

ATIYAH COVERING INDEX THEOREM FOR RIEMANNIAN FOLIATIONS

December 20, 2018

MOULAY-TAHAR BENAMEUR AND JAMES L. HEITSCH

ABSTRACT. We use the symbol calculus for foliations developed in [BH17a] to derive a cohomological formula for the Connes-Chern character of the Type II spectral triple given in [BH17b]. The same proof works for the Type I spectral triple of Connes-Moscovici. The cohomology classes of the two Connes-Chern characters induce the same map on the image of the maximal Baum-Connes map in K-theory, thereby proving an Atiyah L^2 covering index theorem.

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1. INTRODUCTION

In [CM95], Connes and Moscovici developed “a general, in some sense universal, local index formula for arbitrary spectral triples of finite summability degree, in terms of the Dixmier trace and its residue-type extension.” Because of its extremely wide applicability, their formulas are quite complicated. See also [CM98]. In [K97], Kordyukov considered a Type I Connes-Moscovici spectral triple naturally associated to a compact foliated manifold with a transversely elliptic operator with holonomy invariant principal symbol. He proved that the spectral triple is finite dimensional, and that its spectrum is simple and contained in the set $\{m \in \mathbb{N} \mid m \leq q\}$, where q is the co-dimension of the foliation. In [BH17b], we extend Kordyukov’s result to certain non-compact manifolds, including non-compact Galois coverings of compact foliated manifolds, which yields a Type II spectral triple (also called a semi-finite spectral triple), see [B03, BF06], with the same properties as the Type I Connes-Moscovici spectral triple.

In this paper we use the symbol calculus for foliations developed in [BH17a] to derive a cohomological formula for the Connes-Chern characters of these Type I and Type II spectral triples. In particular, we assume that we have a Riemannian foliation F of a compact manifold M with a transverse Dirac operator, and that there is a complementary foliation transverse to F . The formulas we obtain are similar to that for the classical case of a Dirac operator on a compact manifold given in [BIF90]. The same proof works for both the Type I spectral triple of Connes-Moscovici and Kordyukov and the Type II spectral triple of [BH17b]. The Connes-Chern characters we obtain induce the same map on the image of the Baum-Connes map in K-theory, [BC00], so we obtain an Atiyah L^2 type covering index theorem for compact foliated manifolds with transversely elliptic operators.

MSC 2010: 19K56, 46L80, 58B34.

Key words: Connes-Chern character, spectral triples, foliations.

It is possible to dispense with the assumption that M admits a transverse foliation. We show how to do this in Section 7 using the standard Morita reduction to transversals, and we consider spectral triples associated to a chosen complete transversal of F . There are several reasons why we prefer to work on M . Working on M doesn't involve choices while working on T does, so working on M is more natural. Moreover, for important generalizations, see [BH17b], we want a formula in terms of characteristic classes on the ambient manifold M . There are natural examples where the restriction to a transversal doesn't work, but our techniques here do. See Example 7.2.

Acknowledgements. It is a pleasure to thank A. Carey, P. Carrillo-Rouse, T. Fack, G. Hector, and P. Piazza for helpful discussions. We are also indebted to the referee for his cogent suggestions. MB wishes to thank the french National Research Agency for support via the project ANR-14-CE25-0012-01 (SINGSTAR).

2. BACKGROUND

We assume that the reader is familiar with the paper [BH17a], and we will freely use the notations of that paper. In particular, F is a smooth Riemannian foliation of the smooth closed Riemannian manifold M . Then (M, F) is of bounded geometry, so all the leaves of F and all the bundles associated to M and F are of bounded geometry. Denote by TM and T^*M the tangent and cotangent bundles of M . The dimension of F is p and the dimension of M is n , so the codimension of F is $q = n - p$, which we assume is even. If q is not even, we replace M by $M \times \mathbb{S}^1$ and F by the obvious p dimensional foliation it determines on $M \times \mathbb{S}^1$. The normal bundle of F is denoted ν ($= TF^\perp$), and its conormal bundle is ν^* . The induced metrics on ν and ν^* are assumed to be bundle like. The leaf of F through a point x is denoted L_x . The homotopy groupoid of F is \mathcal{G} , with the source and target maps $s, r : \mathcal{G} \rightarrow M$. F_s is the foliation of \mathcal{G} whose leaves are given by $\tilde{L}_x = s^{-1}(x)$, which is also denoted \mathcal{G}_x . F_r is the foliation of \mathcal{G} whose leaves are given by $\tilde{L}^y = r^{-1}(y)$, which is also denoted \mathcal{G}^y . We denote $\mathcal{G}_x \cap \mathcal{G}^y$ by \mathcal{G}_x^y . Note that $r : \tilde{L}_x \rightarrow L_x$ is the simply connected covering map. If $\gamma \in \mathcal{G}_x$, the holonomy along γ from $x = s(\gamma)$ to $r(\gamma)$ in L_x is denoted h_γ .

Denote by $E \rightarrow M$ a smooth complex vector bundle over M with a Hermitian structure. We assume that E is basic, which means that there is an action on it by the *holonomy groupoid* and hence also by the homotopy groupoid, and that the Hermitian connection on it, denoted ∇^E , is locally projectable, so its curvature Ω^E is basic, that is locally a pull-back from a transversal.

Assume that F is transversely spin. Fix a spin structure on ν^* , and denote by $\mathcal{S}_\nu = \mathcal{S}_\nu^+ \oplus \mathcal{S}_\nu^-$ the associated spin bundle, with its natural splitting. The bundles \mathcal{S}_ν^\pm are automatically basic vector bundles over M . The Hilbert space of L^2 -sections of the bundle $\mathcal{S}_\nu \otimes E$ over M is denoted \mathcal{H} , which is \mathbb{Z}_2 graded since $\mathcal{S}_\nu \otimes E$ is \mathbb{Z}_2 -graded.

The Levi-Civita connection on ν^* (respectively ν) is denoted ∇^ν , which is locally the pull-back of the Levi-Civita connection on a transversal. It is also known as a Bott or basic connection. The connection ∇^ν induces a connection on \mathcal{S}_ν , which when combined with the connection ∇^E on E gives a connection, also denoted ∇^ν , on $\mathcal{S}_\nu \otimes E$. The context should make clear on which bundle ∇^ν is acting.

The transverse Dirac operator D on \mathcal{H} is given as follows. Choose a local orthonormal basis f_1, \dots, f_q of ν^* with dual orthonormal basis e_1, \dots, e_q of ν . For $u \in \mathcal{H}$, set

$$\tilde{D}(u) = \sum_{1 \leq i \leq q} f_i \cdot \nabla_{e_i}^\nu u,$$

where $f_i \cdot$ is the operator $c(f_i)$, Clifford multiplication by f_i . In general, \tilde{D} is not self-adjoint. The mean curvature vector field of F is $\mu = \sum_{j=1}^p p_\nu(\nabla_{X_j} X_j)$ where X_1, \dots, X_p is a local orthonormal framing of TF , ∇ is the Levi-Civita connection on M , and $p_\nu : TM \rightarrow \nu$ is the projection. When we think of μ as a covector (the isomorphism $\nu \simeq \nu^*$ being given by the inner product), then we denote Clifford multiplication by it by $c(\mu)$, and it is given by, see [BH17a],

$$c(\mu) = \sum_{j=1}^p \sum_{i=1}^q \langle [e_i, X_j], X_j \rangle f_i.$$

The transverse Dirac operator D associated to F is

$$D = \tilde{D} - \frac{1}{2}c(\mu),$$

which is self-adjoint. See [K07, GlK91].

3. STATEMENT OF THE MAIN RESULT

Recall the type I spectral triple introduced in [CM95] and [K97]:

Definition 3.1. *The Connes-Moscovici Type I spectral triple associated to F is $(C_c^\infty(\mathcal{G}), \mathcal{H}, D)$. The usual operator trace yields the super-trace, denoted Tr_s , constructed using the Berezin integral [G83, BGV92], and $a \in C_c^\infty(\mathcal{G})$ acts on \mathcal{H} by*

$$a(u)(x) = \int_{\mathcal{G}_x} a(\gamma^{-1})h_{\gamma^{-1}}(u(r(\gamma)))d\gamma.$$

The main theorem of [K97] is (see also [CM95]),

Theorem 3.2. *$(C_c^\infty(\mathcal{G}), \mathcal{H}, D)$ is a regular even spectral triple with simple dimension spectrum contained in the set $\{k \in \mathbb{N} \mid k \leq q\}$.*

We now recall an important special case of the main result of [BH17b], which will be used in the sequel. In [BH17b], we were concerned with a general bounded geometry foliation $(\widehat{M}, \widehat{\mathcal{F}})$ endowed with a proper action of a discrete group Γ , and with the associated transverse noncommutative geometry as provided by Connes' formalism of regular spectral triples, [C87], and their semi-finite version as introduced in [BF06]. By using a classical reduction method, we were led to the corresponding problem for the class of *Riemannian bifoliations* $(\widehat{M}, \widehat{\mathcal{F}} \subset \widehat{\mathcal{F}}')$, see again [BH17b].

In the case of a free and proper action, we get a bounded geometry Γ -covering $(\widehat{M}, \widehat{\mathcal{F}} \subset \widehat{\mathcal{F}}') \rightarrow (M, \mathcal{F} \subset \mathcal{F}')$ of bounded geometry bifoliations. The main result of [BH17b] is the following.

Theorem 3.3. *Assume that \widehat{D} is a transversely elliptic Γ -invariant pseudodifferential operator in the Connes-Moscovici sense for the larger foliation $\widehat{\mathcal{F}}'$, which is essentially self-adjoint with holonomy invariant transverse principal symbol. Then the triple $(C_c^\infty(\mathcal{G}), (\mathcal{M}, \tau), \widehat{D})$ is a semi-finite spectral triple which is finitely summable of dimension equal to the Beals-Greiner codimension of $\widehat{\mathcal{F}}'$.*

Here \widehat{D} acts on the Hilbert space $\widehat{\mathcal{H}}$ of L^2 -sections of a bundle over \widehat{M} . \mathcal{G} is the monodromy groupoid of \mathcal{F} , $\mathcal{M} = B(\widehat{\mathcal{H}})^\Gamma$ is the Atiyah von Neumann algebra of Γ -invariant operators on $\widehat{\mathcal{H}}$, with its semi-finite trace τ as defined for instance in [At76].

For the special case of a Γ -covering $(\widehat{M}, \widehat{\mathcal{F}}) \rightarrow (M, \mathcal{F})$, (so $\widehat{\mathcal{F}}' = \widehat{\mathcal{F}}$), of a Riemannian foliation, this theorem becomes

Theorem 3.4. *Assume that \widehat{D} is a transversely elliptic Γ -invariant pseudodifferential operator on \widehat{M} , which is essentially self-adjoint with holonomy invariant transverse principal symbol. Then the triple $(C_c^\infty(\mathcal{G}), (\mathcal{M}, \tau), \widehat{D})$ is a type II spectral triple which is finitely summable of dimension equal to the codimension of the foliation, with simple semi-finite dimension spectrum contained in the set $\{k \in \mathbb{N} \mid k \leq q\}$.*

For technical reasons which will be explained later, we now assume that the normal bundle ν is integrable, that is M has a q dimensional foliation which is transverse to F . This foliation is denoted F_\natural . Since F is transversely spin, F_\natural is a spin foliation.

