, $\hat{v} \in \text{span}\{1, x_k\}$ by the third set of degrees of freedom, $\hat{\sigma}_{ij} = 0$ on the faces $x_i = 0, 1$. Similarly, on the faces $x_j = 0, 1$, $\hat{\sigma}_{ij} \in \mathcal{P}_{2,1}(x_i, x_k), \ k \notin \{i, j\}, \ i < j, i, j = 1, 2, 3 \text{ and by the same argument, } \hat{\sigma}_{ij} = 0 \text{ as}$ well on the faces $x_i = 0, 1$.

We therefore have $\hat{\sigma}_{ij} = x_j(1-x_j)x_i(1-x_i)\tilde{\sigma}_{ij}$. Since $\hat{\sigma}_{ij} \in \mathcal{P}_{2,2,1}(x_i,x_j,x_k)$, we have $\tilde{\sigma}_{ij} \in \mathcal{P}_{0,0,1}$.

With the last set of degrees of freedom, we conclude that $\hat{\sigma}_{ij} = 0, i < j, i, j = 1, 2, 3$.

Let us write

$$\hat{\sigma} = \tilde{\sigma} + \sum_{k=1}^{8} \alpha^{12,(k)} \sigma^{12,(k)} + \alpha^{13,(k)} \sigma^{13,(k)} + \alpha^{23,(k)} \sigma^{23,(k)},$$

$$\alpha^{ij,(k)} \in \mathbb{R}, i < j, i, j = 1, 2, 3 \text{ and } k = 1, \dots, 8, \tilde{\sigma} \in \Sigma_{\text{cf}}^{3D}(K),$$

and recall that $\hat{\sigma}_{12} = \hat{\sigma}_{13} = \hat{\sigma}_{23} = 0$. We show that $\alpha^{12,(k)} = \alpha^{13,(k)} = \alpha^{23,(k)} = 0$, $k = 1, \ldots, 8$. This is a consequence of the tensor product construction as the following argument shows. By (2.3), for i < j

$$\hat{\sigma}_{ij} = \tilde{\sigma}_{ij} + \sum_{k=1}^{8} \alpha^{ij,(k)} \sigma_{ij}^{ij,(k)}.$$

Explicitly

$$\begin{split} \hat{\sigma}_{12} &= \tilde{\sigma}_{12} + \alpha^{12,(1)} x_1 x_2 + \alpha^{12,(2)} x_1 x_2 x_3 + \alpha^{12,(3)} x_1^2 x_2 + \alpha^{12,(4)} x_1^2 x_2 x_3 + \alpha^{12,(5)} x_1 x_2^2 + \alpha^{12,(6)} x_1 x_2^2 x_3 \\ &\quad + \alpha^{12,(7)} x_1^2 x_2^2 + \alpha^{12,(8)} x_1^2 x_2^2 x_3, \\ \hat{\sigma}_{13} &= \tilde{\sigma}_{13} + \alpha^{13,(1)} x_1 x_3 + \alpha^{13,(2)} x_1 x_2 x_3 + \alpha^{13,(3)} x_1^2 x_3 + \alpha^{13,(4)} x_1^2 x_2 x_3 + \alpha^{13,(5)} x_1 x_3^2 + \alpha^{13,(6)} x_1 x_2 x_3^2 \\ &\quad + \alpha^{13,(7)} x_1^2 x_3^2 + \alpha^{13,(8)} x_1^2 x_2 x_3^2, \\ \hat{\sigma}_{23} &= \tilde{\sigma}_{23} + \alpha^{23,(1)} x_2 x_3 + \alpha^{23,(2)} x_1 x_2 x_3 + \alpha^{23,(3)} x_2^2 x_3 + \alpha^{23,(4)} x_1 x_2^2 x_3 + \alpha^{23,(5)} x_2 x_3^2 + \alpha^{23,(6)} x_1 x_2 x_3^2 \\ &\quad + \alpha^{23,(7)} x_2^2 x_3^2 + \alpha^{23,(8)} x_1 x_2^2 x_3^2, \end{split}$$

