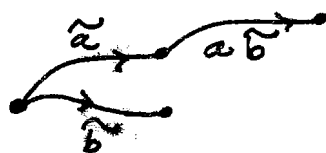


Notes on Alexander Module and Alexander Polynomial
by LK ①

$$\widetilde{ab} = \widetilde{a} + a\widetilde{b}$$



Basic rule for lifting path into covering space.

If start with $\mathbb{G} = (\kappa_1, \dots, \kappa_n |)$ a free group then use base space

$$B = \bigvee_{i=1}^n S^1 = \text{a wedge of circles}$$

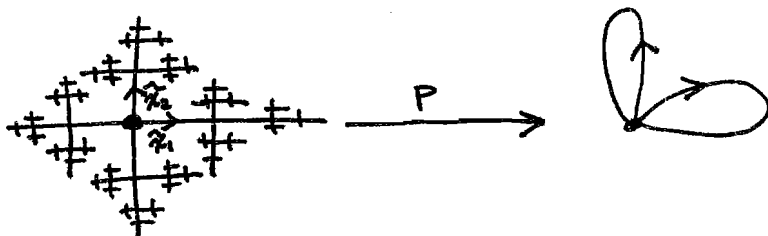
circles cover to give $\kappa_1, \dots, \kappa_n$.

$p: E \rightarrow B$ univ covering space.

E is an infinite tree, acted on freely by \mathbb{G} with fundamental 1-cells

$$\widetilde{\kappa}_1, \dots, \widetilde{\kappa}_n.$$

$n=2$:



e.g.

$$\begin{aligned} \widetilde{\kappa_1 \kappa_2 \kappa_1 \kappa_2} &= \widetilde{\kappa_1} + \kappa_1 (\widetilde{\kappa_2 \kappa_1 \kappa_2}) \\ &= \widetilde{\kappa_1} + \kappa_1 (\widetilde{\kappa_2} + \kappa_2 (\widetilde{\kappa_1} + \kappa_1 \widetilde{\kappa_2})) \\ &= \widetilde{\kappa_1} + \kappa_1 \widetilde{\kappa_2} + \kappa_1 \kappa_2 \widetilde{\kappa_1} + \kappa_1 \kappa_2 \kappa_1 \widetilde{\kappa_2} \\ &= (1 + \kappa_1 \kappa_2) \widetilde{\kappa_1} + (\kappa_1 + \kappa_1 \kappa_2 \kappa_1) \widetilde{\kappa_2} \end{aligned}$$

Defn. $\mathbb{G} = (\kappa_1, \dots, \kappa_n |)$, $\mathbb{Z}[\mathbb{G}] =$ integral group ring of \mathbb{G} . $w \in \mathbb{G}$ then

$$\widetilde{w} = \frac{\partial w}{\partial \kappa_1} \widetilde{\kappa}_1 + \dots + \frac{\partial w}{\partial \kappa_n} \widetilde{\kappa}_n$$

defines operators

$$D_i = \frac{\partial}{\partial \kappa_i} : \mathbb{G} \rightarrow \mathbb{Z}[\mathbb{G}].$$

If $D = D_i$ for some i then

$$D(ab) = D(a) + aD(b)$$

(follows from $\widetilde{ab} = \widetilde{a} + a\widetilde{b}$)

and $D(1) = 0$, $\frac{\partial \kappa_i}{\partial \kappa_i} = \delta_{ii}$.

These derivatives can then be used to find boundaries of lifts of 2-cells in other covering spaces. ③

$E \xrightarrow{P} B$ be as above

$$B' = B \cup \{2\text{-cells } \sigma_j\} \quad \partial \sigma_j = \tau_j \in \mathbb{G}$$

i.e. $\tau_j =$ ordered product of group gens in free group \mathbb{G} .

Thus $\pi_1(B') = (\tau_1, \tau_2, \dots, \tau_n \mid \tau_1, \tau_2, \dots, \tau_m)$.

Let $E' \xrightarrow{P'} B'$ be a regular covering space corresp to a normal subgroup $H \triangleleft \mathbb{G}' = \pi_1(B')$.

Then if $\text{Sym}(E'/B') = \{f: E' \rightarrow E' \mid P'f = P', f \text{ homeom}\}$

(the group of deck transfs B' for E') then

$$\text{Sym}(E'/B') = \mathbb{G}'/H. \quad \boxed{a^{\psi} = \psi(a)}_{\det}$$

Let $\psi: \mathbb{G} = (\tau_1, \tau_2, \dots, \tau_n) \longrightarrow \mathbb{G}'/H$
be canonical map to this quotient.

