A RIGOROUS NUMERICAL ANALYSIS OF THE TRANSFORMED
FIELD EXPANSION METHOD FOR DIFFRACTION BY PERIODIC,
LAYERED STRUCTURES *

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Abstract. Boundary perturbation methods have received considerable attention in recent years
due to their ability to simulate solutions of differential equations of applied interest in a stable,
robust, and highly accurate fashion. In this contribution we study the rigorous numerical analysis of
a recently proposed High–Order Perturbation of Surfaces method for scattering of electromagnetic
waves by a doubly–layered, periodic medium in Transverse Electric polarization. The algorithm in
question is a Transformed Field Expansion method which is discretized with a Fourier–Legendre–
Galerkin, Taylor series approach. We prove not only results on existence and uniqueness of solutions,
but also theorems indicating that solutions of our scheme converge to these solutions with High–Order
Spectral accuracy.

Key words. High–Order Perturbation of Surfaces methods, High–Order Spectral methods,
Helmholtz equation, diffraction gratings, layered media.

AMS subject classifications. 65N35, 78A45, 78B22.

1. Introduction. We consider here the scattering of a time–harmonic elec-
 tromagnetic plane wave by a periodically corrugated grating structure [25]. The scat-
tering of linear waves involving periodic layered media plays a crucial role in a wide
range of engineering and physics applications, e.g., materials science [10], nondestruc-
tive testing [31], sensing [12], geophysics [32], imaging [19], oceanography [2], and
nanoplasmonics [27].

A number of computational methods have been developed for problems of scat-
tering by periodic gratings. The most popular approaches to these problems are
volumetric methods such as Finite Differences and Finite/Spectral Element methods
[7, 1] but these methods are greatly disadvantaged with an unnecessarily large number
of unknowns for piecewise homogeneous grating problems [20]. Interfacial methods
based on Integral Equations (IEs) [6, 4, 17] are a natural alternative but these also face
several challenges. First, for periodic problems, the relevant Green function must be
periodized which greatly increases the computational cost. Additionally, these non–
local IEs produce dense, non–symmetric positive definite systems of linear equations
which must be inverted with each simulation.

A High–Order Perturbation of Surfaces (HOPS) approach can avoid these con-
cerns, such as the method of Transformed Field Expansions (TFE) [21, 22] which
we study here. These high–order algorithms were first developed by Bruno and Re-
itich for the two–dimensional scalar case [3] and later enhanced and stabilized by
Nicholls and Reitich [21, 22], and Nicholls and Malcolm [18]. HOPS approaches are
compelling as they maintain the advantageous properties of classical IE formulations
(e.g., surface formulation and exact enforcement of far–field boundary conditions)
while avoiding many of their shortcomings. For instance, since HOPS schemes utilize
complex exponentials as basis functions in the lateral variable, the quasi–periodicity

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of solutions does not need to be explicitly enforced. In addition, due to the nature of
the scheme, at every perturbation order one need only invert a single, sparse operator
corresponding to the flat–interface, order–zero approximation of the problem. The
TFE method studied in this contribution was generalized by Nicholls and Shen to the
case of irregular bounded obstacles in two [23] and three dimensions [9]. They later
delivered a rigorous numerical analysis of the method [24] and we follow their strategy
in this contribution. Subsequently, in [11, 14, 13, 15] the algorithms were extended to
the case of periodic gratings separating multiply layered materials, whose solutions
are governed by either Helmholtz equations or the full Maxwell equations.

In this paper, we conduct a rigorous numerical analysis of the method developed
by the authors [14, 13, 15] in the case of a doubly layered material with solutions
satisfying a pair of Helmholtz equations coupled via the boundary conditions at the
interface between the two. The TFE algorithm we derived is not only a stable and
high–order numerical scheme, but it can also be used to directly establish the exis-
tence, uniqueness, and analyticity of solutions, as we presently demonstrate. For this
purpose we establish a classical, but non–trivial, elliptic existence, uniqueness, and
regularity theory by using the Green function and a priori estimates. The proof of
our main result is based upon analyticity estimates for the TFE expansions coupled
to the convergence of the Fourier–Legendre–Galerkin method. Our developments il-
lustrate the power and flexibility of the TFE approach for both numerical simulation
and theoretical analysis.

2. Governing equations. To specify the problem and its geometry we consider
the two–dimensional Helmholtz problem which governs the scattering of electromagnetic
waves in Transverse Electric (TE) polarization [25]

\begin{align}
(2.1a) & \quad \Delta u + k_1^2 u = 0, \quad \text{in } z > g(x), \\
(2.1b) & \quad \Delta v + k_2^2 v = 0, \quad \text{in } z < g(x), \\
(2.1c) & \quad u - v = -u^{inc}, \quad \text{at } z = g(x), \\
(2.1d) & \quad \partial_N u - \partial_N v = -\partial_N u^{inc}, \quad \text{at } z = g(x), \\
(2.1e) & \quad \text{OWC}[u] = 0, \quad \text{as } z \to \infty, \\
(2.1f) & \quad \text{OWC}[v] = 0, \quad \text{as } z \to -\infty, \\
(2.1g) & \quad u(x + d, z) = e^{i\alpha d} u(x, z), \\
(2.1h) & \quad v(x + d, z) = e^{i\alpha d} v(x, z),
\end{align}

where $u^{inc} = e^{i\alpha x - i\gamma z}$, $\partial_N$ is an upward pointing normal derivative, and “OWC”
connotes the Outgoing Wave Condition which we make precise presently.

2.1. Transparent Boundary Conditions. The usual procedure when imple-
menting the TFE method is to truncate (if necessary) the unbounded problem domain
to one of finite extent. For this we introduce artificial boundaries above and be-
low the structure, and enforce transparent boundary conditions to equivalently solve
(2.1). Introducing the planes $\{z = a > |g|_{L^\infty}\}$ and $\{z = b < -|g|_{L^\infty}\}$ we show that
transparent boundary conditions can be enforced at these with Dirichlet–Neumann
Operators (DNOs) derived from the Rayleigh expansions [25]. These expansions are
relevant as they are the explicit solutions (obtained from Separation of Variables) of
the problems on $\{z > a\}$ and $\{z < b\}$ upon specification of Dirichlet data at the
artificial boundaries, \( \{ z = a \} \) and \( \{ z = b \} \). More specifically, it is known [25] that

\[
\begin{align*}
  u(x, z) &= u(x, z) = \sum_{p = -\infty}^{\infty} \hat{\zeta}_p e^{i\alpha_p x} e^{i\gamma_{l,p}(z-a)}, \quad z > a, \\
  v(x, z) &= v(x, z) = \sum_{p = -\infty}^{\infty} \hat{\psi}_p e^{i\alpha_p x} e^{i\gamma_{2,p}(b-z)}, \quad z < b,
\end{align*}
\]

where

\[
\alpha_p := \alpha + \frac{2\pi}{d} p, \quad \gamma_{l,p} = \begin{cases} 
  \sqrt{k_l^2 - \alpha_p^2}, & \alpha_p^2 \leq k_l^2, \\
  i\sqrt{\alpha_p^2 - k_l^2}, & \alpha_p^2 > k_l^2,
\end{cases}
\]

for \( l = 1, 2 \). We note that, upon evaluating at the artificial boundaries,

\[
\begin{align*}
  u(x, a) &= \sum_{p = -\infty}^{\infty} \hat{\zeta}_p e^{i\alpha_p x} =: \zeta(x), \quad v(x, b) = \sum_{p = -\infty}^{\infty} \hat{\psi}_p e^{i\alpha_p x} =: \psi(x),
\end{align*}
\]

and from these we can compute the Neumann data at the artificial boundaries,

\[
\begin{align*}
  \partial_z u(x, a) &= \sum_{p = -\infty}^{\infty} (i\gamma_{1,p}) \hat{\zeta}_p e^{i\alpha_p x}, \quad \partial_z v(x, b) = \sum_{p = -\infty}^{\infty} (-i\gamma_{2,p}) \hat{\psi}_p e^{i\alpha_p x}.
\end{align*}
\]

With these we define the DNOs

\[
\begin{align*}
  T_1[\zeta] &= T_1[u(x, a)] := \sum_{p = -\infty}^{\infty} (i\gamma_{1,p}) \hat{\zeta}_p e^{i\alpha_p x}, \\
  T_2[\psi] &= T_2[v(x, b)] := \sum_{p = -\infty}^{\infty} (-i\gamma_{2,p}) \hat{\psi}_p e^{i\alpha_p x},
\end{align*}
\]

which are order–one Fourier multipliers. Using these we can state (2.1) equivalently on the bounded domain \( \{ b < z < a \} \) as

\[
\begin{align*}
  &\Delta u + k_l^2 u = 0, \quad \text{in } g(x) < z < a, \\
  &\Delta v + k_l^2 v = 0, \quad \text{in } b < g(x) < z, \\
  &u - v = -u^{inc}, \quad \text{at } z = g(x), \\
  &\partial_N u - \partial_N v = -\partial_N u^{inc}, \quad \text{at } z = g(x), \\
  &\partial_z u - T_1[u] = 0, \quad \text{at } z = a, \\
  &\partial_z v - T_2[v] = 0, \quad \text{at } z = b, \\
  &u(x + d, z) = e^{i\alpha_d} u(x, z), \\
  &v(x + d, z) = e^{i\alpha_d} v(x, z).
\end{align*}
\]