and recall that by definition

$$\begin{split} &\tilde{\sigma}_{12} \in \text{ span}\{1, x_1, x_2, x_1^2, x_2^2, x_3, x_1x_3, x_2x_3, x_1^2x_3, x_2^2x_3\}, \\ &\tilde{\sigma}_{13} \in \text{ span}\{1, x_1, x_3, x_1^2, x_3^2, x_2, x_1x_2, x_2x_3, x_1^2x_2, x_2x_3^2\}, \\ &\tilde{\sigma}_{23} \in \text{ span}\{1, x_2, x_3, x_2^2, x_3^2, x_1, x_1x_2, x_1x_3, x_1x_2^2, x_1x_3^2\}, \end{split}$$

and the result follows since each of the monomial appearing in the expansion of $\hat{\sigma}_{ij}$, i < j, appears only one time and for each expansion, they are all different.

We conclude that $\hat{\sigma} \in \Sigma_{\text{ncf}}^{3D}(K)$. By the second set of degrees of freedom, since $\hat{\sigma}_{ii} \in \mathcal{P}_{111}$, $\hat{\sigma}_{ii} = 0$ on each face $x_i = 0, 1$. Therefore, $\hat{\sigma}_{ii} = x_i(1 - x_i)\tilde{\sigma}_{ii}$ for some polynomial $\tilde{\sigma}_{ii}$ so $\hat{\sigma}_{ii} = 0$ for all i = 1, 2, 3.

If we denote by n_l , l = 1, 2, 3 the normal in the lth direction, each face is given by $x_i = 0, 1$ for some $i \in \{1, 2, 3\}$, an edge is the intersection of two faces $x_i = 0, 1, x_j = 0, 1$ for some i, j. We can then rewrite the degrees of freedom as follows

- (1') $\int_{\hat{e}} \hat{\sigma} n_i \cdot n_j \, \hat{v} \, dx$, $\hat{v} \in \mathcal{P}_1(\hat{e})$, for each edge with normals n_i and n_j , $i \neq j, i, j = 1, 2, 3$.
- (2') $\int_{\hat{f}} \hat{\sigma} n \cdot n \, \hat{v} \, dx, \, \hat{v} \in \mathcal{P}_{1,1}(\hat{f}),$
- (3') $\int_{\hat{f}} \hat{\sigma} n_i \cdot n_l \, \hat{v} \, dx$, $\hat{v} \in \text{span}\{1, x_k\}$, $k \notin \{i, l\}$, for each face \hat{f} , with normal n_i and $i \neq l$ and all l = 1, 2, 3.

Note that for
$$\hat{v} \in V(\hat{K}), \epsilon(\hat{v}) \in \begin{pmatrix} 0 & \text{span}\{1, x_3\} & \text{span}\{1, x_2\} \\ \text{span}\{1, x_3\} & 0 & \text{span}\{1, x_1\} \\ \text{span}\{1, x_2\} & \text{span}\{1, x_1\} & 0 \end{pmatrix}_{\mathbb{S}}$$
, where

the zero vector space is simply denoted 0. We can thus write the fourth set of degrees of freedom as

(4')
$$\int_{\hat{K}} \hat{\sigma} : \epsilon(\hat{v}), \ \hat{v} \in V(\hat{K}).$$

Given the discussion before the above lemma, the degrees of freedom in Lemma 2.1 imply that for $\hat{\sigma} \in \Sigma(\hat{K})$, $\hat{\sigma}n$ is uniquely determined by the degrees of freedom on the face with normal n.

On the rectangular partition \mathcal{T}_h , we define Σ_h as the space of matrix fields which belong piecewise to $\Sigma(K)$ subject to the continuity conditions that τn for $\tau \in \Sigma_h$ is continuous across faces. We denote by V_h the space of vector fields which belongs to V_K for each rectangle $K \in \mathcal{T}_h$.