Then: If $\tilde{\sigma}_j =$ lift of the cell σ_j to E' , then

$$\partial \tilde{\sigma}_j = (\partial \sigma_j)^{\psi} = \left(\frac{\partial \tau_1}{\partial \tau_1}\right)^{\psi} \tilde{\tau}_1 + \dots + \left(\frac{\partial \tau_m}{\partial \tau_m}\right)^{\psi} \tilde{\tau}_m$$

where $\tilde{\tau}_i =$ lift of τ_i into E' .

This formula shows how the Jacobian matrix $\left(\frac{\partial \tau_i}{\partial \tau_j}\right)^{\psi}$ can be used to compute

$$H_1(E').$$

Application to knot theory (and more generally)

$$H = [\mathbb{G}', \mathbb{G}'] = [\pi_1(B'), \pi_1(B')]$$

commutator subgroup.

Then $\mathbb{G}'/H =$ abelianization (\mathbb{G}') .

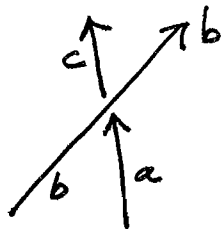
If we start with $\mathbb{G}' = \pi_1(S^3 - K)$

where $K \subset S^3$ is a knot (1-component)

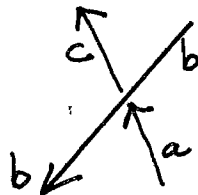
then $\psi: \mathbb{G}' \longrightarrow C_{\infty} = (\mathbb{Z} \mid)$

In the case of the Wirtinger presentation (3)

we have



$$ba = cb \quad \text{or} \quad c = bab^{-1}$$



$$ab = bc \quad \text{or} \quad c = b^{-1}ab$$

Each meridional generator $a, b, c, \dots \mapsto \mathcal{L}$ under ψ .

Each relation in $\pi_1(S^3 - K)$ corresponds to a relation in the Λ -module $H_1(E_\infty)$

where $\Lambda = \mathbb{Z}[\mathcal{L}, \mathcal{L}^{-1}]$ and $E_\infty = \text{coveringspace of } (S^3 - K) \text{ corres to } [\pi_1(S^3 - K), \pi_1(S^3 - K)]$.

$$(1) \quad ba = cb \Rightarrow \widetilde{ba} = \widetilde{cb}$$

$$\widetilde{b} + b\widetilde{a} = \widetilde{c} + c\widetilde{b}$$

$$\widetilde{c} = b\widetilde{a} + (1-c)\widetilde{b}$$

applying ψ : $\widetilde{c} = \mathcal{L}\widetilde{a} + (1-\mathcal{L})\widetilde{b}$

$$(2) \quad ab = bc \Rightarrow \widetilde{ab} = \widetilde{bc}$$

$$\widetilde{a} + a\widetilde{b} = \widetilde{b} + b\widetilde{c}$$

$$b\widetilde{c} = \widetilde{a} + (a-1)\widetilde{b}$$

$$\widetilde{c} = b^{-1}\widetilde{a} + (b^{-1}a - b^{-1})\widetilde{b}$$

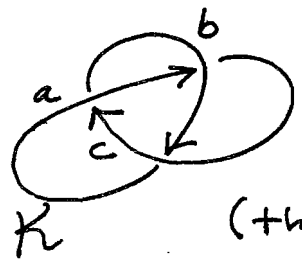
applying ψ : $\widetilde{c} = \mathcal{L}^{-1}\widetilde{a} + (1-\mathcal{L}^{-1})\widetilde{b}$

These then give the generators and relations for the Alex module $H_1(E_\infty)$.

Using Wirtinger presentation, it is easy to see that the cycles in $H_1(E_\infty)$ are gen by diffs $\mathcal{K}_i - \mathcal{K}_i$ (fixing \mathcal{K}_i). We have one redundant relation as well.