Now suppose again that $\rho: \widehat{M} \rightarrow M$ is a Galois covering, and consider $M \subset \widehat{M}$ as a fundamental domain. Set $\widehat{F} = \rho^{-1}F$, with leaves \widehat{L} , $\widehat{\nu} = \rho^*\nu$, and $\widehat{\nu}^* = \rho^*\nu^*$. We also have the corresponding pulled back objects $\widehat{\mathcal{S}}_\nu \otimes \widehat{E}$, $\widehat{\mathcal{H}} = L^2(\widehat{M}, \widehat{\mathcal{S}}_\nu \otimes \widehat{E})$, \widehat{D} , $\widehat{\mathcal{G}}$, etc.

Note that in general the inverse image of a leaf L of F may consist of a number of leaves of \widehat{F} , and that for each of those leaves $\rho: \widehat{L} \rightarrow L$ is a connected covering of L . Since the simply connected covering $\widetilde{\widehat{L}} \rightarrow \widehat{L}$ of the covering $\widehat{L} \rightarrow L$ is just the simply connected covering $\widetilde{L} \rightarrow L$, the following is obvious.

Lemma 3.5. *Given $\gamma \in \mathcal{G}_{\rho(\hat{x})}$, there exists a unique $\hat{\gamma} \in \widehat{\mathcal{G}}_{\hat{x}}$ so that $\rho \circ \hat{\gamma} = \gamma$.*

This lemma is not true in general if we use holonomy groupoids instead of homotopy groupoids.

Definition 3.6. *Let $a \in C_c^\infty(\mathcal{G})$ act on $\hat{u} \in \widehat{\mathcal{H}}$ by*

$$a(\hat{u})(\hat{x}) = \int_{\mathcal{G}_{\rho(\hat{x})}} a(\gamma^{-1})h_{\hat{\gamma}^{-1}}(\hat{u}(\hat{r}(\hat{\gamma})))d\gamma = \int_{\widehat{\mathcal{G}}_{\hat{x}}} a((\rho \circ \hat{\gamma})^{-1})h_{\hat{\gamma}^{-1}}(\hat{u}(\hat{r}(\hat{\gamma})))d\hat{\gamma}.$$

Definition 3.7. *The Atiyah-Connes spectral triple associated to a (possibly non-compact) Galois foliation cover $\rho : (\widehat{M}, \widehat{F}) \rightarrow (M, F)$, with covering group Γ as above, is $(C_c^\infty(\mathcal{G}), B(\widehat{\mathcal{H}})^\Gamma, \widehat{D})$. The trace is the Atiyah-trace which yields the super trace, denoted τ_s , constructed using the Berezin integral, but restricted to the fundamental domain $M \subset \widehat{M}$.*

Note the following.

1) Since we are using the homotopy groupoid, we want the algebra of smooth functions on the homotopy groupoid $\widehat{\mathcal{G}}$ of \widehat{F} , which are Γ -invariant and Γ -compactly supported. This is precisely the algebra downstairs, that is $C_c^\infty(\mathcal{G})$.

2) The von Neumann algebra $B(\widehat{\mathcal{H}})^\Gamma$ is the algebra used by Atiyah in the L^2 covering index theorem.

3) The Dirac operator \widehat{D} on \widehat{M} is defined with respect to Γ -invariant data and hence is Γ -invariant.

4) For a compact covering, this is a slight modification of the Type I Connes-Moscovici spectral triple above.

Denote by d_ν the transverse de Rham differential for the foliated manifold (\mathcal{G}, F_s) , [T97], and by dx_F the volume form along the leaves of F . Our main result is the following.

Theorem 3.8. *Let $(\widehat{M}, \widehat{F})$ be a covering foliation of (M, F) . Assume that F has even codimension, is transversely spin, and that its normal bundle ν is integrable. Let $a_0, \dots, a_k \in C_c^\infty(\mathcal{G})$. The sequences $(\phi_k)_{k \geq 0}$, and $(\hat{\phi}_k)_{k \geq 0}$, k even, where*

$$\phi_k(a_0, \dots, a_k) = \frac{1}{k!} \int_M \sum_{\gamma \in \mathcal{G}_x} (a_0 d_\nu a_1 \cdots d_\nu a_k)(\gamma^{-1}) \wedge \text{ch}(E) \wedge \widehat{A}(\nu^*) \wedge dx_F,$$

and

$$\hat{\phi}_k(a_0, \dots, a_k) = \frac{1}{k!} \int_{M \subset \widehat{M}} \sum_{\hat{\gamma} \in \widehat{\mathcal{G}}_x} (a_0 d_\nu a_1 \cdots d_\nu a_k)((\rho \circ \hat{\gamma})^{-1}) \wedge \text{ch}(\widehat{E}) \wedge \widehat{A}(\widehat{\nu}^*) \wedge dx_F$$

are cocycles in the (b, B) bicomplex for $C_c^\infty(\mathcal{G})$.

The cohomology classes of (ϕ_k) and $(\hat{\phi}_k)$ are respectively the Connes-Chern characters of $(C_c^\infty(\mathcal{G}), \mathcal{H}, D)$ and $(C_c^\infty(\mathcal{G}), B(\widehat{\mathcal{H}})^\Gamma, \widehat{D})$. These classes induce the same map on the image of the maximal Baum-Connes assembly map in $K_*(C_{\max}^*(\mathcal{G}))$, [BC00]. That is, we have an Atiyah L^2 covering index theorem for foliation coverings.

Remark 3.9. *The maximal assembly map referred to in this paper is the map induced by leafwise index maps, see for instance [C94]. If \mathcal{G} is torsion-free, this map is conjectured to be an isomorphism, see [BC00]. When \mathcal{G} has torsion, it has a modified definition, and it is not an isomorphism in general. See again [BC00] and also [Tu99].*

Note that the sums $\sum_{\mathcal{G}_x}$ and $\sum_{\widehat{\mathcal{G}}_x}$ are all finite since $a_0 d_\nu a_1 \cdots d_\nu a_k$ has compact support. Note also that the product in $a_0 d_\nu a_1 \cdots d_\nu a_k$ is the convolution wedge product on the groupoid, while the wedges are the simple pointwise wedge product. Finally note that the periodic cyclic cohomology classes of (ϕ_k) and $(\hat{\phi}_k)$ do not depend on either the metric on M (so also on \widehat{M}), nor the choice of normal bundle ν .

4. SYMBOLS AND THE TRANSVERSE SYMBOL SPACE

To begin, we construct a generalized exponential function $\widehat{\text{exp}}_x : TM_x \rightarrow M$. Suppose $X = (X_\nu, X_F) \in \nu_x \oplus TF_x$. The leaves of F and F_\cap through x are denoted L_x and L_x^\cap , and their leafwise exponentials are denoted exp^F and exp^\cap . Denote parallel translation from x to $\text{exp}^F(X_F)$ along $\text{exp}^F(tX_F)$, $t \in [0, 1]$, in L_x of the leafwise flat bundle ν by $h_{\text{exp}^F(X_F)}$. Set

$$\widehat{\text{exp}}(X) = \text{exp}^\cap(h_{\text{exp}^F(X_F)}(X_\nu)).$$

What we are doing is exponentiating first in the leaf direction using the usual exponential function for L_x , and then exponentiating in the transverse direction using the usual exponential for $L_{\text{exp}^F(X_F)}^\cap$.

Remark 4.1. *The map $\widehat{\text{exp}}$ has the essential properties of the usual exponential function: it is a local diffeomorphism from $0 \in TM_x$ to a neighborhood of x , whose differential at 0 is the identity. For $X \in \nu_x$, $\widehat{\text{exp}}(X)$ is a geodesic in the leaf of F_\cap through x .*

There are two natural foliations on TM_x , namely the planes parallel to TF_x and ν_x respectively. Since F is Riemannian, $\widehat{\text{exp}}$ maps these foliations onto the foliations F and F_\cap , respectively.

The results of [BH17a] still hold for the transverse symbol space $SC_\cap^\ell(M, E)$ defined below if $\widehat{\text{exp}}$ is used in place of the usual exponential function.

Choose a smooth bump function α on $M \times M$ which is supported in a neighborhood of the diagonal, and equals one on a neighborhood of the diagonal. We require that the support of α is close enough to the diagonal that $(\pi_M, \widehat{\text{exp}})^{-1} : \text{Supp}(\alpha) \rightarrow TM$ is a diffeomorphism onto the component of $(\pi_M, \widehat{\text{exp}})^{-1}(\text{Supp}(\alpha))$ which contains the zero section, where $\pi_M : TM \rightarrow M$ is the projection. Suppose that $x' = \widehat{\text{exp}}(X)$ where $X \in TM_x$. Given $u_x \in (\mathcal{S}_\nu \otimes E)_x$, denote by $\mathcal{T}_{x,x'}(u_x)$ the parallel translation of u_x along the radial line $t \rightarrow \widehat{\text{exp}}(tX)$ from x to x' .

Definition 4.2. *Given an operator P on sections of $\mathcal{S}_\nu \otimes E$, its symbol $\zeta(P)$ is defined as follows. Let $x \in M$, $\xi \in T^*M_x$, and $u_x \in (\mathcal{S}_\nu \otimes E)_x$. Set*

$$\zeta(P)(x, \xi)(u_x) = P\left(x' \mapsto e^{i(\widehat{\text{exp}}_x^{-1}(x'), \xi)} \alpha(x, x') \mathcal{T}_{x,x'}(u_x)\right) \Big|_{x'=x}.$$

Similarly for \widehat{P} acting on sections of $\widehat{\mathcal{S}}_\nu \otimes \widehat{E}$.

If we write $T^*M = \nu^* \oplus T^*F$, then $\xi \in T^*M$ may be written as $\xi = (\eta, \zeta)$, and we may also write $\zeta(P)(x, \xi) = \zeta(P)(x, \eta, \zeta)$. Denote by $\pi : \nu^* \rightarrow M$ the projection.

Definition 4.3. *The symbol space $S_\cap^\ell(M, E)$ consists of all $p(x, \eta) \in C^\infty(\nu^*, \pi^*(\text{End}(\mathcal{S}_\nu \otimes E)))$ so that, for any multi-indices α and β , there is a constant $C_{\alpha, \beta} > 0$, such that*

$$\|\partial_\eta^\alpha \partial_x^\beta p(x, \eta)\| \leq C_{\alpha, \beta} (1 + |\eta|)^{\ell - |\alpha|}.$$

The topology on $S_\cap^\ell(M, E)$ is given by the semi-norms

$$\rho_{\alpha, \beta} = \inf \{C_{\alpha, \beta} \mid \|\partial_\eta^\alpha \partial_x^\beta p(x, \eta)\| \leq C_{\alpha, \beta} (1 + |\eta|)^{\ell - |\alpha|}\}.$$

If we replace the variable η with the variable σ , this is a subspace of the space $S^{0, \ell}(M, E)$ of [BH17a], for the transverse foliation F_\cap . We shall consider it so, while retaining the use of the variable η .

Of course, $\partial_\eta^\alpha \partial_x^\beta p(x, \eta)$ only makes sense if we specify local coordinates. We will use the so-called ‘‘normal coordinates’’ on M . Normal coordinates at a point $x \in M$ are given by choosing a neighborhood U_x of $0 \in TM_x$ on which $\widehat{\text{exp}} : TM_x \rightarrow M$ is a diffeomorphism, and orthonormal bases of ν_x and TF_x , which define coordinates (x_1, \dots, x_n) on $\nu_x \oplus TF_x = TM_x$. This then defines coordinates (also denoted (x_1, \dots, x_n)) in the neighborhood $\widehat{\text{exp}}(U_x)$ of x . In addition, it also induces coordinates on ν_x^* , T^*F_x , and T^*M_x .