3. Transformed Field Expansions. We now recall the TFE method [21, 22] which begins with a domain flattening change of variables (also known as \( \sigma \)–coordinates [26] in the geophysical literature and the C–method [5] in the electromagnetics community). Subsequently, we make a boundary perturbation expansion which is solved recursively at each perturbation order.
3.1. The Change of Variables. We define the change of variables \( x' = x, \)
\[
z_1 = a \left( \frac{z-g}{a-g} \right) \quad \text{for} \quad g < z < a, \quad z_2 = b \left( \frac{g-z}{g-b} \right) \quad \text{for} \quad b < z < g,
\]
and define
\[
U_1(x', z_1) := u(x(x'), z(x', z_1, z_2)), \quad U_2(x', z_2) := v(x(x'), z(x', z_1, z_2)).
\]
Using this change of variables, a long computation (see Section 5) transforms (2.2) to
the following system of equations
\[
\begin{align*}
\Delta_1 U_1 + k_1^2 U_1 &= \frac{1}{G_1} (\partial_x R_1^x + \partial_z R_1^z + R_1^0) =: R_1, \quad \text{in} \quad 0 < z_1 < a, \\
\Delta_2 U_2 + k_2^2 U_2 &= \frac{1}{G_2} (\partial_x R_2^x + \partial_z R_2^z + R_2^0) =: R_2, \quad \text{in} \quad b < z_2 < 0, \\
U_1 - U_2 &= \xi_1, \quad \text{at} \quad z_1 = z_2 = 0, \\
\partial_z U_1 - \partial_z U_2 &= \xi_2, \quad \text{at} \quad z_1 = z_2 = 0, \\
\partial_z U_1 - T_1[U_1] &= -\frac{g}{a} T_1[U_1] =: J_1, \quad \text{at} \quad z_1 = a, \\
\partial_z U_2 - T_2[U_2] &= -\frac{g}{b} T_2[U_2] =: J_2, \quad \text{at} \quad z_2 = b, \\
U_1(x' + d, z_1) &= e^{i\omega d} U_1(x', z_1), \\
U_2(x' + d, z_2) &= e^{i\omega d} U_2(x', z_2),
\end{align*}
\]
where the Laplacian operator \( \Delta \) is defined by \( \Delta = \partial_x^2 + \partial_z^2 \), for \( l = 1, 2 \). We refer
the reader to Section 5 for the specific formulas for \( R_l \) and \( \xi_l \).

3.2. A High–Order Perturbation of Surfaces Method. We now introduce
a boundary perturbation method to solve the transformed governing equations, (3.1).
To begin, we assume that the deformation has the form
\[
g(x') = \varepsilon f(x'), \quad f = O(1),
\]
and expand the fields in power series
\[
\{U_1, U_2\} = \sum_{n=0}^{\infty} \{U_{1,n}, U_{2,n}\} \varepsilon^n.
\]
Inserting these expansion into (3.1) and equating at order \( O(\varepsilon^n) \) delivers
\[
\begin{align*}
\Delta_1 U_{1,n} + k_1^2 U_{1,n} &= R_{1,n}, \quad \text{in} \quad 0 < z_1 < a, \\
\Delta_2 U_{2,n} + k_2^2 U_{2,n} &= R_{2,n}, \quad \text{in} \quad b < z_2 < 0, \\
U_{1,n} - U_{2,n} &= \xi_{1,n}, \quad \text{at} \quad z_1 = z_2 = 0, \\
\partial_z U_{1,n} - \partial_z U_{2,n} &= \xi_{2,n}, \quad \text{at} \quad z_1 = z_2 = 0, \\
\partial_z U_{1,n} - T_1[U_{1,n}] &= -\frac{f}{a} T_1[U_{1,n-1}] =: J_{1,n}, \quad \text{at} \quad z_1 = a, \\
\partial_z U_{2,n} - T_2[U_{2,n}] &= -\frac{f}{b} T_2[U_{2,n-1}] =: J_{2,n}, \quad \text{at} \quad z_2 = b, \\
U_{1,n}(x' + d, z_1) &= e^{i\omega d} U_{1,n}(x', z_1), \\
U_{2,n}(x' + d, z_2) &= e^{i\omega d} U_{2,n}(x', z_2).
\end{align*}
\]
Again, we refer the reader to the Section 5 for the specific formulas for the right hand sides $R_{l,n}$ and $\xi_{l,n}$.

Considering the quasiperiodicity of solutions, we propose the following generalized Fourier (Floquet) series expansions

\begin{align}
(3.3a) \quad U_{l,n}(x', z_l) &= \sum_{p=-\infty}^{\infty} \hat{U}^{(p)}_{l,n}(z_l)e^{i\alpha_p x'}, \quad R_{l,n}(x', z_l) &= \sum_{p=-\infty}^{\infty} \hat{R}^{(p)}_{l,n}(z_l)e^{i\alpha_p x'}, \\
(3.3b) \quad J_{l,n}(x') &= \sum_{p=-\infty}^{\infty} \hat{J}^{(p)}_{l,n}e^{i\alpha_p x'}, \quad \xi_{l,n}(x') &= \sum_{p=-\infty}^{\infty} \hat{\xi}^{(p)}_{l,n}e^{i\alpha_p x'},
\end{align}

for $l = 1, 2$. Inserting these expansions into (3.2), the governing equations are reduced to the one-dimensional boundary value problems

\begin{align}
(3.4a) \quad \partial_{z_1}^2 \hat{U}^{(p)}_{1,n} + (k_1^2 - \alpha_p^2)\hat{U}^{(p)}_{1,n} &= \hat{R}^{(p)}_{1,n}, \quad \text{in } 0 < z_1 < a, \\
(3.4b) \quad \partial_{z_2}^2 \hat{U}^{(p)}_{2,n} + (k_2^2 - \alpha_p^2)\hat{U}^{(p)}_{2,n} &= \hat{R}^{(p)}_{2,n}, \quad \text{in } b < z_2 < 0, \\
(3.4c) \quad \xi^{(p)}_{1,n} = \hat{U}^{(p)}_{1,n} - \hat{U}^{(p)}_{2,n} = \xi^{(p)}_{2,n}, \quad \text{at } z_1 = z_2 = 0, \\
(3.4d) \quad \partial_{z_1} \hat{J}^{(p)}_{1,n} - \partial_{z_2} \hat{J}^{(p)}_{2,n} &= \hat{\xi}^{(p)}_{2,n}, \quad \text{at } z_1 = z_2 = 0, \\
(3.4e) \quad \partial_{z_1} \hat{J}^{(p)}_{1,n} - i\gamma_{1,p}\hat{U}^{(p)}_{1,n} &= -\frac{f}{a}(i\gamma_{1,p})\hat{U}^{(p)}_{1,n-1} = : \hat{J}^{(p)}_{1,n}, \quad \text{at } z_1 = a, \\
(3.4f) \quad \partial_{z_2} \hat{J}^{(p)}_{2,n} + i\gamma_{2,p}\hat{U}^{(p)}_{2,n} &= -\frac{f}{b}(-i\gamma_{2,p})\hat{U}^{(p)}_{2,n-1} = : \hat{J}^{(p)}_{2,n}, \quad \text{at } z_2 = b.
\end{align}

4. Function Spaces. In order to use these TFE recursions in a direct proof of the existence, uniqueness, and analyticity of the solutions $\{u, v\}$ of (2.2), we must define our function spaces and state properties of these. To start, we recall, for the $L^2$ function $f = f(x')$, the classical Sobolev norm for any real $s \geq 0$ [16],

$$
\|f\|_{H^s}^2 := \sum_{p=-\infty}^{\infty} \langle p \rangle^{2s} |\hat{f}_p|^2, \quad \langle p \rangle^2 := 1 + |p|^2, \quad \hat{f}_p := \frac{1}{a} \int_0^d f(x')e^{-i\alpha_p x'} dx'.
$$

For the $L^2$ function $w = w(x', z)$ we require the classical Sobolev norm for any integer $s \geq 0$ [8],

$$
\|w\|_{H^s}^2 := \sum_{k=0}^{s} \sum_{p=-\infty}^{\infty} \langle p \rangle^{2(s-k)} \int_b^a |\partial_z^k \hat{w}_p(z)|^2 dz.
$$

With these norms we define the function spaces, for real $s \geq 0$,

$$
H^s([0,d]) := \{ f(x') \in L^2([0,d]) \mid \|f\|_{H^s} < \infty \},
$$

and, for integer $s \geq 0$,

$$
H^s([0,d] \times [a,b]) := \{ w(x', z) \in L^2([0,d] \times [a,b]) \mid \|w\|_{H^s} < \infty \}.
$$

Additionally, we will require their duals, $H^{-s}$ [8].