3. Stability and Error analysis

Our finite element spaces satisfy $\Sigma_h \subset H(\operatorname{div}, \mathbb{S})$ and $V_h \subset L^2(\Omega, \mathbb{R})$ with $\operatorname{div} \Sigma_h \subset V_h$. We next define an interpolation operator $\Pi_h : H^1(\Omega, \mathbb{S}) \to \Sigma_h$ with the required properties. Because of the edge degrees of freedom, the canonical interpolation operator for Σ_h is not well-defined on $H^1(\Omega, \mathbb{S})$. As in [3], we first define an interpolation operator $\Pi_{\hat{K}}^0 : H^1(\hat{K}, \mathbb{S}) \to \Sigma(\hat{K})$ by

$$\begin{split} &\int_{\hat{e}} (\Pi_{\hat{K}}^{0}(\hat{\sigma})) n_{i} \cdot n_{j} \, \hat{v} \, dx = 0 \quad \forall \hat{v} \in \mathcal{P}_{1}(\hat{e}) \text{ for } \hat{e} \text{ with normals } n_{i} \text{ and } n_{j}, \, i \neq j, i, j = 1, 2, 3, \\ &\int_{\hat{f}} (\hat{\sigma} - \Pi_{\hat{K}}^{0}(\hat{\sigma})) n \cdot n \, \hat{v} \, dx = 0 \quad \forall \hat{v} \in \mathcal{P}_{1,1}(\hat{f}), \\ &\int_{\hat{f}} (\hat{\sigma} - \Pi_{\hat{K}}^{0}(\hat{\sigma})) n_{i} \cdot n_{l} \, dx = 0 \quad \text{for each face } \hat{f}, \text{ with normal } n_{i} \text{ and } i \neq l \text{ and all } l = 1, 2, 3, \\ &\int_{\hat{K}} (\hat{\sigma} - \Pi_{\hat{K}}^{0}(\hat{\sigma})) : \epsilon(\hat{v}) = 0 \quad \forall \hat{v} \in V(\hat{K}). \end{split}$$

By Lemma 2.1, $\Pi^0_{\hat{K}}$ is well defined. Next define $\Pi^0_K: H^1(K,\mathbb{S}) \to \Sigma(K)$ by

$$\Pi_K^0 \tau(x) = B \Pi_{\hat{K}}^0 \hat{\tau}(\hat{x}) B^T,$$

for each rectangle K of \mathcal{T}_h and define $\Pi_h^0: H^1(\Omega, \mathbb{S}) \to \Sigma_h$ by

$$(\Pi_h^0 \tau)|_K = \Pi_K \tau.$$

A standard scaling argument gives

where c does not depend on h since Π_h^0 reproduces piecewise constant matrix fields. It follows that Π_h^0 is bounded on $H^1(\Omega, \mathbb{S})$.

The proof of the commutativity property $\operatorname{div} \Pi_h^0 \tau = P_h \operatorname{div} \tau$ for all $\tau \in H^1(\Omega, \mathbb{S})$ is similar to the one given in [3], pp. 1423–1424. Next, let R_h be a Clement interpolation

operator [13] which maps $L^2(\Omega, \mathbb{S})$ into the subspace of Σ_h of continuous matrix fields whose components are piecewise in $P_{1,1,1}$. We have

$$||R_h \tau - \tau||_j \le ch^{m-j} ||\tau||_m, \quad 0 \le j \le 1, \quad j \le m \le 2,$$

with c independent of h. The interpolation operator Π_h is defined by

(3.2)
$$\Pi_h = \Pi_h^0 (I - R_h) + R_h,$$

and as in [3], one establishes its boundedness and the commutativity property. We conclude that the pair (Σ_h, V_h) is stable.

For the error analysis, we recall that the projection operator P_h satisfies the error estimate

$$(3.3) ||P_h v - v||_0 \le ch^m ||v||_m, \text{ for all } v \in H^m(\Omega), \quad 0 \le m \le 1,$$

since V_h contains piecewise constant vector fields. We then have the following theorem

Theorem 3.1. Let (σ, u) and (σ_h, u_h) be the unique solutions of (1.1) and (2.1) respectively. Then

$$\begin{aligned} ||\operatorname{div} \sigma - \operatorname{div} \sigma_h||_0 &\leq ch^m ||\operatorname{div} \sigma||_m, \quad 0 \leq m \leq 1, \\ ||\sigma - \sigma_h||_0 &\leq ch ||\sigma||_1, \\ ||u - u_h||_0 &\leq ch ||u||_2. \end{aligned}$$