The upshot is that we get exactly the relation matrix for $H_1(E_\infty)$ by taking out one row and one column from $(\partial \kappa_i / \partial x_j)$.

ex:

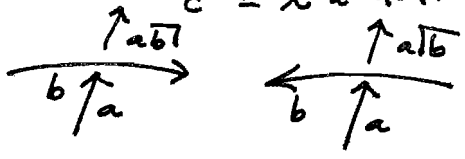


$$c = a \overline{b} = \tau a + (1-\tau)b$$

$$b = c \overline{a} = \tau c + (1-\tau)a$$

$$a = b \overline{c} = \tau b + (1-\tau)c$$

(this formalism replaces $\hat{c} = \tau \hat{a} + (1-\tau) \hat{b}$ etc.)



a	b	c
τ	$1-\tau$	-1
$1-\tau$	-1	τ
-1	τ	$1-\tau$

$$\Rightarrow \begin{pmatrix} -1 & \tau \\ \tau & 1-\tau \end{pmatrix} = \mathcal{J} \text{ relation matrix for } H_1(E_\infty)$$

We can determine the structure of $H_1(E_\infty)$ by doing row and col ops on \mathcal{J} over $\Lambda = \mathbb{Z}[\tau, \tau^{-1}]$.

$$\begin{pmatrix} -1 & \tau \\ \tau & 1-\tau \end{pmatrix} \xrightarrow{r} \begin{pmatrix} -1 & \tau \\ 0 & 1-\tau+\tau^2 \end{pmatrix} \xrightarrow{c} \begin{pmatrix} 1 & 0 \\ 0 & 1-\tau+\tau^2 \end{pmatrix}$$

This shows that $H_1(E_\infty)$ is a cyclic module: $H_1(E_\infty) \cong \Lambda / (1-\tau+\tau^2)\Lambda$.

We say that $\Delta_K(\tau) = \tau^2 - \tau + 1$ is the Alexander polynomial of K .

More generally, $\Delta_K(\tau)$ is determined up to $a \doteq b$ where $a \doteq b$ means $a = \pm \tau^N b$, and $\Delta_K(\tau)$ is defined to be the generator (in Λ) of the ideal gen by $(n-1) \times (n-1)$ minors in $(\frac{\partial \kappa_i}{\partial x_j})$.

where $(\kappa_1, \dots, \kappa_n \mid \kappa_1, \dots, \kappa_m) = \mathbb{G}'$ and $\mathbb{G}' / [\mathbb{G}', \mathbb{G}'] = C_\infty$.

It is an interesting fact of life that for knots and links, $H_1(E_\infty)$ can often be understood by using more geometric pictures of the covering space. (5)

$$K = \partial F \subset S^3$$

F an orientable spanning surface for $K \subset S^3$.

$$X = S^3 - K \text{ split along } F.$$

$$\partial X = F_- \cup F_+$$

$$E_\infty = \dots \tau^{-2}X \cup \tau^{-1}X \cup X \cup \tau X \cup \tau^2X \cup \dots$$

$$\dots \boxed{\tau^{-1}X \quad X \quad \tau X \quad \tau^2X} \dots$$

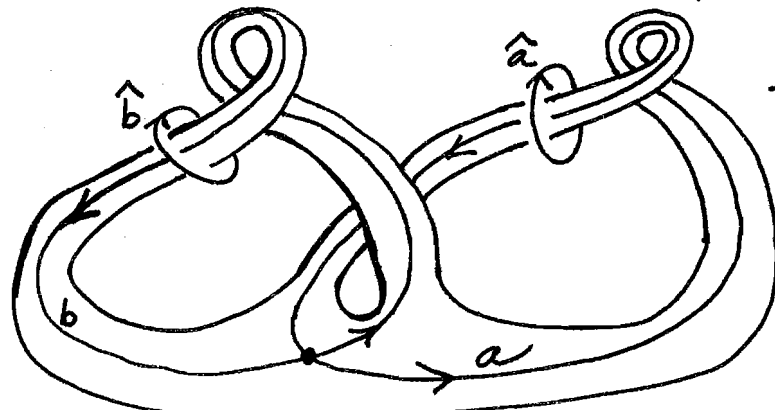
$$F_- \quad F_+$$

$$\boxed{\tau F_- = F_+}$$

Note that $X = S^3 - \text{Nbhd}(F)$. Thus we need to look at $H_1(S^3 - F)$.

$$\text{lk} : H_1(F) \times H_1(S^3 - F) \longrightarrow \mathbb{Z}$$

Alexander duality pairing



$$\begin{aligned} \text{lk}(\hat{a}, a) &= 1 \\ \text{lk}(\hat{a}, b) &= 0 \\ &\text{etc.} \end{aligned}$$

$$\boxed{H_1(S^3 - F) \cong H_1(F)}$$

Let $i : F \rightarrow S^3 - F$ via push by small amount along positive normal.