Denote by $\wedge^* \nu^*$ the complexified Grassmann algebra bundle. Then

$$C^\infty(\nu^*, \pi^*(\text{End}(\mathcal{S}_\nu \otimes E))) \cong C^\infty(\nu^*, \pi^*(\wedge^* \nu^* \otimes \text{End}(E))),$$

as $\wedge^* \nu^* \cong \text{Cliff}(\nu^*)$, where $\text{Cliff}(\nu^*)$ is the Clifford algebra. Since q is even, $\text{End}(\mathcal{S}_\nu) \cong \text{Cliff}(\nu^*)$. The reader should note carefully that when we represent endomorphisms of \mathcal{S}_ν as elements of $C^\infty(\wedge^* \nu^*)$ and we

compose them, the operation we use is Clifford multiplication and not wedge product. Note also that we use the convention $\omega \cdot \omega = -\langle \omega, \omega \rangle = -\|\omega\|^2$ for co-vectors.

Following [G83], we treat Clifford multiplication by a *normal* k -co-vector as a differential operator of order k . Thus if $p \in C^\infty(\nu^*, \pi^*(\wedge^k \nu^* \otimes \text{End}(E))) \cap S_{\mathfrak{h}}^{\ell-k}(M, E)$, we say p has grading ℓ .

Definition 4.4. *The transverse symbol space $SC_{\mathfrak{h}}^\ell(M, E)$ is*

$$SC_{\mathfrak{h}}^\ell(M, E) = \sum_{k=0}^q C^\infty(\nu^*, \pi^*(\wedge^k \nu^* \otimes \text{End}(E))) \cap S_{\mathfrak{h}}^{\ell-k}(M, E).$$

Set $SC_{\mathfrak{h}}^\infty(M, E) = \bigcup_\ell SC_{\mathfrak{h}}^\ell(M, E)$ and $SC_{\mathfrak{h}}^{-\infty}(M, E) = \bigcap_\ell SC_{\mathfrak{h}}^\ell(M, E)$.

In the notation of [BH17a], $SC_{\mathfrak{h}}^\ell(M, E)$ is a subspace of the space $SC^{0,\ell}(M, E)$ for the foliation $F_{\mathfrak{h}}$. Thus the results of [BH17a] apply to $SC_{\mathfrak{h}}^\ell(M, E)$.

Remark 4.5. *Recall the quantization map θ^α of [BH17a], which involves an exponential function. If we use $\widehat{\text{exp}}$ in defining θ^α , and $p \in SC_{\mathfrak{h}}^\infty(M, E)$ and u is a section of $\mathcal{S}_\nu \otimes E$, then $(\theta^\alpha(p)(u))(x)$ depends only on u restricted to $L_x^{\mathfrak{h}}$. This is not true in general if we use the usual exponential function. In particular,*

$$\theta^\alpha(p)(u)(x) = (2\pi)^{-q} \int_{\nu_x \times \nu_x^*} e^{-i\langle X, \eta \rangle} p(x, \eta) \alpha(x, \widehat{\text{exp}}_x(X)) \mathcal{T}_{x, \widehat{\text{exp}}_x(X)}^{-1}(u(\widehat{\text{exp}}_x(X))) dX d\eta.$$

Definition 4.6. *A family $p(t) \in SC_{\mathfrak{h}}^\ell(M, E)$ is an asymptotic symbol if there are symbols $p_k \in SC_{\mathfrak{h}}^{\ell-k}(M, E)$, so that the following asymptotic expansion holds as $t \rightarrow 0$,*

$$p(t) \sim \sum_{k=0}^{\infty} t^k p_k.$$

The leading symbol of $p(t)$ is p_0 .

Note that $p(t) \sim \sum_{k=0}^{\infty} t^k p_k$ means that given any $N > 0$, $\lim_{t \rightarrow 0} t^{-N} \left(p(t) - \sum_{k=0}^N t^k p_k \right) = 0$ in $SC_{\mathfrak{h}}^{\ell-N-1}(M, E)$. It does not imply that $\sum_{k=0}^{\infty} t^k p_k$ converges.

If the family $p(t) \in C^\infty(\nu^*, \pi^*(\wedge^k \nu^* \otimes \text{End}(E))) \cap S_{\mathfrak{h}}^\infty(M, E)$, then for $t > 0$, we set

$$p(\eta, t)_t = t^k p(t\eta, t),$$

and extend to all of $SC_{\mathfrak{h}}^\infty(M, E)$ by linearity.

Definition 4.7. *A transverse asymptotic pseudodifferential operator ($\mathfrak{h}A\Psi DO$) is a family of operators P_t on sections of $\mathcal{S}_\nu \otimes E$ so that there is an asymptotic symbol $p(t) \in SC_{\mathfrak{h}}^\ell(M, E)$, with $P_t \sim \theta^\alpha(p(t)_t)$.*

If $p(t) \sim \sum_{k=0}^{\infty} t^k p_k$, the leading symbol of P_t is the symbol p_0 .

Recall that $P_t \sim Q_t$ if for all $N \geq 0$, and all s, k , $\lim_{t \rightarrow 0} t^{-N} \|P_t - Q_t\|_{s,k} = 0$, where, $\|\cdot\|_{s,k}$ is the norm of an operator from the usual s Sobolev space associated to $\mathcal{S}_\nu \otimes E$ to the usual k Sobolev space.

Of course, as all these constructions are local, we may also make them for \widehat{M} , \widehat{F} and \widehat{E} .

5. PRELIMINARY RESULTS

Our calculations here are local so work equally well for both (M, F) and $(\widehat{M}, \widehat{F})$. For simplicity of notation, we will work on (M, F) .

For $a \in C_c^\infty(\mathcal{G}; \wedge^k r^*(\nu^*))$, set $a_t = t^k a$. Below we will encounter expressions of the form $a\mathcal{D}$, where \mathcal{D} is a transverse differential operator of order ℓ exactly, that is it will differentiate exactly ℓ times, no more and no fewer. For such operators, we set $(a\mathcal{D})_t = t^{k+\ell} a\mathcal{D}$. The reader should note that with this definition we have $\varsigma(\mathcal{D})_t = t^\ell \varsigma(\mathcal{D})$, since $\varsigma(\mathcal{D})$ will be a polynomial in η homogenous of degree ℓ .

Suppose $a_0, \dots, a_k \in C_c^\infty(\mathcal{G})$. To prove Theorem 3.8, we need certain facts about

$$t^{k+2|\ell|} (a_0(\delta a_1)^{(\ell_1)} \dots (\delta a_k)^{(\ell_k)})_{1/t} \quad \text{and} \quad \varsigma(e^{-t^2 D^2})_{1/t}.$$

Here, $\ell = (\ell_1, \dots, \ell_k) \in \mathbb{N}^k$, $\delta a_j = [D, a_j]$, $\mathfrak{D}(a_j) = [D^2, a_j]$, and $(\delta a_j)^{(\ell_j)} = \mathfrak{D}^{\ell_j}(\delta a_j)$. If we set

$$a^{k,\ell}(D) = a_0(\delta a_1)^{(\ell_1)} \dots (\delta a_k)^{(\ell_k)},$$

then,

$$t^{k+2|\ell|} a_0(\delta a_1)^{(\ell_1)} \dots (\delta a_k)^{(\ell_k)} = a^{k,\ell}(tD).$$

5.1. Results on $\delta a = [D, a]$.

First we recall some notation. For $x \in M$, $\mathcal{G}_x = s^{-1}(x)$, that is homotopy equivalence classes of paths, with end points fixed, in the leaf L_x starting at x . The holonomy along an element $\gamma \in \mathcal{G}$ is denoted h_γ . So h_γ transports objects at $s(\gamma)$ to $r(\gamma)$ along the path γ in the leaf $L_{s(\gamma)}$. An element $a \in C_c^\infty(\mathcal{G}; \wedge^k r^*(\nu^*))$ acts on sections of $\mathcal{S}_\nu \otimes E$ by

$$a(u)(x) = \int_{\mathcal{G}_x} a(\gamma^{-1}) h_{\gamma^{-1}} u(r(\gamma)) d\gamma.$$

The operator \tilde{D} acts on a as follows. Choose double foliation coordinate charts $U, (y, z)$ and $V, (y', z')$ respectively of $s(\gamma)$ and $r(\gamma)$, where y and y' are transverse coordinates and z and z' are tangential coordinates. The neighborhood (U, γ, V) of γ consists of all leafwise paths which start in U , end in V , and are parallel with γ . See [HL90] for more details. Coordinates on (U, γ, V) are given by (y, z, z') . In essence, we identify the transverse coordinates using the holonomy h_γ of γ . The operator \tilde{D} is preserved by h_γ , since h_γ preserves the metrics on ν and ν^* and the connection ∇^ν . In addition, \tilde{D} is completely determined by its restriction to the leaves of the transverse foliation. So \tilde{D} does not depend on the tangential coordinates at all, and we may assume that γ , so also $\gamma^{-1} \in (V, \gamma^{-1}, U)$, depends only on the coordinate y , that is we may view a as a section of $\wedge^k \nu^*$ defined on the transverse leaf through $x \in U$. Then $\tilde{D}(a)(\gamma^{-1})$ is well defined, and by definition equals $d_\nu a(\gamma^{-1})$.

Recall that $c(\mu)$ is Clifford multiplication by the mean curvature vector field of F (thought of as a co-vector), and set $\tilde{c}(\gamma) = h_\gamma(c(\mu)(s(\gamma))) - c(\mu)(r(\gamma))$. Then

$$5.2. \quad [tD, a]_{1/t} = d_\nu a + \frac{1}{2} \tilde{c}(\mu) a.$$

We have

$$[tD, a](u)(x) = tD_x \int_{\mathcal{G}_x} a(\gamma^{-1}) h_{\gamma^{-1}} u(r(\gamma)) d\gamma - t \int_{\mathcal{G}_x} a(\gamma^{-1}) h_{\gamma^{-1}} ((Du)(r(\gamma))) d\gamma.$$

Now $D = \tilde{D} - \frac{1}{2}c(\mu)$, and since $h_{\gamma^{-1}}$ preserves the metrics on ν and ν^* and the connection ∇^ν , we have $[\tilde{D}, h_{\gamma^{-1}}] = 0$. Thus

$$[tD, a](u)(x) = t \int_{\mathcal{G}_x} d_\nu a(\gamma^{-1}) h_{\gamma^{-1}} u(r(\gamma)) d\gamma + \frac{t}{2} \int_{\mathcal{G}_x} a(\gamma^{-1}) (h_{\gamma^{-1}} c(\mu)(r(\gamma)) - c(\mu)(s(\gamma))) h_{\gamma^{-1}} u(r(\gamma)) d\gamma.$$

Note that

$$h_{\gamma^{-1}} c(\mu)(r(\gamma)) - c(\mu)(s(\gamma)) = h_{\gamma^{-1}} c(\mu)(s(\gamma^{-1})) - c(\mu)(r(\gamma^{-1})) = \tilde{c}(\mu)(\gamma^{-1}),$$

so

$$[tD, a] = td_\nu a + \frac{t}{2} \tilde{c}(\mu) a,$$

and

$$[tD, a]_{1/t} = d_\nu a + \frac{1}{2} \tilde{c}(\mu) a,$$

where the normal one form parts act by Clifford multiplication.