We recall the following algebra property of Sobolev spaces (see, e.g., [21]) which allows us to estimate products of elements in these classes.
Lemma 4.1. Given any integer \( s \geq 0 \) and any \( \sigma > 0 \), there exists a constant \( \kappa = \kappa(s, \sigma) \) such that if \( f \in C^s([0, d]) \) and \( w \in H^s([0, d] \times [b, a]) \) then

\[
\|f w\|_{H^s} \leq \kappa \|f\|_{C^s} \|w\|_{H^s},
\]

and if \( \tilde{f} \in C^{s+1/2+\sigma}([0, d]) \) and \( \tilde{w} \in H^{s+1/2}([0, d]) \) then

\[
\|\tilde{f} \tilde{w}\|_{H^{s+1/2}} \leq \kappa \|\tilde{f}\|_{C^{s+1/2+\sigma}} \|\tilde{w}\|_{H^{s+1/2}}.
\]

We also recall an elementary property of \( H^s \).

Lemma 4.2. Given any integer \( s \geq 0 \), if \( F \in H^s([0, d] \times [b, a]) \), then \( (a-z)F \in H^s([0, d] \times [b, a]) \) and there exists a positive constant \( Z_a = Z_a(s) \) such that

\[
\|(a-z)F\|_{H^s} \leq Z_a \|F\|_{H^s}.
\]

As we shall see, the key tool for establishing our result is the following elliptic estimate which allows us to show that unique solutions exist to the prototype problem above, (3.2), in an appropriate Sobolev space.

Theorem 4.3. Given any integer \( s \geq 0 \), if \( F_1, F_2 \in H^{s-1}([0, d] \times [b, a]); \xi \in H^{s+1/2}([0, d]), \) and \( \nu, K_1, K_2 \in H^{s-1/2}([0, d]) \), then there exists a unique solution pair \( \{u, v\} \in H^{s+1}([0, d] \times [b, a]) \times H^{s+1}([0, d] \times [b, a]) \) of

\[
\begin{align*}
(4.3a) & \quad \Delta_1 u + k_1^2 u = F_1, & 0 < z_1 < a, \\
(4.3b) & \quad \Delta_2 v + k_2^2 v = F_2, & b < z_1 < 0, \\
(4.3c) & \quad u - v = \xi, & z_1 = z_2 = 0, \\
(4.3d) & \quad \partial_{z_1} u - \partial_{z_2} v = \nu, & z_1 = z_2 = 0, \\
(4.3e) & \quad \partial_{z_1} u - T_1 [u] = K_1, & z_1 = a, \\
(4.3f) & \quad \partial_{z_2} v - T_2 [v] = K_2, & z_2 = b, \\
(4.3g) & \quad u(x' + d, z) = e^{i\sigma d} u(x', z), \\
(4.3h) & \quad v(x' + d, z) = e^{i\sigma d} v(x', z),
\end{align*}
\]

such that, for a universal constant \( K_\varepsilon \),

\[
\max \{\|u\|_{H^{s+1}}, \|v\|_{H^{s+1}}\} \leq K_\varepsilon \{\|F_1\|_{H^{s-1}} + \|F_2\|_{H^{s-1}} + \|\xi\|_{H^{s+1/2}} + \|\nu\|_{H^{s-1/2}} + \|K_1\|_{H^{s-1/2}} + \|K_2\|_{H^{s-1/2}}\}.
\]

We give the proof in Appendix A.

5. Existence, Uniqueness, and Analyticity. To study the existence, uniqueness, and analyticity of solutions we recall (3.2) and present precise expressions for the terms on the right hand sides. Recalling that \( R_{l,n} = \partial_{x'} \xi_{l,n}^2 + \partial_{z_1} \eta_{l,n}^2 + R_{l,n}^0 \), it
can be shown that
\[
R_{1,n}^0 = \frac{a}{2} f \partial_x U_{1,n-1} + \frac{a - z_1}{a} (\partial_x f) \partial_z U_{1,n-1} - \frac{f^2}{a^2} \partial_x U_{1,n-2} - \frac{a - z_1}{a} f(\partial_x f) \partial_z U_{1,n-2},
\]
\[
R_{1,n}^1 = \frac{a}{2} (\partial_x f) \partial_z U_{1,n-1} - \frac{a - z_1}{a} (\partial_x f) \partial_z U_{1,n-2} - \frac{(a - z_1)^2}{a^2} (\partial_x f)^2 \partial_z U_{1,n-2},
\]
\[
R_{1,n}^0 = -\frac{1}{a} (\partial_x f) U_{1,n-1} + (k_1^2) \frac{2f}{a} U_{1,n-1} + \frac{1}{a^2} f(\partial_x f) \partial_x U_{1,n-2} + \frac{a - z_1}{a^2} (\partial_x f)^2 \partial_z U_{1,n-2} - k_1^2 f^2 U_{1,n-2},
\]
similarly for \( R_{2,n} \), and
\[
\xi_{1,n} = (-1)^{n+1} \frac{(i\gamma f)^n}{n!} e^{i\alpha x}, \quad \xi_{2,n} = \frac{Q_{1,n} + Q_{2,n}}{ab},
\]
where
\[
Q_{1,n} = -iab\gamma \xi_{1,n} - iabo(\partial_x f)\xi_{1,n-1} - i\gamma(a - b) f \xi_{1,n-1} - i\alpha(a - b) f(\partial_x f) \xi_{1,n-2} + i\gamma f^2 \xi_{1,n-2} + i\alpha(\partial_x f) f^2 \xi_{1,n-3},
\]
\[
Q_{2,n} = -af \partial_z U_{1,n-1} + ab(\partial_x f) \partial_x U_{1,n-1} + (a - b) f(\partial_x f) \partial_x U_{1,n-2} - ab(\partial_x f)^2 \partial_z U_{1,n-2} - (\partial_x f)^2 \partial_x U_{1,n-3} - a(\partial_x f) f(\partial_x f)^2 \partial_x U_{1,n-3} - b(\partial_x f) f(\partial_x f)(\partial_x f)^2 \partial_x U_{2,n-2} + ab(\partial_x f)^2 \partial_z U_{2,n-2} + (\partial_x f)^2 \partial_x U_{2,n-3} - b(\partial_x f)^2 \partial_x U_{2,n-3}.
\]

To begin our demonstration we establish the analyticity of the Dirichlet data.

**Lemma 5.1.** Given any integer \( s \geq 0 \), if \( f \in C^{s+2}([0,d]) \), then
\[
\| \xi_{1,n} \|_{H^{s+3/2}} \leq K_\xi B^n_\xi,
\]
for constants \( K_\xi, B_\xi > 0 \).

**Proof.** We note that \( \xi_{1,n} = -i\gamma f \xi_{1,n-1}/n \) and use induction to prove this lemma. We begin at \( n = 0 \) and set
\[
K_\xi := \| \xi_{1,0} \|_{H^{s+3/2}}.
\]
We now assume (5.1) for all \( n < \bar{n} \) and consider \( \bar{n} > 1 \) where we bound
\[
\| \xi_{1,\bar{n}} \|_{H^{s+3/2}} \leq |\gamma| M |f|_{C^{s+3/2+\bar{n}}} \| \xi_{1,\bar{n}-1} \|_{H^{s+3/2}} \leq |\gamma| M |f|_{C^{s+2}} K_\xi B^{\bar{n}-1}_\xi.
\]
By choosing \( B_\xi > M|\gamma| |f|_{C^{s+2}} \) the lemma follows.

We now provide the key inductive lemma which enables the proof of our result.

**Lemma 5.2.** Given any integer \( s \geq 0 \), if \( f \in C^{s+2}([0,d]) \) and
\[
\| U_{1,n} \|_{H^{s+2}} + \| U_{2,n} \|_{H^{s+2}} \leq KB^n, \quad \forall n < \bar{n},
\]
for constants $K, B > 0$, there exists a constant $C > 0$ such that

$$\max \left\{ \|R_{l,n}\|_{H^s}, \|J_{l,n}\|_{H^{s+1/2}}, \|\xi_{2,n}\|_{H^{s+1/2}} \right\} \leq KC \left( B^{n-1} + B^{n-2} + B^{n-3} \right),$$

for $l = 1, 2$.