$$\text{Write } i(x) = x^*$$

$$\left. \begin{aligned} a^* &= i(a) = c_1 \hat{a} + d_1 \hat{b} \\ b^* &= i(b) = c_2 \hat{a} + d_2 \hat{b} \end{aligned} \right\} \Rightarrow \begin{aligned} \text{lk}(a^*, b) &= d_1 \\ \text{lk}(a^*, a) &= c_1 \\ \text{lk}(b^*, b) &= d_2 \\ \text{lk}(b^*, a) &= c_2 \end{aligned}$$

Define $\theta: H_1(F) \times H_1(F) \longrightarrow \mathbb{Z}$

(6)

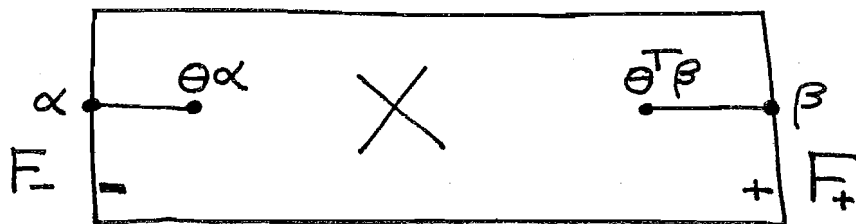
$$\theta(x, y) = \text{lk}(x^*, y).$$

- say $\{a_1, \dots, a_k\}$ basis for $H_1(F)$
 $\{\hat{a}_1, \dots, \hat{a}_k\}$ basis for $H_1(S^3 - F)$.

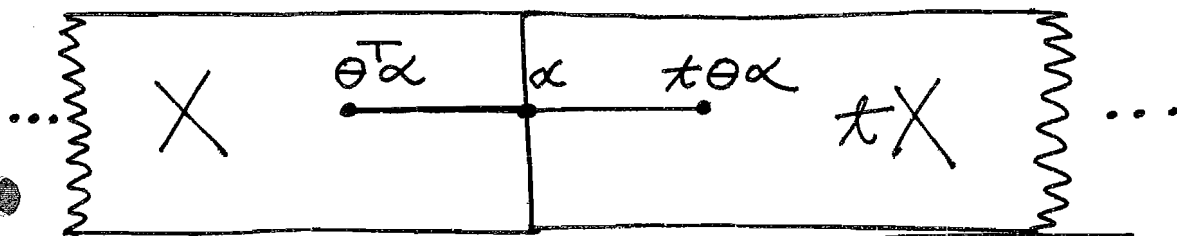
$$a_i^* = \hat{i}_i(a_j) = \sum_j M_{ji} \hat{a}_j$$

$$\Rightarrow \text{lk}(a_i^*, a_j) = M_{ji}$$

So $\theta^T = M =$ matrix of $\hat{i}: H_1(F) \rightarrow H_1(S^3 - F)$ with respect to the dual bases $\{a_i^*\}, \{\hat{a}_j\}$.



In E_{∞} we have the bases $\{\hat{a}_i\}, \{\tau \hat{a}_i\}, \dots$



Thus in $H_1(E_{\infty})$: $\theta^T \alpha = \tau \theta \alpha$.

In other words,

$$\theta^T - \tau \theta$$

is a relation matrix for $H_1(E_{\infty})$ over $\mathcal{A} = \mathbb{Z}[\tau, \tau^{-1}]$.

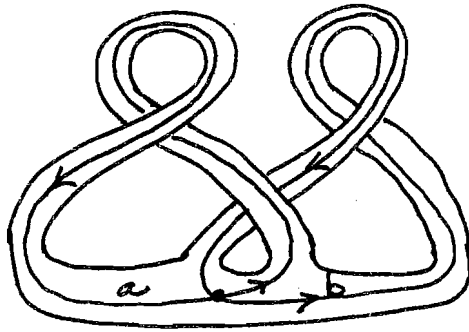
Thus we can obtain the structure of $H_1(E_{\infty})$ from $\theta^T - \tau \theta$

and

$$\Delta_K(\tau) \doteq \text{Det}(\theta^T - \tau \theta).$$

Ex.

(7)



θ	a	b
a	-1	0
b	1	-1

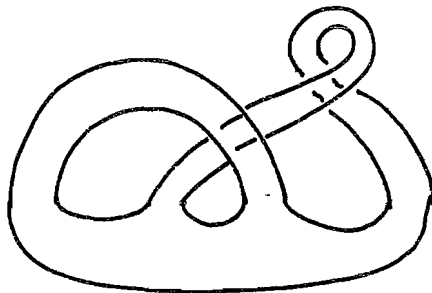


$$\theta = \begin{pmatrix} -1 & 0 \\ 1 & -1 \end{pmatrix}, \quad \theta^T = \begin{pmatrix} -1 & 1 \\ 0 & -1 \end{pmatrix}$$

$$\theta^T - t\theta = \begin{pmatrix} -1+t & -t \\ 1 & -1+t \end{pmatrix}$$

$$|\theta^T - t\theta| = 1 - 2t + t^2 + t = t^2 - t + 1$$

This recomputes the Alex polynomial of the trefoil knot.