5.3. Results on $t^2 D^2$.

Choose dual orthonormal bases, f_1, \dots, f_q of ν^* and e_1, \dots, e_q of ν , which are parallel (using ∇^ν) along radial lines through x . From the proof of Proposition 3.10 of [BH17a], we have

$$D^2 = \tilde{D}^2 - \frac{1}{2}\tilde{D}(c(\mu)) + \nabla_\mu + \frac{1}{4}|\mu|^2,$$

and

$$\tilde{D}^2 = -\sum_j \nabla_{e_j}^\nu \nabla_{e_j}^\nu + \sum_{j < k} f_j \wedge f_k \nabla_{[e_j, e_k]}^\nu + \sum_{j < k} f_j \wedge f_k \Omega_E(e_j, e_k) + \frac{1}{8}\kappa.$$

The term κ is the scalar curvature of the leaves of the transverse foliation F_{th} , and Ω_E is the curvature of the bundle E over M .

The term $\sum_{j < k} f_j \wedge f_k \nabla_{[e_j, e_k]}^\nu$ is problematic, since its symbol may have grading three, and in order to apply the results of [BH17a], we need the terms of the symbol of D^2 to have grading at most two. This is why we assume that ν is integrable.

Lemma 5.4. *If ν is integrable, $[e_j, e_k](x) = 0$, so also $\sum_{j < k} f_j \wedge f_k \nabla_{[e_j, e_k]}^\nu(x) = 0$.*

Proof. Denote by P_x^ν the plaque of the foliation determined by ν through the point x . The e_j are tangent to P_x^ν , since $TP_x^\nu = \nu|_{P_x^\nu}$, and parallel translation by ∇^ν preserves ν . Recall that ∇^ν is the pull back of the Levi-Civita connection ∇^P on P_x^ν under the local projection to P_x^ν along the leaves of F . We have

$$0 = \left(\nabla_{e_i}^\nu e_j - \nabla_{e_j}^\nu e_i \right)(x) = \left(\nabla_{e_i}^P e_j - \nabla_{e_j}^P e_i \right)(x) = [e_i, e_j](x),$$

since ∇^P is torsion free. □

Note carefully that all the remaining terms in \tilde{D}^2 are completely determined by their restriction to F_{th} , and are projectable, that is invariant under the holonomy.

Recall, [BH17a], that

$$\Xi = -\sum_{i: j \neq k} f_k \wedge f_j \otimes e_k(\langle [e_j, X_i], X_i \rangle),$$

where X_1, \dots, X_p is a local orthonormal framing of TF , is the grading two part of the Clifford form $-\tilde{D}(c(\mu))$. Now,

$$t^2 D^2 = t^2 \left(\sum_{j < k} f_j \wedge f_k \Omega_E(e_j, e_k) - \frac{1}{2}\tilde{D}(c(\mu)) + \frac{1}{8}\kappa - \sum_j \nabla_{e_j}^\nu \nabla_{e_j}^\nu + \nabla_\mu + \frac{1}{4}|\mu|^2 \right),$$

which we may write as

$$\mathbf{5.5.} \quad t^2 D^2 = t^2 \left(\sum_{j < k} f_j \wedge f_k \Omega_E(e_j, e_k) + \frac{1}{2}\Xi \right) + t^2 \left(\frac{1}{8}\kappa - \sum_j \nabla_{e_j}^\nu \nabla_{e_j}^\nu + \nabla_\mu + \frac{1}{4}|\mu|^2 - \frac{1}{2}(\tilde{D}(c(\mu)) + \Xi) \right).$$

The two forms occurring in the first term are acting by Clifford multiplication. All the operators in the second term have grading at most one, in the sense that when applied to an element of $C_c^\infty(\mathcal{G}; r^*(\wedge \nu^*))$, e. g. $[tD, a]$, they raise its grading by at most one, e. g.

$$[\nabla_{e_j}^\nu \nabla_{e_j}^\nu, a] = \nabla_{e_j}^\nu (\nabla_{e_j}^\nu(a)) + 2\nabla_{e_j}^\nu(a) \nabla_{e_j}^\nu.$$

The first term on the right has grading the same as a , while the second has grading one more. In particular,

$$[t^2 \nabla_{e_j}^\nu \nabla_{e_j}^\nu, [tD, a]]_{1/t}$$

is polynomial in t and has no constant term, so $\lim_{t \rightarrow 0} [t^2 \nabla_{e_j}^\nu \nabla_{e_j}^\nu, [tD, a]]_{1/t} = 0$, and the same is true for all the other operators in the second term of Equation 5.5. Thus, when we rescale by $1/t$, the t^2 in the first term will disappear, but not in the second term, so that term (when applied to an element of $C_c^\infty(\mathcal{G}; r^*(\wedge \nu^*))$) will not play a role in the computation of the trace. (However, the term $\sum_j \nabla_{e_j}^\nu \nabla_{e_j}^\nu$, or rather its symbol, will play a crucial role in that computation.) For simplicity we write $t^2 D^2$ as

$$t^2 D^2 = \frac{t^2}{2}(\Omega_E + \Xi) + t^2 \Lambda.$$

Set

$$\tilde{\Xi}(\gamma) = \Xi(r(\gamma)) - h_\gamma \Xi(s(\gamma)).$$

$$5.6. \quad [t^2 D^2, [tD, a]]_{1/t} = \left(\frac{1}{2} \tilde{\Xi}\right)(d_\nu a + \frac{1}{2} \tilde{c}(\mu) a) + f_1(t),$$

where f_1 is polynomial in all variables, and $f_1(0) = 0$.

By the comments above about $\Lambda = \frac{1}{8} \kappa - \sum_j \nabla_{e_j}^\nu \nabla_{e_j}^\nu + \nabla_\mu + \frac{1}{4} |\mu|^2 - \frac{1}{2} (\tilde{D}(c(\mu)) + \Xi)$, we have

$$[t^2 D^2, [tD, a]]_{1/t} = \frac{1}{2} [\Omega_E + \Xi, d_\nu a + \frac{1}{2} \tilde{c}(\mu) a] + [t^2 \Lambda, [D, a]]_{1/t},$$

where the last term is polynomial in t with no constant term. Now, Ω_E is a basic two form, that is a two form which is constant along leaves of F , so $[\Omega_E, d_\nu a + \frac{1}{2} \tilde{c}(\mu) a] = 0$. The argument used in the calculation of $[tD, a]$ shows that

$$[t^2 D^2, [tD, a]]_{1/t} = \left(\frac{1}{2} \tilde{\Xi}\right)(d_\nu a + \frac{1}{2} \tilde{c}(\mu) a) + f_1(t),$$

where f_1 is polynomial in all variables, and $f_1(0) = 0$.

A straight forward induction argument gives

$$5.7. \quad [t^2 D^2, [\dots [t^2 D^2, [tD, a]] \dots]]_{1/t} = \left(\frac{1}{2} \tilde{\Xi}\right)^\ell (d_\nu a + \frac{1}{2} \tilde{c}(\mu) a) + f_\ell(t),$$

where there are ℓ copies of $t^2 D^2$, and f_ℓ is polynomial in all variables, and $f_\ell(0) = 0$. Thus we have

Lemma 5.8.

$$a^{k,\ell}(tD)_{1/t} = a_0 \left[\prod_{j=1}^k \left(\frac{1}{2} \tilde{\Xi}\right)^{\ell_j} (d_\nu a_j + \frac{1}{2} \tilde{c}(\mu) a_j) \right] + f_{k,\ell}(t),$$

where $f_{k,\ell}$ is polynomial in all variables, and $f_{k,\ell}(0) = 0$.

Note that here the $d_\nu a_j + \frac{1}{2} \tilde{c}(\mu) a_j$ are now acting as differential forms, *not* as Clifford forms.

It is very important to note the following. Each term of $f_{k,\ell}$, for $\ell \neq (0, \dots, 0)$, is the composition of elements of $C_c^\infty(\mathcal{G}, \wedge r^*(\nu^*))$ and $\natural\text{A}\Psi\text{DOs}$ (actually differential operators) whose leading symbols are 0. For example, $[t^2 D^2, [t^2 D^2, [tD, a]]]_{1/t}$ contains elements of $C_c^\infty(\mathcal{G}, \wedge r^*(\nu^*))$ multiplied by the elements

$$t \nabla_{e_j}^\nu (d_\nu a) \nabla_{e_k}^\nu \nabla_{e_k}^\nu \nabla_{e_j}^\nu \quad \text{and} \quad t \nabla_{e_j}^\nu (d_\nu a) \nabla_{e_j}^\nu \nabla_{e_k}^\nu \nabla_{e_k}^\nu.$$

Both $t \nabla_{e_k}^\nu \nabla_{e_k}^\nu \nabla_{e_j}^\nu$ and $t \nabla_{e_j}^\nu \nabla_{e_k}^\nu \nabla_{e_k}^\nu$ are transverse asymptotic differential operators whose leading symbols are 0.

5.9. A digression.

For later use, we recall the symbol of $t^2 D^2$ given in Proposition 3.10 of [BH17a], and we assume that we are using the Dominguez metric on M . As we are assuming that the normal bundle is integrable, the term $i\langle \vartheta_\nu, \xi \rangle = 0$ in that calculation. Thus

$$\begin{aligned} \varsigma(D^2)(x, \eta) &= |\eta|^2 - i \sum_j e_{j,x} e_j \langle \widehat{\exp}_x^{-1}(x'), \xi \rangle + \frac{1}{8} \kappa + \sum_{j < k} f_j \wedge f_k \otimes \Omega_E(e_j, e_k) + \\ &\quad \frac{1}{2} \sum_{i,j,k} f_j \wedge f_k \otimes e_k \langle ([e_j, X_i], X_i) \rangle + i\langle \mu, \eta \rangle - \frac{1}{4} |\mu|^2. \end{aligned}$$

Consider the term $e_{j,x} e_j \langle \widehat{\exp}_x^{-1}(x'), \xi \rangle$. Note that the three other terms in the first row are projectable. The three terms in the second row all come from μ , which is projectable if we use the Dominguez metric. As noted in the proof of Proposition 6.3 below, the symbol of D^2 is projectable if we use the Dominguez metric. It follows that $e_{j,x} e_j \langle \widehat{\exp}_x^{-1}(x'), \xi \rangle$ must be projectable, and so equals $e_{j,x} e_j \langle \widehat{\exp}_x^{-1}(x'), \eta \rangle$, which has grading one. Set

5.10.

$$p(x, \eta, t) = |\eta|^2 + \sum_{j < k} f_j \wedge f_k \otimes \Omega_E(e_j, e_k) + \frac{1}{2} \sum_{i; j \neq k} f_j \wedge f_k \otimes e_k(\langle [e_j, X_i], X_i \rangle) +$$

$$t \left(i \langle \mu, \eta \rangle - i \sum_j e_{j,x} e_j \langle \widehat{\exp}_x^{-1}(x'), \eta \rangle \right) + t^2 \left(\frac{1}{8} \kappa - \frac{1}{4} |\mu|^2 - \frac{1}{2} \sum_{i,k} e_k(\langle [e_k, X_i], X_i \rangle) \right).$$

Then $p(x, \eta, t)_t = \varsigma(t^2 D^2)$, so $t^2 D^2$ is an \hbar A Ψ DO with leading symbol

$$p_0^{t^2 D^2}(x, \eta) = |\eta|^2 + \sum_{j < k} f_j \wedge f_k \otimes \Omega_E(e_j, e_k) + \frac{1}{2} \sum_{i; j \neq k} f_i \wedge f_k \otimes e_k(\langle [e_i, X_j], X_j \rangle),$$

which we write as

$$p_0^{t^2 D^2} = |\eta|^2 + \frac{1}{2} (\Omega_E + \Xi),$$

for short.