Proof. For $l = 1, 2$, we recall that $R_{l,n} = \partial_{x^r} R_{l,n}^x + \partial_{x^z} R_{l,n}^z + R_{l,n}^0$, so that one can deduce

$$\|R_{l,n}\|_{H^s} \lesssim \|R_{l,n}^x\|_{H^{s+1}} + \|R_{l,n}^z\|_{H^{s+1}} + \|R_{l,n}^0\|_{H^s},$$

where $\|A\| \lesssim \|B\|$ means that there exists a constant $C$, independent of all variables of importance, such that $\|A\| \leq C \|B\|$. With the estimates

$$\|R_{l,n}^x\|_{H^{s+1}} \lesssim \left| f \partial_{x^r} U_{l,n-1} \right|_{H^{s+1}} + \left| \partial_{x^r} f \partial_{x^z} U_{l,n-1} \right|_{H^{s+1}} + \left| f \partial_{x^r} \partial_{x^z} U_{l,n-2} \right|_{H^{s+1}},$$

and

$$\|R_{l,n}^z\|_{H^{s+1}} \lesssim \left| f \partial_{x^r} \partial_{x^z} U_{l,n-1} \right|_{H^{s+1}} + \left| \partial_{x^r} f \partial_{x^z} U_{l,n-2} \right|_{H^{s+1}},$$

and

$$\|R_{l,n}^0\|_{H^s} \lesssim \left| \partial_{x^r} f \partial_{x^z} U_{l,n-1} \right|_{H^s} + \left| f \partial_{x^r} U_{l,n-1} \right|_{H^s} + \left| f \partial_{x^r} \partial_{x^z} U_{l,n-2} \right|_{H^s},$$

and

$$\|R_{l,n}^0\|_{H^s} \lesssim \left| f \partial_{x^r} U_{l,n-1} \right|_{H^s} + \left| f \partial_{x^r} \partial_{x^z} U_{l,n-2} \right|_{H^s},$$

we find that

$$\|R_{l,n}\|_{H^s} \lesssim K \left( \left| f \right|_{C^{s+2}} B^{n-1} + \left| f \right|_{C^{s+2}}^2 B^{n-2} \right).$$

For $J_{l,n}$ we can show that

$$\|J_{l,n}\|_{H^{s+1/2}} \lesssim \left| f \right|_{C^{s+1/2} \sigma} U_{l,n-1} \|_{H^{s+3/2}} \lesssim \left| f \right|_{C^{s+1/2} \sigma} B^{n-1}.$$

Hence, we deduce that

$$\max \left\{ \|R_{l,n}\|_{H^s}, \|J_{l,n}\|_{H^{s+1/2}} \right\} \leq KC \left( f \right|_{C^{s+2}} B^{n-1} + \left| f \right|_{C^{s+2}}^2 B^{n-2} \right) \lesssim K \left( B^{n-1} + B^{n-2} \right).$$

It remains to estimate $\xi_{2,n}$ and, for this, we use Lemma 5.1 which implies

$$\|\xi_{1,n}\|_{H^{s+1/2}} \leq K \xi B^n,$$
and the lemma follows.

In addition, we find, for $l = 1, 2$,

\[ \|Q_{2,n}\|_{H^{s+1/2}} \lesssim \|\partial_z U_{1,n-1}\|_{H^{s+1/2}} + \|\partial_x U_{1,n-1}\|_{H^{s+1/2}} + \|\partial_x U_{1,n-2}\|_{H^{s+1/2}} \]
\[ + \|\partial_z U_{1,n-2}\|_{H^{s+1/2}} + \|\partial_x U_{1,n-3}\|_{H^{s+1/2}} + \|\partial_z U_{1,n-3}\|_{H^{s+1/2}} \]
\[ \lesssim \|U_{1,n-1}\|_{H^{s+2}} + \|U_{1,n-2}\|_{H^{s+2}} + \|U_{1,n-3}\|_{H^{s+2}} \]
\[ \lesssim K (B^{n-1} + B^{n-2} - B^{n-3}), \]

and the lemma follows.

We can now state and prove the main theorem of this section.

**Theorem 5.3.** Given any integer $s \geq 0$, if $f \in C^{s+2}([0, d])$ and $\xi_{1,n} \in H^{s+3/2}([0, d])$ such that

\[ \|\xi_{1,n}\|_{s+3/2} \leq K \xi B^n, \]

for constants $K, B > 0$, then $U_{1,n} \in H^{s+2}([0, d] \times [b, a])$ for $l = 1, 2$, and

\[ \|U_{1,n}\|_{H^{s+2}} + \|U_{2,n}\|_{H^{s+2}} \leq KB^n, \]

for some universal constant $K$.

**Proof.** We proceed by induction, and at order $n = 0$ Theorem 4.3 guarantees a unique solution such that

\[ \|U_{1,0}\|_{H^{s+2}} + \|U_{2,0}\|_{H^{s+2}} \leq K_c \|\xi_{1,0}\|_{H^{s+3/2}}, \]

so we choose $K \geq K_c \|\xi_{1,0}\|_{H^{s+3/2}}$. We now assume (5.2) holds for all $n \leq \bar{n}$ and from Theorem 4.3 we deduce that

\[ \|U_{1,n}\|_{H^{s+2}} + \|U_{2,n}\|_{H^{s+2}} \leq C_1 \left( \|R_{1,n}\|_{H^s} + \|R_{2,n}\|_{H^s} \right) \]
\[ + \|J_{1,n}\|_{H^{s+1/2}} + \|J_{2,n}\|_{H^{s+1/2}} + \|J_{2,n}\|_{H^{s+1/2}} \right) + C_2 \|\xi_{1,n}\|_{H^{s+3/2}}. \]

Appealing to Lemmas 5.1 and 5.2 we find that

\[ \|U_{1,n}\|_{H^{s+2}} + \|U_{2,n}\|_{H^{s+2}} \leq 5C_1 K \xi \left( B^{n-1} - B^{n-2} - B^{n-3} \right) + C_2 K \xi B^n \]

Now, upon choosing $K > C_2 K_\xi$ and

\[ B > \max \left\{ B_\xi, 5C_1 \tilde{C}, (5C_1 \tilde{C})^{1/2}, (5C_1 \tilde{C})^{1/3} \right\}, \]

the theorem follows.

**6. Convergence Analysis.** We are now in a position to conduct a numerical analysis of our TFE approach. We recall the TFE recursions (3.2) and note that, in practice, we make use of the Floquet series representation, (3.3), and focus our attention on the reduced problem (3.4). We further specialize by splitting this into two: A homogeneous Helmholtz problem with inhomogeneous coupling ($\tilde{\xi}^{(p)}_{j,n} \neq 0$),
for any \( v \in \Lambda = (\text{interval of variables where } n \geq 0) \), see (B.2). Clearly, the solution of (3.4) is the sum of the solutions of these two problems and, in practical numerical implementations, we need only solve the latter as (B.1) can be solved explicitly via Separation of Variables, e.g., [11, 14, 13]. For this reason we focus upon (B.2) and, for simplicity, we suppress the index \( n \). The weak form of this boundary value problem is:

\[
\text{Find } \tilde{U}^{(p)} \in H^1(b, a) \text{ such that }
\]

\[
B(\tilde{U}^{(p)}, \varphi) = R(\varphi), \quad \forall \varphi \in H^1(b, a),
\]

where

\[
B(\tilde{U}^{(p)}, \varphi) := -i\gamma_{1,p} \tilde{U}_1^{(p)}(a) \varphi_1(a) - i\gamma_{2,p} \tilde{U}_2^{(p)}(b) \varphi_2(b)
\]

\[
+ \int_b^a \partial_z \tilde{U}^{(p)} \partial_z \varphi \ dz - \gamma_p^2 \int_b^a \tilde{U}^{(p)} \varphi \ dz,
\]

\[
R(\varphi) := \tilde{\gamma}_1^{(p)} \varphi(a) - \tilde{\gamma}_2^{(p)} \varphi(b) + \int_b^a (-\tilde{R}^{(p)}) \varphi \ dz.
\]

For our numerical analysis we define the discrete function space

\[
X_{M,p} = \text{span} \{ u \in C(b, a) \mid u|_{(b,0)}, u|_{(b,0)} \in P_M,
\]

\[
(\partial_z u - i\gamma_{1,p} u)(a) = \tilde{\gamma}_1^{(p)}, \ (\partial_z u + i\gamma_{2,p} u)(b) = \tilde{\gamma}_2^{(p)} \},
\]

where \( P_M \) is the space of all complex valued polynomials of degree less than or equal to \( M \). The Legendre–Galerkin approximation of (6.1) is as follows:

\[
\text{Find } \tilde{U}^{(p),M} \in X_{M,p} \text{ such that }
\]

\[
B(\tilde{U}^{(p),M}, \varphi_M) = R(\varphi_M), \quad \forall \varphi_M \in X_{M,p}.
\]

To prove the main theorem of this section the following interpolation result [28] is required for the projection \( \Pi^1_M \) from \( H^1(b, a) \) to \( P_M \) subject to the boundary conditions of the space \( X_{M,p} \).

**Lemma 6.1.** There exists a mapping \( \Pi^1_M : H^1(b, a) \rightarrow X_{M,p} \) such that

\[
(\partial_z (\Pi^1_M V - V), \partial_z \varphi_M) = 0, \quad \forall \varphi_M \in X_{M,p}.
\]

Moreover, for \( 1 \leq l \leq M + 1 \) we have

\[
\left\| \Pi^1_M V - V \right\|_{H^l} \lesssim \sqrt{\frac{(M - l + 1)!}{M!}} (M + 1)^{(l + 1)/2} \left\| \partial_z^l V \right\|_{L^2},
\]

where \( \mu = 0, 1 \).