Proof. We only outline the proof since it is essentially the same as the one in [8]. Define $||.||_A$ by $||\tau||_A^2 = \int_{\Omega} A\tau : \tau \, dx$. Then one shows that

$$||\sigma - \sigma_h||_A \le ||\sigma - \Pi_h \sigma||_A.$$

Since the norm $||.||_A$ is equivalent to the L^2 norm, we obtain using (3.1),

$$||\sigma - \sigma_h||_0 \le ch||\sigma||_1.$$

One also shows that $\operatorname{div} \sigma_h = P_h \operatorname{div} \sigma$. The commutativity property and (3.3) yields the error estimate on $\operatorname{div} \sigma$. For the error estimate for the displacement, again from [8], we have

$$||P_h u - u_h||_0 \le ||\sigma - \sigma_h||_0,$$

so

$$||u - u_h||_0 \le ||u - P_h u||_0 + ||P_h u - u_h||_0 \le ch||u||_1 + ch||\sigma||_1 \le ch||u||_2,$$

since $A\sigma = \epsilon(u)$.

4. Concluding remarks

In three dimensions, the space of rigid body motions \mathbb{T} is the space of linear polynomial functions of the form $x \mapsto a + b \times x$ for some $a, b \in \mathbb{R}^3$. It is contained in V(K) for each element K. Hence, as in [8, 9, 3, 4, 11], we can take $V(K) = \mathbb{T}$ and the stress space as

(4.1)
$$\tau \in \Sigma(K) := \left\{ \begin{pmatrix} \mathcal{P}_{311} & \mathcal{P}_{221} & \mathcal{P}_{212} \\ \mathcal{P}_{221} & \mathcal{P}_{131} & \mathcal{P}_{122} \\ \mathcal{P}_{212} & \mathcal{P}_{122} & \mathcal{P}_{113} \end{pmatrix}_{\mathbb{S}}, \operatorname{div} \tau \in \mathbb{T} \right\}.$$

Since for $v \in \mathbb{T}$, $\epsilon(v) = 0$, the degrees of freedom are the same as the ones in Lemma 2.1 with the last set of degrees of freedom removed. The error analysis is the same. We note that the elements of [20, 21, 15] can be simplified as well this way.

As noted for example in [12], high order mixed finite elements for elasticity usually require a high number of degrees of freedom and typically require to find the dimension and a base for the space

$$\{\tau \in \Sigma(K), \operatorname{div} \sigma = 0, \sigma n = 0, \operatorname{on} \partial \Omega\}.$$

A possible approach is to follow the proof in the simplicial case [4], but this goes well beyond the scope of these remarks. Thus for that reason as well, we did not pursue here the construction of a discrete elasticity sequence.