Note that the leading symbol of $t^2 D^2 - \lambda$, so also $t^2 \widehat{D}^2 - \lambda$, is given by

$$p_0^{t^2 D^2 - \lambda} = |\eta|^2 + \frac{1}{2} (\Omega_E + \Xi) - \lambda.$$

5.11. Results on $\varsigma(e^{-t^2 D^2})_{1/t}$.

Let e_1, \dots, e_q be a local orthonormal basis of ν with dual orthonormal basis f_1, \dots, f_q of ν^* . Denote by Ω_ν the curvature of ∇^ν acting on $\mathcal{S}_\nu \otimes E$, which is a projectable two form. Set

$$\Omega_\nu(e_i, e_j) = \sum_{i,j,k,\ell=1}^q (\Omega_\nu)_{\ell,i,j}^k e_k \otimes f_\ell, \quad \text{that is } (\Omega_\nu)_{\ell,i,j}^k = \langle \Omega_\nu(e_i, e_j)(e_\ell), e_k \rangle,$$

and note that $(\Omega_\nu)_{\ell,i,j}^k$ is skew in the indices i, j (since Ω_ν is a 2-form) as well as the k, ℓ , (since Ω_ν has coefficients in $so_q = spin_q$). Set

$$\Omega_\nu(\partial/\partial\eta, \partial/\partial\eta')(p(x, \eta) \wedge q(x, \eta')) = \sum_{i,j,k,\ell=1}^q (\Omega_\nu)_{\ell,i,j}^k f_k \wedge f_\ell \wedge \frac{\partial p(x, \eta)}{\partial \eta_i} \wedge \frac{\partial q(x, \eta')}{\partial \eta'_j}.$$

Then $e^{-\frac{1}{4}\Omega_\nu(\partial/\partial\eta, \partial/\partial\eta')}$ is a finite sum of compositions of such operators, and the number of compositions is $\leq q/2$ because of the $f_k \wedge f_\ell$. Note that $\Omega_\nu(\partial/\partial\eta, \partial/\partial\eta')$ is identical to the operator used in [G83] and [BIF90]. As in [G83, BIF90, BH17a], the differential operator $a_0(p, q)$, defined on pairs of elements in $SC_{\hbar}^\infty(M, E)$, is given by

$$a_0(p, q)(x, \eta) = e^{-\frac{1}{4}\Omega_\nu(\partial/\partial\eta, \partial/\partial\eta')} p(x, \eta) \wedge q(x, \eta') |_{\eta'=\eta}.$$

It is easy to check that the results of [BIF90] extend to $SC_{\hbar}^\infty(M, E)$ and $SC_{\hbar}^\infty(\widehat{M}, \widehat{E})$, as do those of [BH17a], mutatis mutandis. For $(\widehat{M}, \widehat{E})$, this is because the geometry there is uniformly bounded. In particular, the fact that $p_0^{t^2 D^2} = |\eta|^2 + \frac{1}{2} (\Omega_E + \Xi)$, which is the symbol of a uniformly transversely elliptic differential operator, see [Sh92], implies that $t^2 D^2$ is asymptotically elliptic, in the sense of [BIF90], for the symbol space $SC_{\hbar}^2(M, E)$, and similarly for $t^2 \widehat{D}^2$ and $SC_{\hbar}^2(\widehat{M}, \widehat{E})$. Thus, combining Lemmas 3.13 and 3.14 of [BIF90] in our context here gives

Lemma 5.12. $\varsigma(e^{-t^2 D^2})_{1/t}$ is an asymptotic symbol in $SC_{\hbar}^{-\infty}(M, E)$, with leading symbol $e^{a_0(-p_0^{t^2 D^2}, \cdot)}(1)$, where 1 is the symbol which is the constant function 1.

Next we need to compute, for an arbitrary symbol $q \in SC_{\hbar}^\infty(M, E)$,

$$a_0(-p_0^{t^2 D^2}, q) = e^{-\frac{1}{4}\Omega_\nu(\partial/\partial\eta, \partial/\partial\eta')} ((-p_0^{t^2 D^2})(x, \eta) \wedge q(x, \eta')) |_{\eta'=\eta}.$$

Denote by $\Omega_\nu(\partial/\partial\eta, \partial/\partial\eta)$ the operator which acts as

$$\Omega_\nu(\partial/\partial\eta, \partial/\partial\eta)q(x, \eta) = \sum_{i,j,k,a,b,c,d=1}^q (\Omega_\nu)_{b,j,i}^a (\Omega_\nu)_{d,i,k}^c f_j \wedge f_k \wedge \frac{\partial^2 q(x, \eta)}{\partial\eta_j \partial\eta_k}.$$

Denote by $\Omega_\nu(\eta, \partial/\partial\eta)$ the operator which acts as

$$\Omega_\nu(\eta, \partial/\partial\eta)q(x, \eta) = \sum_{i,j,k,\ell=1}^q (\Omega_\nu)_{\ell,i,j}^k f_k \wedge f_\ell \wedge \eta_i \frac{\partial q(x, \eta)}{\partial\eta_j}.$$

Then, computing as in [G83] and [BIF90], we have

$$a_0(-p_0^{t^2 D^2}, q) = \left[-|\eta|^2 - \frac{1}{2}(\Omega_E + \Xi) + \frac{1}{2}\Omega_\nu(\eta, \partial/\partial\eta) + \frac{1}{16}\Omega_\nu(\partial/\partial\eta, \partial/\partial\eta) \right] q(x, \eta).$$

Thus,

$$e^{a_0(-p_0^{t^2 D^2}, \cdot)}(1) = e^{\left(-|\eta|^2 - \frac{1}{2}(\Omega_E + \Xi) + \frac{1}{2}\Omega_\nu(\eta, \partial/\partial\eta) + \frac{1}{16}\Omega_\nu(\partial/\partial\eta, \partial/\partial\eta)\right)}(1).$$

Note that $\Omega_\nu(\eta, \partial/\partial\eta)(|\eta|^2) = 2 \sum_{i,j,k,\ell=1}^q (\Omega_\nu)_{\ell,i,j}^k f_k \wedge f_\ell \otimes \eta_i \eta_j = 0$, since $(\Omega_\nu)_{\ell,i,j}^k = -(\Omega_\nu)_{\ell,j,i}^k$. (This corrects a typo in [BIF90].) Thus $\Omega_\nu(\eta, \partial/\partial\eta)$ and $|\eta|^2$ commute. Since $\Omega_\nu(\eta, \partial/\partial\eta)$ also commutes with all the other operators in the exponent, this equals

$$e^{\left(-|\eta|^2 - \frac{1}{2}(\Omega_E + \Xi) + \frac{1}{16}\Omega_\nu(\partial/\partial\eta, \partial/\partial\eta)\right)} e^{\frac{1}{2}\Omega_\nu(\eta, \partial/\partial\eta)}(1).$$

As $e^{\frac{1}{2}\Omega_\nu(\eta, \partial/\partial\eta)}(1) = 1$, and Ξ commutes with $|\eta|^2 + \frac{1}{2}\Omega_E - \frac{1}{16}\Omega_\nu(\partial/\partial\eta, \partial/\partial\eta)$, (it is a transverse 2 form, so a nilpotent multiplication operator, which involves no differentiation), we have

$$p_0^{e^{-t^2 D^2}} = e^{a_0(-p_0^{t^2 D^2}, \cdot)}(1) = e^{-\frac{1}{2}\Xi} \wedge e^{\left(-|\eta|^2 - \frac{1}{2}\Omega_E + \frac{1}{16}\Omega_\nu(\partial/\partial\eta, \partial/\partial\eta)\right)}(1),$$

which we write as

$$e^{-\frac{1}{2}\Xi} \wedge e^{-|\eta|^2 - \frac{1}{2}\Omega_E + \frac{1}{16}\Omega_\nu}(1).$$

Note that the second exponent contains only operators which act in normal directions, and is constant along leaves of F .

Finally, the fact that $\varsigma(e^{-t^2 D^2})_{1/t}$ is an asymptotic symbol means that we have

$$\varsigma(e^{-t^2 D^2})_{1/t} \sim e^{-\frac{1}{2}\Xi} \wedge e^{-|\eta|^2 - \frac{1}{2}\Omega_E + \frac{1}{16}\Omega_\nu}(1) + \Upsilon.$$

Here Υ is a power series in t with coefficients in $SC_{\text{fl}}^{-\infty}(M, E)$, in particular, $\Upsilon = \sum_{k=1}^{\infty} t^k p_k(x, \eta)$, where for all $N \geq 1$,

$$\lim_{t \rightarrow 0} t^{-N} \left(\varsigma(e^{-t^2 D^2})_{1/t} - \left[e^{-\frac{1}{2}\Xi} \wedge e^{-|\eta|^2 - \frac{1}{2}\Omega_E + \frac{1}{16}\Omega_\nu}(1) + \sum_{k=1}^N t^k p_k(x, \eta) \right] \right) = 0$$

in $SC_{\text{fl}}^{-\infty}(M, E)$.

Just as in Corollary 3.3, [BIF90] p. 25, we have

Proposition 5.13.

$$\begin{aligned} \lim_{t \rightarrow 0} \left[(2\pi)^{-q} \int_{\nu_x^*} \text{Tr}_s(\varsigma(e^{-t^2 D^2})_{1/t}) d\eta \right] &= (2\pi)^{-q} \int_{\nu_x^*} \lim_{t \rightarrow 0} \text{Tr}_s(\varsigma(e^{-t^2 D^2})_{1/t}) d\eta = \\ &= (2\pi)^{-q} \int_{\nu_x^*} \text{Tr}_s(e^{-\frac{1}{2}\Xi} e^{-|\eta|^2 - \frac{1}{2}\Omega_E + \frac{1}{16}\Omega_\nu}(1)) d\eta = e^{-\frac{1}{2}\Xi} \wedge \text{ch}(E) \wedge \widehat{A}(\nu^*)(x). \end{aligned}$$

The last equality is classical.

6. PROOF OF THE THEOREM 3.8

To prove Theorem 3.8, we need to compute the Connes-Chern character of $(C_c^\infty(\mathcal{G}), \mathcal{H}, D)$, and then note that the same proof works for $(C_c^\infty(\mathcal{G}), B(\widehat{\mathcal{H}})^\Gamma, \widehat{D})$, mutatis mutandis. Combined with Proposition 6.5 below, this gives the proof of the theorem.