**Proof.** We prove this lemma for \( V_1 := V|_{(b,0)} \). By the straightforward change of variables \( x = \frac{2z}{a} - 1 \), the domain of \( V_1 \in H^1(b, a) \) can be transformed to the interval \((-1, 1)\). Thus we establish the result for a real valued function \( v(x) \) on \( \Lambda = (-1, 1) \). Let \( \Pi^1_M \) be the \( H^1 \)-orthogonal projection operator onto \( P_M \times P_M \) and, for any \( v \in H^1(\Lambda) \), we define \( v_*(x) \) by

\[
v(x) = v_*(x) + \left( \frac{1 + x}{2} \right) v(1) + \left( \frac{1 - x}{2} \right) v(-1), \quad v_*(x) \in H^1_0(\Lambda).
\]
Similarly, we define $\varphi_*(x)$ by
\[
\varphi(x) = \varphi_*(x) + \left(\frac{1+x}{2}\right)v(1) + \left(\frac{1-x}{2}\right)v(-1), \quad \varphi_*(x) \in H_0^1(\Lambda),
\]
for any $\varphi \in P_M$. Regarding
\[
\psi \Pi_M^1 v(x) := \Pi_M^{1,0} \varphi_*(x) + \left(\frac{1+x}{2}\right)v(1) + \left(\frac{1-x}{2}\right)v(-1),
\]
we observe that
\[
\left\langle \partial_x(\psi \Pi_M^1 v - v), \partial_x \varphi \right\rangle = \left\langle \partial_x(\Pi_M^{1,0} \varphi_*(x) - v(x)), \partial_x \varphi \right\rangle
\]
\[
+ \left(\frac{v(1)}{2} - \frac{v(-1)}{2}\right) \int_{-1}^1 \partial_x(\Pi_M^{1,0} \varphi_*(x) - v(x)) \, dx
\]
\[
= \left\langle \partial_x(\Pi_M^{1,0} \varphi_*(x) - v(x)), \partial_x \varphi \right\rangle = 0,
\]
for $\varphi \in P_M$. By Theorem 3.39 in [28] we find
\[
\left\| \psi \Pi_M^1 v(x) - v \right\|_{H^1} \leq \left\| \Pi_M^{1,0} \varphi_*(x) - v(x) \right\|_{H^1}
\]
\[
\lesssim \sqrt{\frac{(M - l + 1)!}{M!}} (M + l)^{(1-l)/2} \left\| \partial_x^l \varphi \right\|_{L^2}.
\]
For $l = 1$, by the Poincaré inequality, we derive that
\[
\left\| \partial_x v \right\|_{L^2} \leq \left\| \partial_x v \right\|_{L^2} + c (|v(1)| + |v(-1)|) \leq c \left\| \partial_x v \right\|_{L^2}.
\]
From this the lemma follows. \hfill \Box

Now, we are ready to prove the convergence theorem.

**Theorem 6.2.** Let $\bar{U}(p)$ and $\bar{U}(p,M)$ be the solutions of (6.1) and (6.2), respectively. Then, for $1 \leq l \leq M + 1$, we have
\[
\left\| \bar{U}(p) - \bar{U}(p,M) \right\|_{H^l} + |\gamma| \left\| \bar{U}(p) - \bar{U}(p,M) \right\|_{L^2}
\]
\[
\leq (1 + \gamma^2 M^{-1}) \sqrt{\frac{(M - l + 1)!}{M!}} (M + l)^{(1-l)/2} \left\| \partial_x^l \bar{U}(p) \right\|_{L^2}.
\]

**Proof.** Let
\[
e_M := \bar{U}(p,M) - o \Pi_M^1 \bar{U}(p), \quad \tilde{e}_M = \bar{U}(p) - o \Pi_M^1 \bar{U}(p).
\]
For $\varphi_M \in X_{M,p}$, using (6.1) and (6.2), we find
\[
B(\bar{U}(p) - \bar{U}(p,M), \varphi_M) = 0.
\]
Using Lemma 6.1, we obtain
\[
B(e_M, \varphi_M) = B(\bar{U}(p,M) - \bar{U}(p) + \bar{U}(p) - o \Pi_M^1 \bar{U}(p), \varphi_M)
\]
\[
= B(\bar{U}(p) - o \Pi_M^1 \bar{U}(p), \varphi_M)
\]
\[
= -\gamma^2 \left( e_M, \varphi_M \right) - i \gamma_1 \bar{e}_M \varphi_M(a) \bar{\varphi}_M(a) - i \gamma_2 \bar{e}_M \varphi_M(b) \varphi_M(b).
\]
In view of (6.3), we rewrite (6.1) by replacing \( \{ e_M, -i\gamma_{1,p}e_M(a), i\gamma_{2,p}e_M(b), \gamma_p^2 e_M \} \), respectively. Then, by the regularity result (B.7) from Appendix B we obtain that
\[
\| e_M \|^2_{H^1} + \gamma_p^2 \| e_M \|^2_{L^2} \lesssim \gamma_p^4 \| \tilde{e}_M \|^2_{L^2} + \gamma_{1,p}^2 |\tilde{e}_M(a)|^2 + \gamma_{2,p}^2 |\tilde{e}_M(b)|^2.
\]

By the Gagliardo–Nirenberg interpolation inequality [29] and Lemma 6.1 we find
\[
|\tilde{e}(\pm 1)| \lesssim \| \tilde{e}_M \|_{L^2}^{1/2} \| \tilde{e}_M \|^2_{H^1} \sim \sqrt{\frac{(M - l + 1)!}{M!}} (M + l)^{-l/2} \| \partial_x^l U(p) \|_{L^2}.
\]

Using Lemma 6.1 again, we deduce that
\[
\left\| U^{(p)} - U^{(p),M} \right\|_{H^1} + |\gamma_p| \left\| U^{(p)} - U^{(p),M} \right\|_{L^2} \lesssim (\| e_M \|_{H^1} + |\gamma_p| \| e_M \|_{L^2} + \| \tilde{e}_M \|_{H^1} + |\gamma_p| \| \tilde{e}_M \|_{L^2}) \lesssim \left(1 + \gamma_p^2 M^{-1} + |\gamma_p| M^{-1/2}\right) \sqrt{\frac{(M - l + 1)!}{M!}} (M + l)^{(1-l)/2} \| \partial_x^l U(p) \|_{L^2}.
\]

We now reintroduce the index \( n \) and let
\[
U_{l,n}^{(p),M}(x, z) := \sum_{p=-P}^{P} \tilde{U}_{l,n}^{(p),M}(z)e^{i\alpha_p x}, \quad l = 0, 1,
\]
be the Fourier–Legendre approximation of the solution \( U_{l,n} \) of (3.2). Using the same argument as in Theorem 3.3 in [24], we can prove the following estimate.

**Theorem 6.3.** For any integer \( r \geq 1 \), if \( U \in H^r \) then
\[
\left\| \nabla (U_{l,n} - U_{l,n}^{(p),M}) \right\|_{L^2} + k_l \left\| U_{l,n} - U_{l,n}^{(p),M} \right\|_{L^2} \lesssim \left(P^{1-r} + (1 + k_l^2 M^{-1}) \sqrt{\frac{(M - r + 1)!}{M!}} (M + r)^{(1-r)/2}\right) \| U_{l,n} \|_{H^r}.
\]

Finally if we choose
\[
U_{l,n}^{(p),M}(x, z) := \sum_{n=0}^{N} U_{l,n}^{(p),M}(x, z)e^{i\alpha_p x},
\]
as our approximation to the solution \( U_m \) of (3.2), then, using Theorem 6.2 and Theorem 2.1 of [24], we have the final result.

**Theorem 6.4.** For any integer \( r \geq 2 \) if \( f \in C^r([0, d]), \xi_1 \in H^{r-1/2}([0, d]), \) and \( \xi_2 \in H^{r-3/2}([0, d]) \) then we have, for \( l = 1, 2, \)
\[
\left\| \nabla (U_l - U_{l,n}^{(p),M}) \right\|_{L^2} + k_l \left\| U_l - U_{l,n}^{(p),M} \right\|_{L^2} \lesssim (B\varepsilon)^{N+1} + \| \xi_2 \|_{H^{r-3/2}} \lesssim (P^{1-r} + (1 + k_l^2 M^{-1}) \sqrt{\frac{(M - r + 1)!}{M!}} (M + r)^{(1-r)/2}) \| \xi_1 \|_{H^{r-1/2}}.
\]
7. Conclusions. In this paper we have provided a rigorous numerical analysis of a High–Order Perturbation of Surfaces (HOPS) algorithm for electromagnetic scattering. Introducing Dirichlet–Neumann operators at artificial boundaries placed above the top and below the bottom of the structure, we equivalently reformulated the governing Helmholtz equations for the doubly layered medium on a bounded domain. Using a suitable change of variables, the governing equations on a separable geometry with flat interfaces were derived. Introducing boundary perturbations, we described the scattered field in a Taylor series, more precisely, we derived a sequence of linear boundary value problems to be solved at each perturbation order resulting in the Transformed Field Expansions (TFE) algorithm. Our approach to establishing the convergence and accuracy of the TFE methodology is to combine analyticity theorems with results on Legendre–Galerkin methods. Our developments clearly point towards several extensions of great importance. In particular, our approach must be generalized to the three dimensional vector wave equations of electromagnetics and linear elastodynamics. These extensions are not straightforward as more complicated boundary conditions between layers are required. Hence the algorithmic differences will be significant and we will describe them in a future publication.

Appendix A. Proof of the Elliptic Estimate: Theorem 4.3.

To begin our proof of Theorem 4.3 we state two classic results [22] on solutions of Helmholtz problems on each of the two layers separately.