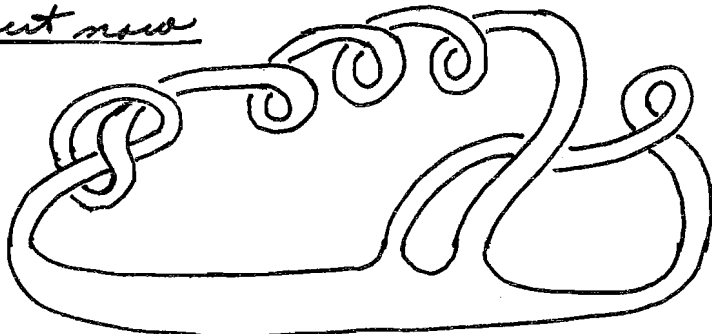
$$F, \quad \partial F \cong \mathbb{O}$$

$$\theta = \begin{pmatrix} 0 & 0 \\ 1 & -1 \end{pmatrix}$$

$$\theta^T - t\theta = \begin{pmatrix} 0 & 1 \\ 0 & -1 \end{pmatrix} - t \begin{pmatrix} 0 & 0 \\ 1 & -1 \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -t & -1+t \end{pmatrix}$$

$\Delta \doteq 1$ as expected since $\partial F = \text{unknot}$

But now



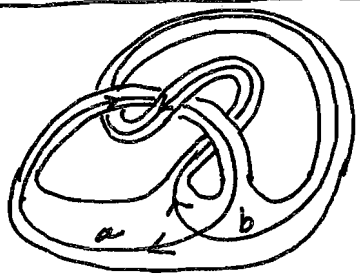
$$K = \partial F'$$

$$\theta = \begin{pmatrix} 0 & 0 \\ 1 & -1 \end{pmatrix}$$

$$\Delta_K \doteq 1$$

This shows how to construct knots with Alexander polynomial = 1.

Showing K knotted involves using the fundamental group or other knot invariants such as the Jones polynomial.



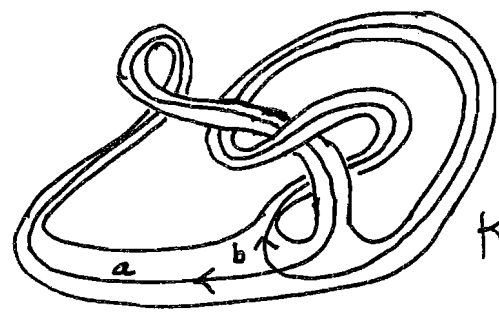
$K = \partial F$

Θ	a	b
a	0	1
b	2	0

$\Theta^T - t\Theta = \begin{pmatrix} 0 & 2-t \\ 1-2t & 0 \end{pmatrix}$

Alex module has two generators.

$\Delta_K \doteq (2-t)(1-2t)$



$K' = \partial F'$

Θ'	a	b
a	1	1
b	2	0

$\Theta'^T - t\Theta' = \begin{pmatrix} 1-t & 2-t \\ 1-2t & 0 \end{pmatrix}$

$\xrightarrow{c} \begin{pmatrix} -1 & 2-t \\ 1-2t & 0 \end{pmatrix}$

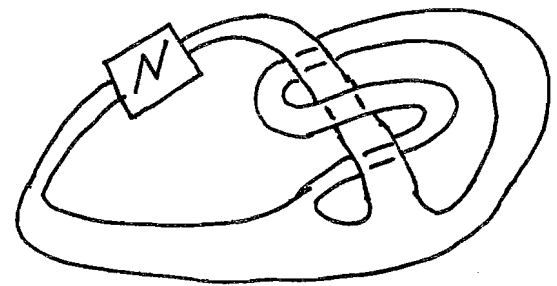
$\xrightarrow{r} \begin{pmatrix} -1 & 2-t \\ 0 & (1-2t)(2-t) \end{pmatrix}$

$\xrightarrow{c} \begin{pmatrix} 1 & 0 \\ 0 & (1-2t)(2-t) \end{pmatrix}$

Alex module has one generator.