Suppose $a_0, \dots, a_k \in C_c^\infty(\mathcal{G})$. In the notation of [CM95], Theorem II.3, p. 230, for $k > 0$, k even, (there is a separate formula for ϕ_0 , see below) the Connes-Chern character of $(C_c^\infty(\mathcal{G}), \mathcal{H}, D)$ is given by the local formula

$$\phi_k(a_0, \dots, a_k) = \sum_{\ell, q} \frac{(-1)^{|\ell|}}{\ell_1! \dots \ell_k!} \alpha_\ell \sigma_q(|\ell| + \frac{k}{2}) \tau_q(\gamma a^{k, \ell}(D)(I+D^2)^{-k/2-|\ell|}).$$

As above, $a^{k, \ell}(D) = a_0(\delta a_1)^{(\ell_1)} \dots (\delta a_k)^{(\ell_k)}$, where $\delta a_j = [D, a_j]$, $\mathfrak{D}(a_j) = [D^2, a_j]$, $(\delta a_j)^{(\ell_j)} = \mathfrak{D}^{\ell_j}(\delta a_j)$, γ is the grading operator, and τ_q is a certain residue. Only $q = 0$ actually occurs, since the dimension spectrum is simple. Note that $\sigma_0(k/2 + |\ell|) = \Gamma(k/2 + |\ell|)$, and $\alpha_\ell^{-1} = (\ell_1 + 1)(\ell_1 + \ell_2 + 2) \dots (\ell_1 + \dots + \ell_k + k)$. Then, for $k > 0$,

$$\begin{aligned} \tau_0(\gamma a^{k, \ell}(D)(I+D^2)^{-k/2-|\ell|}) &= \text{Res}_{z=0} \left[\text{Tr}_s(a^{k, \ell}(D)(I+D^2)^{-k/2-|\ell|-z}) \right] = \\ &= \text{Res}_{z=0} \left[\text{Tr}_s(\Gamma(k/2 + |\ell| + z)^{-1} \int_0^\infty t^{k/2+|\ell|+z-1} a^{k, \ell}(D) e^{-t(I+D^2)} dt) \right] = \\ &= \Gamma(k/2 + |\ell|)^{-1} \text{Res}_{z=0} \left[\text{Tr}_s \left(\int_0^\infty 2t^{k+2|\ell|+2z-1} a^{k, \ell}(D) e^{-t^2(I+D^2)} dt \right) \right] = \\ &= \Gamma(k/2 + |\ell|)^{-1} \text{Res}_{z=0} \left[\int_0^\infty 2t^{2z-1} e^{-t^2} \text{Tr}_s(a^{k, \ell}(tD) e^{-t^2 D^2}) dt \right]. \end{aligned}$$

Lemma 6.1. *For any $\epsilon > 0$,*

$$\text{Res}_{z=0} \left[\int_0^\infty 2t^{2z-1} e^{-t^2} \text{Tr}_s(a^{k, \ell}(tD) e^{-t^2 D^2}) dt \right] = \text{Res}_{z=0} \left[\int_0^\epsilon 2t^{2z-1} e^{-t^2} \text{Tr}_s(a^{k, \ell}(tD) e^{-t^2 D^2}) dt \right].$$

Proof. We need to show that

$$\int_\epsilon^\infty 2t^{2z-1} e^{-t^2} \text{Tr}_s(a^{k, \ell}(tD) e^{-t^2 D^2}) dt = \int_\epsilon^\infty 2t^{2z+k+2|\ell|-1} e^{-t^2} \text{Tr}_s(a^{k, \ell}(D) e^{-t^2 D^2}) dt$$

is finite.

If T is an operator of order j , the operator $[D^2, T] = |D| [|D|, T] + [|D|, T] |D|$ has order $j + 1$. So the operator $a^{k, \ell}(D)$ has order $|\ell|$, and the operator $a^{k, \ell}(D)(1 + D^2)^{-|\ell|/2}$ is bounded. Now

$$|\text{Tr}_s(a^{k, \ell}(D) e^{-t^2 D^2})| = |\text{Tr}_s(a^{k, \ell}(D)(1 + D^2)^{-|\ell|/2} (1 + D^2)^{|\ell|/2} e^{-t^2 D^2})|$$

and $(1 + D^2)^{|\ell|/2} e^{-t^2 D^2}$ is positive and trace class. As $|\text{Tr}_s(AB)| \leq \|A\| \text{Tr}(B)$ for A bounded and B positive and trace class, we have

$$|\text{Tr}_s(a^{k, \ell}(D) e^{-t^2 D^2})| \leq \|a^{k, \ell}(D)(1 + D^2)^{-|\ell|/2}\| \text{Tr}((1 + D^2)^{|\ell|/2} e^{-t^2 D^2}).$$

For $t \geq \epsilon$,

$$\text{Tr}((1 + D^2)^{|\ell|/2} e^{-t^2 D^2}) \leq \text{Tr}((1 + D^2)^{|\ell|/2} e^{-\epsilon^2 D^2}).$$

Therefore, for any $z \in \mathbb{C}$,

$$\begin{aligned} \left| \int_\epsilon^\infty 2t^{2z+k+2|\ell|-1} e^{-t^2} \text{Tr}_s(a^{k, \ell}(D) e^{-t^2 D^2}) dt \right| &\leq \\ \|a^{k, \ell}(D)(1 + D^2)^{-|\ell|/2}\| \text{Tr}((1 + D^2)^{|\ell|/2} e^{-\epsilon^2 D^2}) \int_\epsilon^\infty 2t^{2\text{re}(z)+k+2|\ell|-1} e^{-t^2} dt &< \infty, \end{aligned}$$

where $\text{re}(z)$ is the real part of z . □

Suppose that $(\widehat{M}, \widehat{F})$ is a covering foliation of (M, F) . For $a \in C_c^\infty(\mathcal{G}; \wedge^k r^*(\nu^*))$, set $a_t(\gamma) = t^k a(\gamma)$. We identify $x \in M$ with $\bar{x} \in \mathcal{G}$, the class of the constant path at x . So $a(x) := a(\bar{x})$. Recall that the dual normal bundle of \widehat{F} is $\widehat{\nu}^*$, and that we are thinking of M as being a fundamental domain $M \subset \widehat{M}$. So for a symbol $p \in SC_{\text{rh}}^\ell(\widehat{M}, \widehat{E})$, $p(x, \eta)$ makes sense for $x \in M \subset \widehat{M}$ and $\eta \in \widehat{\nu}_x^* = \nu_x^*$. Denote by dx_F the volume form along the leaves of F .

The trace and supertrace for M are the usual ones, denoted Tr and Tr_s , while those for \widehat{M} , which are the usual traces restricted to $M \subset \widehat{M}$, are denoted τ and τ_s . As noted above, these traces are constructed using the Berezin integral, [G83, BGV92]. This is why only the tangential volume form dx_F appears in the formulas below.

Note that the actions of $a \in C_c^\infty(\mathcal{G}; \wedge^* r^*(\nu^*))$ are smoothing along the leaves of F and \widehat{F} , which is why we may restrict our attention to operators with symbols in $SC_{\text{rh}}^\infty(M, E)$ and $SC_{\text{rh}}^\infty(\widehat{M}, \widehat{E})$.

Now we invoke results of Dominguez, [D98], and Alvarez Lopéz, [AL92], which simplify things considerably. Recall, [CM95], that the cohomology classes of the cocycles (ϕ_k) and $(\widehat{\phi}_k)$ do not depend on the metric on M or the induced metric on \widehat{M} . Dominguez has shown that it is always possible to choose the metric on M (so also on \widehat{M}) so that $c(\mu)$ is a basic form, so Ξ is also, (and in fact $\Xi = d_\nu c(\mu) = dc(\mu)$). Thus, if we use this metric, we have $\widetilde{c}(\mu) = \Xi = 0$. Alvarez Lopéz showed that the basic component of the mean curvature vector field of any Riemannian foliation is closed. Thus $\Xi = 0$ for the Dominguez metric. Similarly for the corresponding operators on $(\widehat{M}, \widehat{F})$. Then Lemma 5.8 immediately gives

$$\text{if } \ell \neq (0, \dots, 0), \quad a^{k, \ell}(tD)_{1/t} = f_{k, \ell}(t), \quad \text{and} \quad a^{k, 0}(tD)_{1/t} = a^{k, 0}(t\widehat{D})_{1/t} = a_0 d_\nu a_1 \cdots d_\nu a_k.$$

As noted above, each term of $f_{k, \ell}$, $\ell \neq (0, \dots, 0)$, can be written as a sum of compositions of an element of $C_c^\infty(\mathcal{G}; \wedge^* r^*(\nu^*))$ and a $\text{hA}\Psi\text{DO}$ (actually differential operator) whose leading symbol is 0. More specifically, denote by e_1, \dots, e_q a local orthonormal basis of ν , and denote by ∇_J a (local) operator which is a finite composition (in any order and with repetitions allowed) of $\nabla_{e_1}^\nu, \dots, \nabla_{e_q}^\nu$ and ∇_μ . Then each term of $f_{k, \ell}$ may be written locally as $t^{b_J} a_J \nabla_J$ for some $a_J \in C_c^\infty(\mathcal{G}; \wedge^* r^*(\nu^*))$, independent of t , and $b_J \geq 1$. Similarly for $a^{k, \ell}(t\widehat{D})_{1/t}$ and the induced operator $t^{b_J} a_J \widehat{\nabla}_J$ on $\widehat{\mathcal{H}}$. Then

$$a^{k, \ell}(tD) = t^{k+2|\ell|} \sum_J a_J \nabla_J = \sum_J (a_J)_t t^{b_J} (\nabla_J)_t,$$

and

$$\mathbf{6.2.} \quad a^{k, \ell}(t\widehat{D}) = t^{k+2|\ell|} \sum_J a_J \widehat{\nabla}_J = \sum_J (a_J)_t t^{b_J} (\widehat{\nabla}_J)_t.$$

Note that both $t^{b_J} (\nabla_J)_t$ and $t^{b_J} (\widehat{\nabla}_J)_t$ are transverse asymptotic differential operators with leading symbols 0.

Proposition 6.3. *For all $t > 0$,*

$$\text{Tr}(a^{k, \ell}(tD)e^{-t^2 D^2}) = (2\pi)^{-q} \int_{\nu^*} t^{k+2|\ell|} \sum_J \sum_{\gamma \in \mathcal{G}_x^x} (a_J)_{1/t}(\gamma^{-1}) \text{Tr} \left(\varsigma(\nabla_J e^{-t^2 D^2})_{1/t}(x, \eta) \right) d\eta dx_F,$$

and

$$\text{Tr}_s(a^{k, \ell}(tD)e^{-t^2 D^2}) = (2\pi)^{-q} \int_{\nu^*} t^{k+2|\ell|} \sum_J \sum_{\gamma \in \mathcal{G}_x^x} (a_J)_{1/t}(\gamma^{-1}) \text{Tr}_s \left(\varsigma(\nabla_J e^{-t^2 D^2})_{1/t}(x, \eta) \right) d\eta dx_F.$$

The corresponding results hold for $\tau(a^{k, \ell}(t\widehat{D})e^{-t^2 \widehat{D}^2})$ and $\tau_s(a^{k, \ell}(t\widehat{D})e^{-t^2 \widehat{D}^2})$, where ν^ is replaced by $\widehat{\nu}^*|M \subset \widehat{M}$, and $\sum_{\gamma \in \mathcal{G}_x^x} (a_J)_{1/t}(\gamma^{-1})$ is replaced by $\sum_{\widehat{\gamma} \in \widehat{\mathcal{G}}_x^x} (a_J)_{1/t}((\rho \circ \widehat{\gamma})^{-1})$.*

Proof. We only do the proof for $\text{Tr}_s(a^{k, \ell}(tD)e^{-t^2 D^2})$, as the other three cases are quit similar.