THEOREM A.1. Given any integer \( s \geq 0 \), if \( F_1 \in H^{s+1}([0,a]) \), \( U \in H^{s+1/2}([0,d]) \), then there exists a unique solution \( u \in H^{s+1}([0,a]) \) of

\[
\begin{align*}
\Delta_1 u + k_1^2 u &= F_1, & 0 < z_1 < a, \\
u &= U, & z_1 = 0, \\
\partial_{z_1} u - T_1[u] &= K_1, & z_1 = a,
\end{align*}
\]

such that

\[ \|u\|_{H^{s+1}} \leq C_u \left\{ \|F_1\|_{H^{s+1/2}} + \|U\|_{H^{s+1/2}} + \|K_1\|_{H^{s+1/2}} \right\}. \]

In addition, if \( \hat{U} = [-\partial_{z_1} u]_{z_1=0} \), and we define the Dirichlet–Neumann Operator (DNO)

\[ \mathcal{G} : (U, K_1, F_1) \rightarrow \hat{U}, \quad \mathcal{G}[U, K_1, F_1] = G^{(0)}[U] + G^{(a)}[K_1] + G^{([0,a])}[F_1], \]

then

\[ \begin{align*}
\|G^{(0)}[U]\|_{H^{s-1/2}} &\leq C_{G^{(0)}} \|U\|_{H^{s+1/2}}, \\
\|G^{(a)}[K_1]\|_{H^{s-1/2}} &\leq C_{G^{(a)}} \|K_1\|_{H^{s-1/2}}, \\
\|G^{([0,a])}[F_1]\|_{H^{s-1/2}} &\leq C_{G^{([0,a])}} \|F_1\|_{H^{s-1}}.
\end{align*} \]

Proof. For clarity of presentation we drop the “1” subscript on all variables. Due to the quasiperiodic boundary conditions we posit expansions

\[ \{u, F\}(x, z) = \sum_{p=-\infty}^{\infty} \{\hat{u}_p, \hat{F}_p\}(z) e^{i\alpha_p x}, \quad \{U, K\}(x) = \sum_{p=-\infty}^{\infty} \{\hat{U}_p, \hat{K}_p\} e^{i\alpha_p x}, \]
and (A.1) delivers the two–point boundary value problem
\[ \partial_z^2 \hat{u}_p + \gamma_p^2 \hat{u}_p = \hat{F}_p, \quad 0 < z < a, \]
\[ \hat{u}_p(0) = \hat{U}_p, \]
\[ \partial_z \hat{u}_p(a) - (i\gamma_p) \hat{u}_p(a) = \hat{K}_p, \]
where
\[ \gamma_p = \begin{cases} \gamma'_p := \sqrt{k^2 - \alpha_p^2}, & \alpha_p^2 < k^2, \\ 0, & \alpha_p^2 = k^2, \\ i\gamma''_p := i\sqrt{\alpha_p^2 - k^2}, & \alpha_p^2 > k^2, \end{cases} \]
\[ \gamma'_p, \gamma''_p \in \mathbb{R}, \quad \gamma'_p, \gamma''_p > 0. \]

It is not difficult to show that the unique solution of this problem is given by
\[ \hat{u}_p(z) = \hat{U}_p \Phi_0(z; p) + \hat{K}_p e^{i\gamma_p \alpha} \Phi_a(z; p) - I_0[\hat{F}_p](z) - I_a[\hat{F}_p](z), \]
where
\[ \Phi_0(z; p) = e^{i\gamma_p z} := \begin{cases} e^{i\gamma'_p z}, & \alpha_p^2 < k^2, \\ 1, & \alpha_p^2 = k^2, \\ e^{-i\gamma''_p z}, & \alpha_p^2 > k^2, \end{cases} \]
and
\[ \Phi_a(z; p) = \frac{\sinh(\gamma_p z)}{\gamma_p} := \begin{cases} \frac{\sin(\gamma'_p z)}{\gamma'_p}, & \alpha_p^2 < k^2, \\ z, & \alpha_p^2 = k^2, \\ \frac{\sinh(\gamma''_p z)}{\gamma''_p}, & \alpha_p^2 > k^2, \end{cases} \]
and
\[ I_0[\hat{F}_p](z) := \int_0^z \Phi_0(z; s) \Phi_a(s; p) \hat{F}_p(s) \, ds, \]
\[ I_a[\hat{F}_p](z) := \int_z^a \Phi_0(s; p) \Phi_a(z; s) \hat{F}_p(s) \, ds. \]

It is straightforward to compute that
\[ \partial_z I_0[\hat{F}_p](z) = \Phi_0(z; p) \Phi_a(z; p) \hat{F}_p(z) + \int_0^z (\partial_z \Phi_0(z; p)) \Phi_a(s; p) \hat{F}_p(s) \, ds, \]
\[ \partial_z I_a[\hat{F}_p](z) = -\Phi_0(z; p) \Phi_a(z; p) \hat{F}_p(z) + \int_z^a \Phi_0(s; p) (\partial_z \Phi_a(z; p)) \hat{F}_p(s) \, ds. \]

Noting the cancellation in the sum of the terms \( \partial_z I_0 \) and \( \partial_z I_a \) we realize
\[ \partial_z \hat{u}_p(z) = \hat{U}_p \partial_z \Phi_0(z; p) + \hat{K}_p e^{i\gamma_p \alpha} \partial_z \Phi_a(z; p) - \hat{I}_0[\hat{F}_p](z) - \hat{I}_a[\hat{F}_p](z), \]
where
\[ \hat{I}_0[\hat{F}_p](z) := \int_0^z (\partial_z \Phi_0(z; p)) \Phi_a(s; p) \hat{F}_p(s) \, ds, \]
\[ \hat{I}_a[\hat{F}_p](z) := \int_z^a \Phi_0(s; p) (\partial_z \Phi_a(z; p)) \hat{F}_p(s) \, ds. \]
If we evaluate this at \( z = 0 \) we find

\[
-\partial_z \hat{u}_p(0) = -\hat{U}_p \partial_z \Phi_0(0; p) - \hat{K}_p e^{i\gamma_{p-a}} \partial_z \Phi_a(0; p) + \tilde{I}_0[\hat{F}_p](0) + \tilde{I}_a[\hat{F}_p](0)
\]

\[
= -\hat{U}_p (i\gamma_p) - \hat{K}_p e^{i\gamma_{p-a}} + \int_0^a e^{i\gamma_{p-s}} \cosh(\gamma_{p-s}) \hat{F}_p(s) \, ds.
\]

With these it is easy to see that

\[
G^{(0)}[U] = -\sum_{p=-\infty}^{\infty} (\partial_z \Phi_0)(0; p) \hat{U}_p e^{i\alpha_{p-x}} = \sum_{p=-\infty}^{\infty} (-i\gamma_p) \hat{U}_p e^{i\alpha_{p-x}},
\]

and

\[
G^{(a)}[K] = -\sum_{p=-\infty}^{\infty} e^{i\gamma_{p-a}} (\partial_z \Phi_a)(0; p) \hat{K}_p e^{i\alpha_{p-x}} = \sum_{p=-\infty}^{\infty} (-e^{i\gamma_{p-a}}) \hat{K}_p e^{i\alpha_{p-x}},
\]

and

\[
G^{(0,a)}[F] = \sum_{p=-\infty}^{\infty} \int_0^a \left( e^{i\gamma_{p-s}} \cosh(\gamma_{p-s}) \hat{F}_p(s) \, ds \right) e^{i\alpha_{p-x}}.
\]

Regarding the estimates, these follow from the asymptotic estimates of \( \| \Phi_0 \|_{L^2(dz)} \), \( \| \Phi_a \|_{L^2(dz)} \), \( \| I_0[F] \|_{L^2(dz)} \), \( \| I_a[F] \|_{L^2(dz)} \), \( \| \tilde{I}_0[F] \|_{L^2(dz)} \), and \( \| \tilde{I}_a[F] \|_{L^2(dz)} \).

The analogue in the lower layer is the following result. It is established in an almost identical fashion as Theorem A.1.

**Theorem A.2.** Given any integer \( s \geq 0 \), if \( F_2 \in H^{s-1}([0, d] \times [-a, 0]), V \in H^{s+1/2}([0, d]), K_2 \in H^{s+1/2}([0, d]), \) then there exists a unique solution \( v \in H^{s+1}([0, d] \times [b, 0]) \) of

(A.2a) \[ \Delta_2 v + k_2^2 v = F_2, \quad b < z_2 < 0, \]

(A.2b) \[ v = V, \quad z_2 = 0, \]

(A.2c) \[ \partial_{z_2} v - T_2[v] = K_2, \quad z_2 = b, \]

such that

\[ \| v \|_{H^{s+1}} \leq C_v \left( \| F_2 \|_{H^{s-1}} + \| V \|_{H^{s+1/2}} + \| K_2 \|_{H^{s+1/2}} \right). \]

In addition, if \( \tilde{V} = [\partial_{z_2} v]_{z_2=0} \), and we define the Dirichlet–Neumann Operator (DNO)

\[ \mathcal{J} : (V, K_2, F_2) \rightarrow \tilde{V}, \quad \mathcal{J}[V, K_2, F_2] = J^{(0)}[V] + J^{(b)}[K_2] + J^{(b,0)}[F_2], \]

then

\[ \| J^{(0)}[V] \|_{H^{s+1/2}} \leq C_{J^{(0)}} \| V \|_{H^{s+1/2}}, \]

\[ \| J^{(b)}[K_2] \|_{H^{s-1/2}} \leq C_{J^{(b)}} \| K_2 \|_{H^{s-1/2}}, \]

\[ \| J^{(b,0)}[F_2] \|_{H^{s-1/2}} \leq C_{J^{(b,0)}} \| F_2 \|_{H^{s-1}}. \]
In addition we require the following result on the boundary conditions which couple \( u \) and \( v \) at the interface \( z_1 = z_2 = 0 \).