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References

- 1. Scot Adams and Bernardo Cockburn, A mixed finite element method for elasticity in three dimensions, J. Sci. Comput. **25** (2005), no. 3, 515–521.
- 2. M. Amara and J. M. Thomas, Equilibrium finite elements for the linear elastic problem, Numer. Math. **33** (1979), no. 4, 367–383.
- 3. Douglas N. Arnold and Gerard Awanou, Rectangular mixed finite elements for elasticity, Math. Models Methods Appl. Sci. 15 (2005), no. 9, 1417–1429.
- 4. Douglas N. Arnold, Gerard Awanou, and Ragnar Winther, Finite elements for symmetric tensors in three dimensions, Math. Comp. 77 (2008), no. 263, 1229–1251.
- Douglas N. Arnold, Franco Brezzi, and Jr Douglas, J., PEERS: a new mixed finite element for plane elasticity, Japan J. Appl. Math. 1 (1984), no. 2, 347–367.
- 6. Douglas N. Arnold, Richard S. Falk, and Ragnar Winther, Finite element exterior calculus, homological techniques, and applications, Acta Numer. 15 (2006), 1–155.
- 7. ______, Mixed finite element methods for linear elasticity with weakly imposed symmetry, Math. Comp. **76** (2007), no. 260, 1699–1723 (electronic).
- 8. Douglas N. Arnold and Ragnar Winther, *Mixed finite elements for elasticity*, Numer. Math. **92** (2002), no. 3, 401–419.
- 9. _____, Nonconforming mixed elements for elasticity, Math. Models Methods Appl. Sci. 13 (2003), no. 3, 295–307, Dedicated to Jim Douglas, Jr. on the occasion of his 75th birthday.
- 10. Gerard Awanou, Rectangular mixed elements for elasticity with weakly imposed symmetry condition, Preprint, 2010.
- 11. _____, A rotated nonconforming rectangular mixed element for elasticity, Calcolo **46** (2009), no. 1, 49–60.
- 12. Gerard Awanou, Symmetric matrix fields in the finite element method, Symmetry 2 (2010), 1375–1389.
- 13. C. Bernardi and V. Girault, A local regularization operator for triangular and quadrilateral finite elements, SIAM J. Numer. Anal. **35** (1998), no. 5, 1893–1916 (electronic).
- 14. Daniele Boffi, Franco Brezzi, Leszek F. Demkowicz, Ricardo G. Durán, Richard S. Falk, and Michel Fortin, *Mixed finite elements, compatibility conditions, and applications*, Lecture Notes in Mathematics, vol. 1939, Springer-Verlag, Berlin, 2008, Lectures given at the C.I.M.E. Summer School held in Cetraro, June 26–July 1, 2006, Edited by Boffi and Lucia Gastaldi.
- 15. Shao-Chun Chen and Ya-Na Yang, Conforming rectangular mixed finite elements for elasticity, J. of Scientific Computing (2010), Submitted.

- 16. Bernardo Cockburn, Jayadeep Gopalakrishnan, and Johnny Guzmán, A new elasticity element made for enforcing weak stress symmetry, Math. Comp. **79** (2010), no. 271, 1331–1349.
- J. Gopalakrishnan and J. Guzmán, A second elasticity element using the matrix bubble, Submitted, 2010.
- 18. _____, Symmetric non-conforming mixed finite elements for linear elasticity, Submitted.
- 19. J. Guzmán, A unified analysis of several mixed methods for elasticity with weak stress symmetry, J. Sci. Comput. 44 (2010), no. 2, 156–169.
- 20. Jun Hu and Zhong-Ci Shi, Lower order rectangular nonconforming mixed finite elements for plane elasticity, SIAM J. Numer. Anal. 46 (2007/08), no. 1, 88–102.
- 21. Hong-Ying Man, Jun Hu, and Zhong-Ci Shi, Lower order rectangular nonconforming mixed finite element for the three-dimensional elasticity problem, Math. Models Methods Appl. Sci. 19 (2009), no. 1, 51–65.
- 22. Mary E. Morley, A family of mixed finite elements for linear elasticity, Numer. Math. **55** (1989), no. 6, 633–666.
- Serge Nicaise, Katharina Witowski, and Barbara I. Wohlmuth, An a posteriori error estimator for the Lamé equation based on equilibrated fluxes, IMA J. Numer. Anal. 28 (2008), no. 2, 331–353.
- 24. R. Stenberg, On the construction of optimal mixed finite element methods for the linear elasticity problem, Numer. Math. 48 (1986), no. 4, 447–462.
- 25. _____, A family of mixed finite elements for the elasticity problem, Numer. Math. 53 (1988), no. 5, 513–538.
- 26. _____, Two low-order mixed methods for the elasticity problem, The mathematics of finite elements and applications, VI (Uxbridge, 1987), Academic Press, London, 1988, pp. 271–280.
- 27. S.-Y. Yi, Nonconforming mixed finite element methods for linear elasticity using rectangular elements in two and three dimensions, Calcolo 42 (2005), no. 2, 115–133.
- 28. Son-Young Yi, A new nonconforming mixed finite element method for linear elasticity, Math. Models Methods Appl. Sci. 16 (2006), no. 7, 979–999.

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