First we prove it without applying the rescaling operator $(\cdot)_{1/t}$. From [BH17b] we have

Lemma 6.4. *For all $t > 0$, $a_J \nabla_J e^{-t^2 D^2}$ is a smoothing operator on \mathcal{H} . Similarly for $a_J \widehat{\nabla}_J e^{-t^2 \widehat{D}^2}$ acting on $\widehat{\mathcal{H}}$.*

Thus we have

$$\mathrm{Tr}_s(a_J \nabla_J e^{-t^2 D^2}) = \int_M \mathrm{Tr}_s(K(x, x)) dx_F,$$

where K is the Schwartz kernel of $a_J \nabla_J e^{-t^2 D^2}$, and we use the Berezin integral for spinors, [G83, BGV92].

Denote by δ_x , δ_0 , δ_0^F , and δ_x^L the Dirac delta operators at $x \in M$, $0 \in T^*M_x$, $0 \in T^*F_x$, and $x \in L_x$, respectively. For $u_x \in (\mathcal{S}_\nu \otimes E)_x$,

$$K(x, x)u_x = \left(a_J \nabla_J e^{-t^2 D^2} (\delta_x(x') \mathcal{T}_{x, x'}(u_x)) \right)(x) = \left(a_J \nabla_J e^{-t^2 D^2} (\delta_0(\widehat{\mathrm{exp}}_x^{-1}(x')) \alpha(x, x') \mathcal{T}_{x, x'}(u_x)) \right)(x).$$

Now the Fourier Transform of δ_0 is the constant function 1. So

$$\begin{aligned} \mathrm{Tr}_s(K(x, x)) &= \mathrm{Tr}_s \left((a_J \nabla_J e^{-t^2 D^2} [(2\pi)^{-n} \int_{T^*M_x} e^{i\langle \widehat{\mathrm{exp}}_x^{-1}(x'), \xi \rangle} \alpha(x, x') \mathcal{T}_{x, x'} d\xi]) (x) \right) = \\ &(2\pi)^{-q} \int_{\nu_x^*} \mathrm{Tr}_s \left((a_J \nabla_J e^{-t^2 D^2} [(2\pi)^{-p} \int_{T^*F_x} e^{i\langle \widehat{\mathrm{exp}}_x^{-1}(x'), (\zeta, 0) \rangle} e^{i\langle \widehat{\mathrm{exp}}_x^{-1}(x'), (0, \eta) \rangle} \alpha(x, x') \mathcal{T}_{x, x'} d\zeta]) (x) \right) d\eta, \end{aligned}$$

where we are able to move the integration outside by the argument in the proof of Theorem 3.7 of [G83]. Since the symbol of $\nabla_J e^{-t^2 D^2}$ depends only on η , we may write this as

$$\begin{aligned} (2\pi)^{-q} \int_{\nu_x^*} \mathrm{Tr}_s \left((a_J [(2\pi)^{-p} \int_{T^*F_x} e^{i\langle \widehat{\mathrm{exp}}_x^{-1}(x'), (\zeta, 0) \rangle} d\zeta \nabla_J e^{-t^2 D^2} (e^{i\langle \widehat{\mathrm{exp}}_x^{-1}(x'), (0, \eta) \rangle} \alpha(x, x') \mathcal{T}_{x, x'})]) (x) \right) d\eta = \\ (2\pi)^{-q} \int_{\nu_x^*} \mathrm{Tr}_s \left(\int_{\mathcal{G}_x} a_J(\gamma^{-1}) \right. \\ \left. h_{\gamma^{-1}} \left[\left((2\pi)^{-p} \int_{T^*F_x} e^{i\langle \widehat{\mathrm{exp}}_x^{-1}(x'), (\zeta, 0) \rangle} d\zeta \nabla_J e^{-t^2 D^2} (e^{i\langle \widehat{\mathrm{exp}}_x^{-1}(x'), (0, \eta) \rangle} \alpha(x, x') \mathcal{T}_{x, x'}) \right) (r(\gamma)) \right] d\gamma \right) d\eta. \end{aligned}$$

Now, at $\gamma \in \mathcal{G}_x$,

$$\left((2\pi)^{-p} \int_{T^*F_x} e^{i\langle \widehat{\mathrm{exp}}_x^{-1}(x'), (\zeta, 0) \rangle} d\zeta \right) (r(\gamma)) = \delta_0^F((\widehat{\mathrm{exp}}_x|_{L_x})^{-1}(r(\gamma))) \alpha(x, r(\gamma)) = \delta_x^L(r(\gamma)).$$

Thus,

$$\begin{aligned} \mathrm{Tr}_s(K(x, x)) &= (2\pi)^{-q} \int_{\nu_x^*} \mathrm{Tr}_s \left(\int_{\mathcal{G}_x} \delta_x^L(r(\gamma)) a_J(\gamma^{-1}) \right. \\ &\quad \left. h_{\gamma^{-1}} \left[\nabla_J e^{-t^2 D^2} \left((e^{i\langle \widehat{\mathrm{exp}}_x^{-1}(x'), (0, \eta) \rangle} \alpha(x, x') \mathcal{T}_{x, x'})|_{L_x} \right) (r(\gamma)) \right] d\gamma \right) d\eta = \\ &(2\pi)^{-q} \int_{\nu_x^*} \sum_{\gamma \in \mathcal{G}_x} \mathrm{Tr}_s \left(a_J(\gamma^{-1}) h_{\gamma^{-1}} \left[\varsigma(\nabla_J e^{-t^2 D^2})(x, \eta) \right] \right) d\eta = \\ &(2\pi)^{-q} \int_{\nu_x^*} \sum_{\gamma \in \mathcal{G}_x} a_J(\gamma^{-1}) \mathrm{Tr}_s \left(\varsigma(\nabla_J e^{-t^2 D^2})(x, \eta) \right) d\eta, \end{aligned}$$

since $r(\gamma) = x$ for $\gamma \in \mathcal{G}_x$, and $h_{\gamma^{-1}} \left[\varsigma(\nabla_J e^{-t^2 D^2})(x, \eta) \right] = \varsigma(\nabla_J e^{-t^2 D^2})(x, \eta)$, as this symbol is invariant under the holonomy action. We thus have

$$\mathrm{Tr}_s(a_J \nabla_J e^{-t^2 D^2}) = (2\pi)^{-q} \int_{\nu_x^*} \sum_{\gamma \in \mathcal{G}_x} a_J(\gamma^{-1}) \mathrm{Tr}_s \left(\varsigma(\nabla_J e^{-t^2 D^2})(x, \eta) \right) d\eta dx_F.$$

When we apply the rescaling operator $(\cdot)_{1/t}$, the integral is unchanged just as in [G83] p. 175. Replacing η by η/t gives a factor of t^q when doing the integration in η . The only part of $(a_J)_{1/t}(x) \varsigma(\nabla_J e^{-t^2 D^2})_{1/t}$ which contributes comes from $C^\infty(\wedge^q \nu^*) \otimes_{C^\infty(M)} C^\infty(T^*M, \mathrm{End}(E))$, since we are using the Berezin integral, and this contributes the factor t^{-q} , which cancels the t^q .

The result now follows by summing over J and multiplying by $t^{k+2|\ell|}$. \square

Now suppose that $\ell \neq (0, \dots, 0)$. Then

$$\int_0^\epsilon 2t^{2z-1} e^{-t^2} \text{Tr}_s(a^{k,\ell}(tD)e^{-t^2 D^2}) dt = \int_0^\epsilon 2t^{2z-1} e^{-t^2} (2\pi)^{-q} \int_{\nu^*} t^{k+2|\ell|} \sum_J \sum_{\gamma \in \mathcal{G}_x^J} (a_J)_{1/t}(\gamma^{-1}) \text{Tr}_s \left(\varsigma(\nabla_J e^{-t^2 \widehat{D}^2})(x, \eta)_{1/t} \right) d\eta dx_F dt.$$

Set

$$f(t) = e^{-t^2} (2\pi)^{-q} \int_{\nu^*} t^{k+2|\ell|} \sum_J \sum_{\gamma \in \mathcal{G}_x^J} (a_J)_{1/t}(\gamma^{-1}) \text{Tr}_s \left(\varsigma(\nabla_J e^{-t^2 \widehat{D}^2})(x, \eta)_{1/t} \right) d\eta dx_F.$$

The results in Section 5 imply that $f(t)$ is continuous at $t = 0$. So

$$\tau_0(\gamma a^{k,\ell}(D)(I+D)^{-k/2-|\ell|}) = \Gamma(k/2 + |\ell|)^{-1} \text{Res}_{z=0} \left[\int_0^\epsilon 2t^{2z-1} f(t) dt \right],$$

which is approximated by (and, because it is independent of ϵ , actually equals)

$$\begin{aligned} \Gamma(k/2 + |\ell|)^{-1} \text{Res}_{z=0} \left[\int_0^\epsilon 2t^{2z-1} f(0) dt \right] &= \Gamma(k/2 + |\ell|)^{-1} \text{Res}_{z=0} \left[\frac{t^{2z}}{z} f(0) \Big|_0^\epsilon \right] = \\ &= \Gamma(k/2 + |\ell|)^{-1} \text{Res}_{z=0} \left[\frac{\epsilon^{2z}}{z} f(0) \right] = \Gamma(k/2 + |\ell|)^{-1} \lim_{t \rightarrow 0} f(t) = \\ &= \Gamma(k/2 + |\ell|)^{-1} \lim_{t \rightarrow 0} \left[e^{-t^2} (2\pi)^{-q} \int_{\nu^*} t^{k+2|\ell|} \sum_J \sum_{\gamma \in \mathcal{G}_x^J} (a_J)_{1/t}(\gamma^{-1}) \text{Tr}_s \left(\varsigma(\nabla_J e^{-t^2 \widehat{D}^2})(x, \eta)_{1/t} \right) d\eta dx_F \right]. \end{aligned}$$

By the results in Section 5, we may interchange the limit and the integration to get

$$\Gamma(k/2 + |\ell|)^{-1} (2\pi)^{-q} \int_{\nu^*} \lim_{t \rightarrow 0} \left[e^{-t^2} t^{k+2|\ell|} \sum_J \sum_{\gamma \in \mathcal{G}_x^J} (a_J)_{1/t}(\gamma^{-1}) \text{Tr}_s \left(\varsigma(\nabla_J e^{-t^2 \widehat{D}^2})(x, \eta)_{1/t} \right) \right] d\eta dx_F.$$

By Theorem 6.1 of [BH17a], Equation 6.2, and the results of Section 5,

$$\lim_{t \rightarrow 0} t^{k+2|\ell|} \sum_J (a_J)_{1/t} \varsigma(\nabla_J e^{-t^2 \widehat{D}^2})(x, \eta)_{1/t} = \lim_{t \rightarrow 0} \sum_J a_J e^{-\frac{1}{4}\Omega_\nu(\partial/\partial\eta, \partial/\partial\eta')} p_0^J(x, \eta) \wedge q_0(x, \eta')|_{\eta'=\eta},$$

where $p_0^J(x, \eta)$ is the leading symbol of $t^{b_J} \nabla_J$ and $q_0(x, \eta')$ is the leading symbol of $e^{-t^2 \widehat{D}^2}$, with η' substituted for η . Since $b_J \geq 1$, $p_0^J(x, \eta) = 0$, and

$$\lim_{t \rightarrow 0} t^{k+2|\ell|} \sum_J (a_J)_{1/t} \varsigma(\nabla_J e^{-t^2 \widehat{D}^2})(x, \eta)_{1/t} = 0.$$

Thus, if $\ell \neq (0, \dots, 0)$,

$$\tau_0(\gamma a^{k,\ell}(D)(I+D^2)^{-k/2-|\ell|}) = 0.$$

For the case $k \neq 0$, and $\ell = (0, \dots, 0)$, we have $t^k \sum_J (a_J)_{1/t} = a^{k,0}(tD)_{1/t} = a_0 d_\nu a_1 \cdots d_\nu a_k$. The results in Section 5 give

$$\lim_{t \rightarrow 0} \text{Tr}_s \left(\varsigma(e^{-t^2 \widehat{D}^2})_{1/t} \right) = e^{-\frac{1}{2}\Xi} \wedge e^{-|\eta|^2 - \frac{1}{2}\Omega_E + \frac{1}{16}\Omega_\nu}(1) = e^{-|\eta|^2 - \frac{1}{2}\Omega_E + \frac{1}{16}\Omega_\nu}(1),$$

since $\Xi = 0$ for the Dominguez metric. Proposition 5.13 then finishes the proof for this case.