$\Delta_{K'} \doteq (2-t)(1-2t)$

The above gives examples of knots K and K' with same Alex polynomial, but different (non-isom) Alexander modules. We can generalize this example to

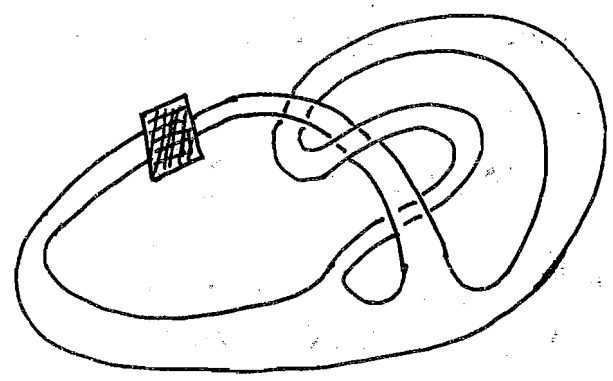


$\Theta = \begin{pmatrix} N & 1 \\ 2 & 0 \end{pmatrix}$

$\Theta^T - t\Theta = \begin{pmatrix} N-tN & 2-t \\ 1-2t & 0 \end{pmatrix}$

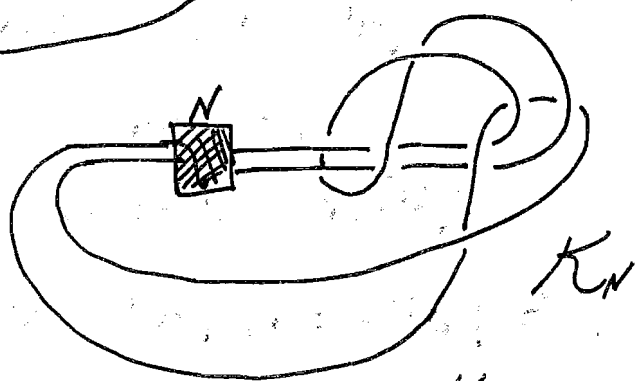
$\xrightarrow{r} \begin{pmatrix} -N & 2-t \\ 1-2t & 0 \end{pmatrix}$

Question. Are there Alex modules all non-isomorphic for $|N| \geq 3$?

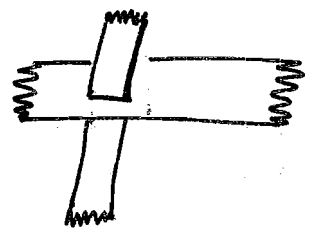


Exercise:
 Demonstrate this isotopy!

\approx

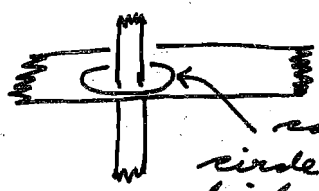


This shows that the knot K_N is a ribbon knot where this means that:
Defn. K is a ribbon knot if $K = \partial D$ where $D \hookrightarrow \mathbb{R}^3$ is an immersed disk with ribbon singularities



where a ribbon singularity consists in a transverse intersection of an interior arc of the disk with an arc that goes between two boundary points.

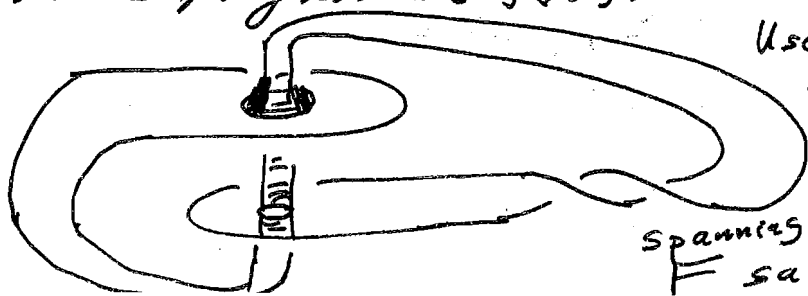
K ribbon $\Rightarrow K = \partial D$, $D \subset \mathbb{R}^4$ (upper 4 space). Pf:



cap off this circle with a disk in \mathbb{R}^4 . //

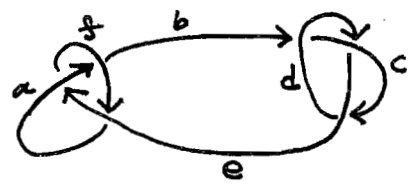
K ribbon $\Rightarrow \Delta_K \doteq f(t) f(1/t)$ for some polynomial $f(t)$.

Pf.