**Theorem A.3.** Given any integer \( s \geq 0 \), if \( Q \in H^{s+1/2}([0,d]) \) and \( R \in H^{s-1/2}([0,d]) \), then there exists a unique solution pair \( U, V \in H^{s+1/2}([0,d]) \) of

\[
\begin{align*}
(A.3a) & \quad U - V = Q, \\
(A.3b) & \quad G^{(0)}[U] + J^{(0)}[V] = R,
\end{align*}
\]

such that

\[
\max \{ \|U\|_{H^{s+1/2}}, \|V\|_{H^{s+1/2}} \} \leq C_0 \{ \|Q\|_{H^{s+1/2}} + \|R\|_{H^{s-1/2}} \}.
\]

**Proof.** The result follows simply from the well–known expressions for the flat–interface DNOs

\[
\begin{align*}
G^{(0)}[U] &= G^{(0)} \left[ \sum_{p=-\infty}^{\infty} \hat{U}_p e^{i\alpha_p x} \right] = \sum_{p=-\infty}^{\infty} (-i\gamma_1,p)\hat{U}_p e^{i\alpha_p x}, \\
J^{(0)}[U] &= J^{(0)} \left[ \sum_{p=-\infty}^{\infty} \hat{V}_p e^{i\alpha_p x} \right] = \sum_{p=-\infty}^{\infty} (-i\gamma_2,p)\hat{V}_p e^{i\alpha_p x},
\end{align*}
\]

so that the governing equations become

\[
\begin{pmatrix}
1 \\
(-i\gamma_1,p) \\
(-i\gamma_2,p)
\end{pmatrix}
\begin{pmatrix}
\hat{U}_p \\
\hat{V}_p
\end{pmatrix} = \begin{pmatrix}
\hat{Q}_p \\
\hat{R}_p
\end{pmatrix}, \quad \forall \ p \in \mathbb{Z}.
\]

These are readily solved

\[
\begin{pmatrix}
\hat{U}_p \\
\hat{V}_p
\end{pmatrix} = \frac{1}{(i\gamma_1,p) + (i\gamma_2,p)} \begin{pmatrix}
(-i\gamma_2,p) \\
(i\gamma_1,p)
\end{pmatrix} \begin{pmatrix}
1 \\
1
\end{pmatrix} \begin{pmatrix}
\hat{Q}_p \\
\hat{R}_p
\end{pmatrix},
\]

and the \( \{\hat{U}_p, \hat{V}_p\} \) have the right decay to verify the conclusions of the theorem.

We can now proceed to our principal result, Theorem 4.3.

**Proof.** [Theorem 4.3] We begin by rewriting (4.3) as

\[
\begin{align*}
(A.4a) & \quad \Delta_1 u + k_1^2 u = F_1, & 0 < z_1 < a, \\
(A.4b) & \quad u = U, & z_1 = 0, \\
(A.4c) & \quad \partial_{z_1} u - T_1 [u] = K_1, & z_1 = a, \\
(A.4d) & \quad \Delta_2 v + k_2^2 v = F_2, & b < z_2 < 0, \\
(A.4e) & \quad v = V, & z_2 = 0, \\
(A.4f) & \quad \partial_{z_2} v - T_2 [v] = K_2, & z_2 = b, \\
(A.4g) & \quad U - V = \xi, & z_1 = z_2 = 0, \\
(A.4h) & \quad \hat{U} + \hat{V} = -\nu, & z_1 = z_2 = 0.
\end{align*}
\]

From Theorem A.1 we see that, provided that \( F_1 \in H^{s-1}, K_1 \in H^{s-1/2}, U \in H^{s+1/2} \) then (A.4a)–(A.4c) delivers a unique solution \( u \in H^{s+1} \) as desired. The functions \( F_1 \) and \( K_1 \) in the correct spaces are provided so we merely need show that \( U \) is in \( H^{s+1/2} \). In a similar fashion, Theorem A.2 guarantees that, if \( F_2 \in H^{s-1}, K_2 \in H^{s-1/2} \),
V ∈ H^{s+1/2} then (A.4d)–(A.4f) provides a unique solution v ∈ H^{s+1}. As before, the functions F_2 and K_2 in the correct spaces are provided so we are left to show that V is in H^{s+1/2}.

Thus, all that remains is to consider (A.4g)–(A.4h) which we write in terms of DNOs as

\[ U - V = \xi, \]

\[
\left(G^{(0)}[U] + G^{(a)}[K_1] + G^{(b,0)}[F_1]\right) + \left(J^{(0)}[V] + J^{(b)}[K_2] + J^{(b,0)}[F_2]\right) = -\nu,
\]
or

\[ U - V = \xi, \]

\[
G^{(0)}[U] + J^{(0)}[V] = -\nu - G^{(a)}[K_1] - G^{(b,0)}[F_1] - J^{(b)}[K_2] - J^{(b,0)}[F_2].
\]

Theorem A.3 delivers the required solutions U, V ∈ H^{1/2} provided that

\[ Q = \xi ∈ H^{s+1/2}, \]

\[ R = -\nu - G^{(a)}[K_1] - G^{(b,0)}[F_1] - J^{(b)}[K_2] - J^{(b,0)}[F_2] ∈ H^{s-1/2}, \]

both of which are true from (i.) our hypotheses on ξ, ν, K_1, K_2, F_1, and F_2; and (ii.) the mapping properties of G^{(a)}, G^{(b,0)}, J^{(b)}, and J^{(b,0)} established in Theorems A.1 and A.2.

Appendix B. Regularity of Solutions of the Weak Formulation.

We now produce an elliptic regularity theory for solutions of the boundary value problem (3.4). (For the sake of simplicity, we drop the indices \{p, n\}.) Noting that \(\gamma^2_{l,p} = k_l^2 - \alpha_p^2\) we split (3.4) into two BVPs: One with inhomogeneous coupling (which, due to the homogeneous Helmholtz equations, we can solve explicitly with Fourier analysis),

(B.1a) \(\partial_{z_2}^2 \bar{U}_1 + \gamma^2_{1,1} \bar{U}_1 = 0,\) \(0 < z_1 < a,\)

(B.1b) \(\partial_{z_2}^2 \bar{U}_2 + \gamma^2_{2,2} \bar{U}_2 = 0,\) \(b < z_2 < a,\)

(B.1c) \(\bar{U}_1 - \bar{U}_2 = \bar{\xi}_1,\) \(z_1 = z_2 = 0,\)

(B.1d) \(\partial_{z_1} \bar{U}_1 - \partial_{z_2} \bar{U}_2 = \bar{\xi}_2,\) \(z_1 = z_2 = 0,\)

(B.1e) \(\partial_{z_2} \bar{U}_1 - i\gamma_1 \bar{U}_1 = 0,\) \(z_1 = a,\)

(B.1f) \(\partial_{z_2} \bar{U}_2 + i\gamma_2 \bar{U}_2 = 0,\) \(z_2 = b,\)

and one with homogeneous coupling (but inhomogeneous Helmholtz equations)

(B.2a) \(\partial_{z_1}^2 \bar{U}_1 + \gamma^2_{1,1} \bar{U}_1 = \bar{R}_1,\) \(0 < z_1 < a,\)

(B.2b) \(\partial_{z_2}^2 \bar{U}_2 + \gamma^2_{2,2} \bar{U}_2 = \bar{R}_2,\) \(b < z_2 < 0,\)

(B.2c) \(\bar{U}_1 - \bar{U}_2 = 0,\) \(z_1 = z_2 = 0,\)

(B.2d) \(\partial_{z_1} \bar{U}_1 - \partial_{z_2} \bar{U}_2 = 0,\) \(z_1 = z_2 = 0,\)

(B.2e) \(\partial_{z_1} \bar{U}_1 - i\gamma_1 \bar{U}_1 = \bar{J}_1,\) \(z_1 = a,\)

(B.2f) \(\partial_{z_2} \bar{U}_2 + i\gamma_2 \bar{U}_2 = \bar{J}_2,\) \(z_2 = b.\)
To study the regularity of solutions of (B.2) we find the variational formulation as in [29, 30],

\[
\begin{align*}
(B.3) \quad \int_\beta^a \partial_\overline{z} \partial_\overline{\varphi} - \gamma^2 \int_\beta^a \overline{\varphi} = i \gamma_1 \overline{U_1(a)} \varphi_1(a) + i \gamma_2 \overline{U_2(b)} \varphi_2(b) \\
= \overline{\tilde{J}_1 \varphi_1(a) - \tilde{J}_2 \varphi_2(b)} + \int_\beta^a (\overline{-R}) \varphi, 
\end{align*}
\]

take \( \varphi = \overline{U} \), and consider the imaginary and real parts, respectively. For the imaginary part we find

\[-\gamma_1 |\overline{U_1(a)}|^2 - \gamma_2 |\overline{U_2(b)}|^2 = \text{Im} \left\{ (\overline{-R}, \overline{U}) \right\} + \text{Im} \left\{ \overline{\tilde{J}_1 \overline{U_1(a)}} \right\} - \text{Im} \left\{ \overline{\tilde{J}_2 \overline{U_2(b)}} \right\},
\]