For the case $k = 0$ we have, following [CM95], p. 230,

$$\begin{aligned} \phi_0(a_0) &= \text{Res}_{z=0} \left[\text{Tr}_s(z^{-1} a_0 (I+D^2)^{-z}) \right] = \text{Res}_{z=0} \left[\text{Tr}_s((z\Gamma(z))^{-1} a_0 \int_0^\infty t^{z-1} e^{-t(I+D^2)} dt) \right] = \\ &= \text{Res}_{z=0} \left[\text{Tr}_s(a_0 \int_0^\infty t^{z-1} e^{-t(I+D^2)} dt) \right], \end{aligned}$$

since $z\Gamma(z) = \Gamma(z+1)$, so $z\Gamma(z)|_{z=0} = 1$. Then proceed just as in the case $k > 0$ to get

$$\phi_0(a_0) = \int_M \sum_{\mathcal{G}_x^x} a_0(\gamma^{-1}) \text{ch}(E) \wedge \widehat{A}(\nu^*) \wedge dx_F.$$

To complete the proof of Theorem 3.8, we have the following.

Proposition 6.5. *Denote the cohomology classes of the sequences $(\phi_k)_{k \geq 0}$, and $(\widehat{\phi}_k)_{k \geq 0}$, k even, by $\text{ch}(M, F)$ and $\text{ch}(\widehat{M}, \widehat{F})$ respectively. Then the following diagram commutes on the image of the Baum-Connes assembly map in $K_*(C_{\max}^*(\mathcal{G}))$.*

$$\begin{array}{ccc} K_*(C_{\max}^*(\mathcal{G})) & \xrightarrow{\text{ch}(M, F)} & \mathbb{Z} \\ & \searrow \text{ch}(\widehat{M}, \widehat{F}) & \swarrow \\ & \mathbb{R} & \end{array}$$

Proof. If X is in the image of the Baum-Connes assembly map, it can be represented by $[e] - [e']$, where e and e' are projections in $M_\infty(\widetilde{C}_c^\infty(\mathcal{G}))$ with the supports of the corresponding elements of $M_\infty(\widetilde{C}_c^\infty(\mathcal{G}))$ in an arbitrarily small neighborhood of the units $\mathcal{G}^{(0)} \sim M \subset \mathcal{G}$. See Remark 3.9. The tilde indicates that a unit has been added. Then the supports of the elements $e(d_\nu e)^k$ and $e'(d_\nu e')^k$ are also contained in an arbitrarily small neighborhood of $\mathcal{G}_0 \sim M \subset \mathcal{G}$.

Apply the first part of Theorem 3.8 to get

$$\langle \text{ch}(M, F), X \rangle = \frac{1}{k!} \int_M \sum_{\gamma \in \mathcal{G}_x^x} [\text{tr}(e(d_\nu e)^k) - \text{tr}(e'(d_\nu e')^k)] (\gamma^{-1}) \wedge \text{ch}(E) \wedge \widehat{A}(\nu^*) \wedge dx_F,$$

and

$$\langle \text{ch}(\widehat{M}, \widehat{F}), X \rangle = \frac{1}{k!} \int_{M \subset \widehat{M}} \sum_{\widehat{\gamma} \in \widehat{\mathcal{G}}_x^x} [\text{tr}(e(d_\nu e)^k) - \text{tr}(e'(d_\nu e')^k)] ((\rho \circ \widehat{\gamma})^{-1}) \wedge \text{ch}(\widehat{E}) \wedge \widehat{A}(\widehat{\nu}^*) \wedge dx_F.$$

Because of the restrictions on the supports, both sums collapse to just the constant path at x , so the two expressions are equal. \square

7. STATEMENT OF THE THEOREM FOR TRANSVERSE SPECTRAL TRIPLES

In this section we will be content to give the statement of the theorem for two transverse spectral triples associated to F . The proof of Theorem 3.8 can be easily extended, using the Getzler calculus [G83], to prove Theorem 7.1 below. The advantage of using these spectral triples is that we do not have to assume that the normal bundle of F is integrable.

The transverse spectral triples are given as follows.

1. The Type I Connes-Moscovici spectral triple: $(C_c^\infty(\mathcal{G}_T^T), \mathcal{H}_T, D_T)$. T is a complete transversal for the foliation F , and the trace is the usual super trace, constructed using the Berezin integral, for operators on $\mathcal{H}_T = L^2(T, (\mathcal{S}_\nu \otimes E)|_T)$. $\mathcal{G}_T^T \subset \mathcal{G}$ is the sub groupoid of elements which start and end in T . D_T can be taken to be either D or \widetilde{D} , restricted to T .

This transverse spectral triple is Morita equivalent to the Type I spectral triple given in Section 3, and so is an even spectral triple with simple dimension spectrum contained in the set $\{k \in \mathbb{N} \mid k \leq q\}$. See [K97].

2. The Type II Atiyah spectral triple: $(C_c^\infty(\mathcal{G}_T^T), B(\widehat{\mathcal{H}}_T)^\Gamma, \widehat{D}_T)$. $\rho : (\widehat{M}, \widehat{F}) \rightarrow (M, F)$ is a (possibly non-compact) Galois foliation cover of M , with covering group Γ . The transversal \widehat{T} is the inverse image of T . \widehat{D}_T and $\widehat{\mathcal{S}}_\nu \otimes \widehat{E}$ are the pull backs of D_T and $\mathcal{S}_\nu \otimes E$. $\widehat{\mathcal{H}}_T = L^2(\widehat{T}, (\widehat{\mathcal{S}}_\nu \otimes \widehat{E})|_{\widehat{T}})$. $B(\widehat{\mathcal{H}}_T)^\Gamma$ is the algebra of bounded operators on $\widehat{\mathcal{H}}_T$ which are Γ -invariant. The trace is the usual super trace, constructed using the Berezin integral, restricted to a fundamental domain of T in \widehat{T} , denoted $T \subset \widehat{T}$.

As in the main result: the algebra is the same as in the Type I case; $B(\widehat{\mathcal{H}}_{\widehat{T}})^{\Gamma}$ is the analog of the von Neuman algebra used by Atiyah in the L^2 covering index theorem; $\widehat{D}_{\widehat{T}}$ is Γ -invariant; for a non-compact covering foliation, this is Morita equivalent to the Type II spectral triple given in Section 3, and so is an even spectral triple with simple dimension spectrum contained in the set $\{k \in \mathbb{N} \mid k \leq q\}$; for a compact covering, this is a slight modification of the Type I Connes-Moscovici spectral triple above.

Theorem 7.1. *Let $\rho : (\widehat{M}, \widehat{F}) \rightarrow (M, F)$ be a Galois covering foliation of (M, F) . Assume that F has even codimension and is transversely spin. Let $a_0, \dots, a_k \in C_c^\infty(\mathcal{G}_T^T)$. The sequences $(\phi_k)_{k \geq 0}$, and $(\widehat{\phi}_k)_{k \geq 0}$, k even, where*

$$\phi_k(a_0, \dots, a_k) = \frac{1}{k!} \int_T \sum_{\gamma \in \mathcal{G}_x^T} (a_0 d_\nu a_1 \cdots d_\nu a_k)(\gamma^{-1}) \wedge \text{ch}(E) \wedge \widehat{A}(\nu^*),$$

and

$$\widehat{\phi}_k(a_0, \dots, a_k) = \frac{1}{k!} \int_{T \subset \widehat{T}} \sum_{\widehat{\gamma} \in \widehat{\mathcal{G}}_x^T} (a_0 d_\nu a_1 \cdots d_\nu a_k)((\rho \circ \widehat{\gamma})^{-1}) \wedge \text{ch}(\widehat{E}) \wedge \widehat{A}(\widehat{\nu}^*)$$

are cocycles in the (b, B) bicomplex for $C_c^\infty(\mathcal{G}_T^T)$.

The cohomology classes of (ϕ_k) and $(\widehat{\phi}_k)$ are the Connes-Chern characters of $(C_c^\infty(\mathcal{G}_T^T), \mathcal{H}_T, D_T)$ and $(C_c^\infty(\widehat{\mathcal{G}}_T^T), B(\widehat{\mathcal{H}}_{\widehat{T}})^{\Gamma}, \widehat{D}_{\widehat{T}})$, respectively. These classes induce the same map on the image of the Baum-Connes assembly map in $K_*(C_{\max}^*(\mathcal{G}_T^T))$, [BC00]. That is, we have an Atiyah L^2 covering index theorem for these spectral triples associated to foliation coverings.

Finally, we point out that the above Morita reduction method of restriction to a complete transversal is not generally accessible when the action of Γ is not free.

Example 7.2. *Recall the spectral triple*

$$(C_c^\infty(\mathcal{G}) \rtimes \Gamma, \mathcal{N} \subset B(M, E \otimes \wedge \nu^*) \otimes l^2\Gamma, D_E \rtimes \Gamma)$$

from [BH17b]. Here Γ is a countable group of diffeomorphisms acting properly, but not freely, on the possibly non-compact manifold M , preserving a foliation F with normal bundle ν^* and homotopy graph \mathcal{G} . M/Γ is assumed to be a compact space. The bundle E is basic, Γ equivariant, and Hermitian. The operator D_E is a twisted transverse Dirac operator for F , and \mathcal{N} is a certain von Neumann algebra. Given $\phi \in C_c^\infty(\widehat{\mathcal{G}}) \rtimes \Gamma$, $\phi(g) \in C_c^\infty(\widehat{\mathcal{G}})$, $g \in \Gamma$, also acts on sections of $E \otimes \wedge^* \nu^*$. Finally, the trace TR used is: for certain $A \in \mathcal{N}$, and all $\phi \in C_c^\infty(\widehat{\mathcal{G}}) \rtimes \Gamma$, $\phi \circ A$ is TR trace class with

$$\text{TR}(\phi \circ A) = \sum_{g \in \Gamma} \text{Tr}(g^{-1} \phi(g) \circ A_{g,e}),$$

where Tr is just the usual trace.

As Γ does not act freely, it is not possible in general to reduce the associated index problem to the corresponding one on a complete transversal. Indeed, a Γ -equivariant complete transversal does not always exist [BH17b], so the Morita reduction method to a transversal is not available, and we must use our global constructions.

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INSTITUT MONTPELLIERAIN ALEXANDER GROTHENDIECK, UMR 5149 DU CNRS, UNIVERSITÉ DE MONTPELLIER
E-mail address: moulay.benameur@umontpellier.fr

MATHEMATICS, STATISTICS, AND COMPUTER SCIENCE, UNIVERSITY OF ILLINOIS AT CHICAGO
E-mail address: heitsch@uic.edu