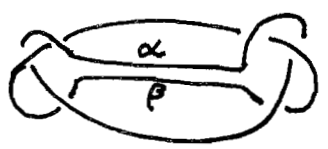
Use tubing trick to create spanning surface F such that

$$S^2 \hookrightarrow S^4$$

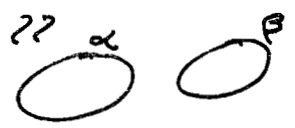


$$\begin{aligned} a\bar{f} &= b & e\bar{a} &= f \\ b\bar{d} &= c & f\bar{e} &= a \\ c\bar{e} &= d & & \\ d\bar{f} &= e & & \end{aligned}$$

⇒ need reln $f=d$



If go both above & below, then same grp.



$$\pi_1(S^3 - T \# T^*) : \begin{aligned} a\bar{f} &= b & e\bar{a} &= f & f\bar{e} &= a \\ b\bar{d} &= c & c\bar{e} &= d & \cancel{d\bar{f}} &= e \end{aligned}$$

Setting $f=d$: $a\bar{d} = b, e\bar{a} = d, d\bar{e} = a$
 $b\bar{d} = c, c\bar{e} = d$

$$\begin{aligned} da\bar{d} &= b, & ae\bar{a} &= d, & ed\bar{e} &= a \\ \bar{d}bd &= c & \bar{e}ce &= d \end{aligned}$$

$$\begin{aligned} \overbrace{da} &= \overbrace{bd}, & \overbrace{ae} &= \overbrace{da}, & \overbrace{ed} &= \overbrace{ae} \\ \overbrace{bd} &= \overbrace{dc}, & \overbrace{ce} &= \overbrace{ed} & & \end{aligned} \Rightarrow \begin{aligned} ae &= ce \\ \Rightarrow \underline{a=c} \end{aligned}$$

⇓

$$\begin{aligned} da &= bd, & ae &= da, & ed &= ae \\ \cancel{bd} &= \cancel{da}, & \cancel{ae} &= \cancel{ed} & d &= \bar{e}ae \end{aligned}$$

$$\underline{\bar{e}ae} = b\bar{e}ae, \quad ae = \underline{\bar{e}aea}$$

$$\Rightarrow b\bar{e} = 1 \Rightarrow \underline{b=e}$$

$$\Rightarrow \bar{e}aea = e\bar{e}ae, \quad ae = \bar{e}aea$$

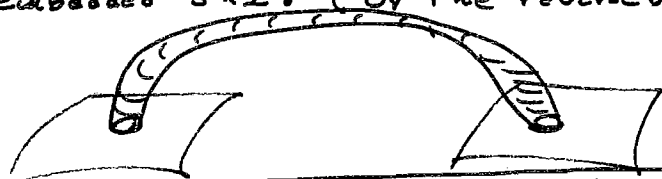
$$\Rightarrow \bar{e}aea = ae$$

$$\Rightarrow \boxed{aea = eae}$$

So get an $S^2 \hookrightarrow S^4$ with $\pi_1(S^4 - S^2) \cong (a, e \mid aea = eae) \cong \pi_1(\mathcal{P})$

S-Equivalence

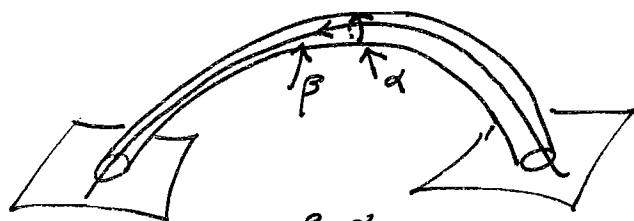
Two surfaces $\subset \mathbb{R}^3 \subset S^3$ are S-equivalent if one can be obtained from the other by a combination of ambient isotopy and tubing. Here tubing means cutting out two disjoint D^2 's + replacing by an embedded $S^1 \times I$. (Or the reverse)



Thm (See handout). $S, S' \subset \mathbb{R}^3 \subset S^3$ surfaces with boundary K, K' resp. Assume $K \cong K'$ (amb. iso.), then S and S' are S-equivalent.

From this we can prove that

- $\nabla_K = |\bar{x}^{-1} \bar{\theta}^T - \bar{x} \bar{\theta}|$ is a precise invariant of the link K : $K \cong K' \Rightarrow \nabla_K = \nabla_{K'}$.
- $\sigma(K) = \text{Signature}(\bar{\theta} + \bar{\theta}^T)$ is a precise invariant of K .



$\tilde{\theta}$ = Seifert pairing when the tube is added.