With this we estimate

\[
\begin{align*}
\gamma_1 |\overline{U_1(a)}|^2 + \gamma_2 |\overline{U_2(b)}|^2 &\leq \frac{\kappa_M \delta_1}{2} \left\| \overline{U} \right\|^2_{L^2} + \frac{1}{2 \delta_1 \kappa_M} \left\| \overline{R} \right\|^2_{L^2} \\
&+ \frac{\gamma_1}{2} |\overline{U_1(a)}|^2 + \frac{\gamma_2}{2} |\overline{U_2(b)}|^2 + \frac{1}{2 \gamma_1} |\overline{\tilde{J}_1}|^2 + \frac{1}{2 \gamma_2} |\overline{\tilde{J}_2}|^2,
\end{align*}
\]

where \( \kappa_M := \max(|\gamma_1|, |\gamma_2|) \) and \( \delta_1 > 0 \) will be chosen later. For the real part we deduce that

\[
\left\| \partial_\overline{U} \right\|^2_{L^2} - \gamma^2 \left\| \overline{U} \right\|^2_{L^2} = \text{Re} \left\{ (\overline{-R}, \overline{U}) \right\} + \text{Re} \left\{ \overline{\tilde{J}_1 \overline{U_1(a)}} \right\} - \text{Re} \left\{ \overline{\tilde{J}_2 \overline{U_2(b)}} \right\},
\]

and this implies

\[
\begin{align*}
\left\| \partial_\overline{U} \right\|^2_{L^2} &\leq \kappa_M \left\| \overline{U} \right\|^2_{L^2} + \delta_2 \kappa_M \left| \overline{U_1(a)} \right|^2 + \delta_2 \kappa_M \left| \overline{U_2(b)} \right|^2 \\
&+ \frac{1}{4 \delta_2 \kappa_M^2} |\overline{\tilde{J}_1}|^2 + \frac{1}{4 \delta_2 \kappa_M^2} |\overline{\tilde{J}_2}|^2 + \frac{\delta_3 \kappa_M^2}{2} \left\| \overline{U} \right\|^2_{L^2} + \frac{1}{2 \delta_3 \kappa_M^2} \left\| \overline{R} \right\|^2_{L^2},
\end{align*}
\]

where \( \delta_2, \delta_3 > 0 \) will also be chosen later. Using (B.4) we deduce that

\[
\kappa_m \left( \left| \overline{U_1(a)} \right|^2 + \left| \overline{U_2(b)} \right|^2 \right) \leq \kappa_M \delta_1 \left\| \overline{U} \right\|^2_{L^2} + \frac{1}{\delta_1 \kappa_M} \left\| \overline{R} \right\|^2_{L^2} + \frac{1}{\gamma_1} |\overline{\tilde{J}_1}|^2 + \frac{1}{\gamma_2} |\overline{\tilde{J}_2}|^2,
\]

where \( \kappa_m := \min(|\gamma_1|, |\gamma_2|) \), and this implies

\[
\begin{align*}
\left| \overline{U_1(a)} \right|^2 + \left| \overline{U_2(b)} \right|^2 &\leq \tau \delta_1 \left\| \overline{U} \right\|^2_{L^2} + \frac{1}{\delta_1 \kappa_M \kappa_m} \left\| \overline{R} \right\|^2_{L^2} + \frac{1}{\gamma_1 \kappa_m} |\overline{\tilde{J}_1}|^2 + \frac{1}{\gamma_2 \kappa_m} |\overline{\tilde{J}_2}|^2,
\end{align*}
\]

where \( \tau = \kappa_M / \kappa_m \). Using (B.6) and (B.5), we derive that

\[
\left\| \partial_\overline{U} \right\|^2_{L^2} \leq \left( \kappa_M^2 + \delta_2 \kappa_M \tau \delta_1 + \frac{\delta_3 \kappa_M^2}{2} \right) \left\| \overline{U} \right\|^2_{L^2} + \left( \frac{\delta_1 \tau^2}{2} + \frac{1}{2 \delta_3 \kappa_M^2} \right) \left\| \overline{R} \right\|^2_{L^2}
\]

\[
+ \left( \delta_2 \tau^2 + \frac{1}{r \delta_2 \kappa_M^2} \right) \left( |\overline{\tilde{J}_1}|^2 + |\overline{\tilde{J}_2}|^2 \right).
\]

Setting \( \delta_2 = \delta_3 / (2 \delta_1 \tau) \), we obtain

\[
\left\| \partial_\overline{U} \right\|^2_{L^2} \leq \left( \kappa_M^2 + \delta_3 \kappa_M \right) \left\| \overline{U} \right\|^2_{L^2} + \left( \frac{\delta_1 \tau^2}{2 \delta_1^2} + \frac{1}{2 \delta_3 \kappa_M^2} \right) \left\| \overline{R} \right\|^2_{L^2}
\]

\[
+ \left( \delta_2 \tau^2 + \frac{1}{r \delta_2 \kappa_M^2} \right) \left( |\overline{\tilde{J}_1}|^2 + |\overline{\tilde{J}_2}|^2 \right).
\]
Regarding (B.3) again, we now take the test function
\[
\varphi = 2z\partial_z \bar{U} = \begin{cases} 
2z\partial_z \bar{U}_1, & z \in (0, a) =: I_1, \\
2z\partial_z \bar{U}_2, & z \in (b, 0) =: I_2.
\end{cases}
\]
Then, the weak form (B.3) becomes
\[
\|\partial_z \bar{U}_1\|_{L^2(I_1)}^2 + \|\partial_z \bar{U}_2\|_{L^2(I_2)}^2 + a \left|\partial_z \bar{U}_1(a)\right|^2 - b \left|\partial_z \bar{U}_2(b)\right|^2 \\
+ \gamma_1^2 \left\|\bar{U}_1\right\|_{L^2(I_1)}^2 + \gamma_2^2 \left\|\bar{U}_2\right\|_{L^2(I_2)}^2 \\
\leq \kappa_M^2 M \left(\left|\bar{U}_1(a)\right|^2 + \left|\bar{U}_2(b)\right|^2\right) + \frac{1}{2} \left\|\partial_z \bar{U}\right\|_{L^2}^2 + 8M^2 \left\|\hat{R}\right\|_{L^2}^2 \\
+ \left(m + \frac{8}{m} 4M^2 \kappa_M^2\right) \left(\left|\partial_z \bar{U}_1(a)\right|^2 + \left|\partial_z \bar{U}_2(b)\right|^2\right) + \frac{32M^2}{m} \left(\left|\hat{J}_1\right|^2 + \left|\hat{J}_2\right|^2\right),
\]
where \(M := \max(|a|, |b|)\) and \(m := \min(|a|, |b|)\). Hence, we deduce that
\[
\frac{1}{2} \left\|\partial_z \bar{U}\right\|_{L^2}^2 + \frac{m}{2} \left(\left|\partial_z \bar{U}_1(a)\right|^2 + \left|\partial_z \bar{U}_2(b)\right|^2\right) + \kappa_m \left\|\bar{U}\right\|_{L^2}^2 \\
\leq \left(\kappa_M^2 M + \frac{32M^2 \kappa_M^2}{m}\right) \left(\tau \delta_1 \left\|\bar{U}\right\|_{L^2}^2 + \frac{1}{\delta_1 \kappa_M \kappa_m} \left\|\hat{R}\right\|_{L^2}^2 + \frac{1}{\kappa_m^2} \left(\left|\hat{J}_1\right|^2 + \left|\hat{J}_2\right|^2\right)\right) \\
+ 8M^2 \left\|\hat{R}\right\|_{L^2}^2 + \frac{32M^2}{m} \left(\left|\hat{J}_1\right|^2 + \left|\hat{J}_2\right|^2\right),
\]
and this implies
\[
\frac{1}{2} \left\|\partial_z \bar{U}\right\|_{L^2}^2 + \frac{m}{2} \left(\left|\partial_z \bar{U}_1(a)\right|^2 + \left|\partial_z \bar{U}_2(b)\right|^2\right) \\
+ \left(\kappa_m^2 - \left(\kappa_M^2 M + \frac{32M^2 \kappa_M^2}{m}\right) \tau \delta_1\right) \left\|\bar{U}\right\|_{L^2}^2 \\
\leq C \left(\left\|\hat{R}\right\|_{L^2}^2 + \left|\hat{J}_1\right|^2 + \left|\hat{J}_2\right|^2\right).
\]
By choosing \(\delta_1 < (1/(2\tau^3))(M + 32M^2/m)^{-1}\), we derive our required estimate
\[
\varphi = 2z\partial_z \bar{U} = \begin{cases} 
2z\partial_z \bar{U}_1, & z \in (0, a) =: I_1, \\
2z\partial_z \bar{U}_2, & z \in (b, 0) =: I_2.
\end{cases}
\]

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smoothly windowed lattice sums for doubly periodic Green functions in three-dimensional