θ = Seifert pairing without the tube.

$$\tilde{\theta} = \begin{pmatrix} \theta & \beta & \alpha \\ \beta^T & b & 0 \\ \alpha & 0 & 1 & 0 \end{pmatrix}$$

Change of basis for Seifert pairing means that you can do an invertible row op paired with the identical column op.

$$\begin{pmatrix} \theta & b & 0 \\ \beta^T & N & 0 \\ 0 & 1 & 0 \end{pmatrix} \xrightarrow{r/c} \begin{pmatrix} \theta & 0 & 0 \\ \beta^T & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix} \text{ use this for } \tilde{\theta}$$

$$\bar{x}^{-1} \tilde{\theta}^T - \bar{x} \tilde{\theta} = \begin{pmatrix} \bar{x}^{-1} \bar{\theta}^T - \bar{x} \bar{\theta} & \bar{x}^{-1} b & 0 \\ -\bar{x} \beta^T & 0 & \bar{x}^{-1} \\ 0 & -\bar{x} & 0 \end{pmatrix}$$

$$\Rightarrow |\bar{x}^{-1} \tilde{\theta}^T - \bar{x} \tilde{\theta}| = |\bar{x}^{-1} \bar{\theta}^T - \bar{x} \bar{\theta}| \underbrace{\begin{vmatrix} 0 & \bar{x}^{-1} \\ -\bar{x} & 0 \end{vmatrix}}_1$$

This proves that ∇_K is precise.

$$\tilde{\Theta} \equiv \left(\begin{array}{c|cc} \Theta & 0 & 0 \\ \hline b^T & 0 & 0 \\ 0 & 1 & 0 \end{array} \right) \quad \tilde{\Theta} + \tilde{\Theta}^T = \left(\begin{array}{c|cc} \Theta & b & 0 \\ \hline b^T & 0 & 1 \\ 0 & 0 & 0 \end{array} \right)$$

$$\xrightarrow{r/c} \left(\begin{array}{c|cc} \Theta & 0 & 0 \\ \hline 0 & 0 & 1 \\ 0 & 1 & 0 \end{array} \right)$$

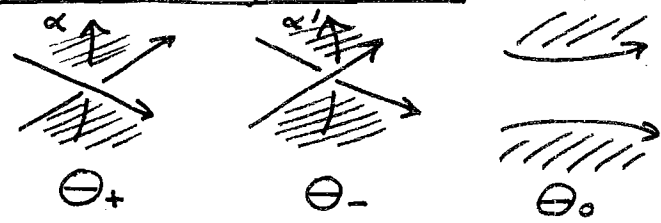
$$\text{Signature}(\tilde{\Theta}) = \text{Signature}(\Theta) + \underbrace{\text{Signature} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}}_{= \phi}$$

This proves that $\text{Signature}(\Theta) = \sigma^2(K)$ is a precise invariant. //

Note also that $|\tilde{\Theta} + \tilde{\Theta}^T| = -|\Theta + \Theta^T|$.

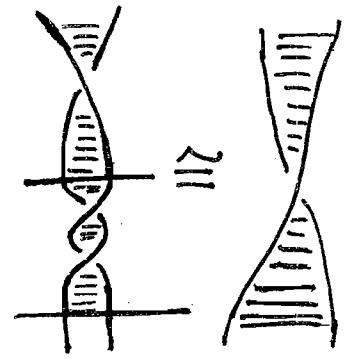
Abs Value $(|\Theta + \Theta^T|) \stackrel{\text{def}}{=} \text{Det}(K)$, the determinant of the knot K , is a precise invariant.

Skein Identity for ∇_K



$$\Theta_+ = \left(\begin{array}{c|c} n & \chi^T \\ \hline \chi & \Theta_0 \end{array} \right)$$

$$\Theta_- = \left(\begin{array}{c|c} n+1 & \chi^T \\ \hline \chi & \Theta_0 \end{array} \right)$$



Will also discuss

- Art Invariant
- Knot Cobordism
- Homfly etc...

(state sums)

$$\Rightarrow \nabla_+ - \nabla_- = [(x^{-1} - x)n - (x^{-1} - x)(n+1)] \nabla_0$$

$$\boxed{\nabla_+ - \nabla_- = (x - x^{-1}) \nabla_0}$$

Conway
Axioms

$$\nabla_{\nearrow} - \nabla_{\searrow} = z \nabla_{\rightarrow}, \quad z = x - x^{-1}$$

$$\nabla_{\bigcirc} = 1$$

We have proved that ∇_K satisfies the axioms for the Conway "potential function".